

**Микро- и наноэлектромеханические системы  
на базе  
сегнетэлектрических структур  
и методы их исследования**

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# Outline

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- **Introduction**
- **Overview of MEMS**
- **Basics of piezoelectric effect in bulk materials and thin films**
- **Piezoelectric measurements: Laser Interferometry**
- **Piezoelectric measurements: Piezoresponse Force Microscopy**
- **Relaxor ferroelectrics: nanoscale vs. macroscopic response**
- **Summary**

# Where to read

## *Reports on Progress in Physics* Sept., 2006

Nanoscale ferroelectrics: processing, characterization and future trends

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### Abstract

This review paper summarizes recent advances in the quickly developing field of nanoscale ferroelectrics, analyses its current status, and considers potential future developments. The paper presents a brief survey of the fabrication methods of ferroelectric nanostructures and investigation of the size effects by means of scanning probe microscopy. One of the focuses of the review will be the study of kinetics of nanoscale ferroelectric switching in inhomogeneous electrical and elastic fields. Another emphasis will be made on tailoring the electrical and mechanical properties of ferroelectrics with a viewpoint of fabrication of nanoscale domain structures.

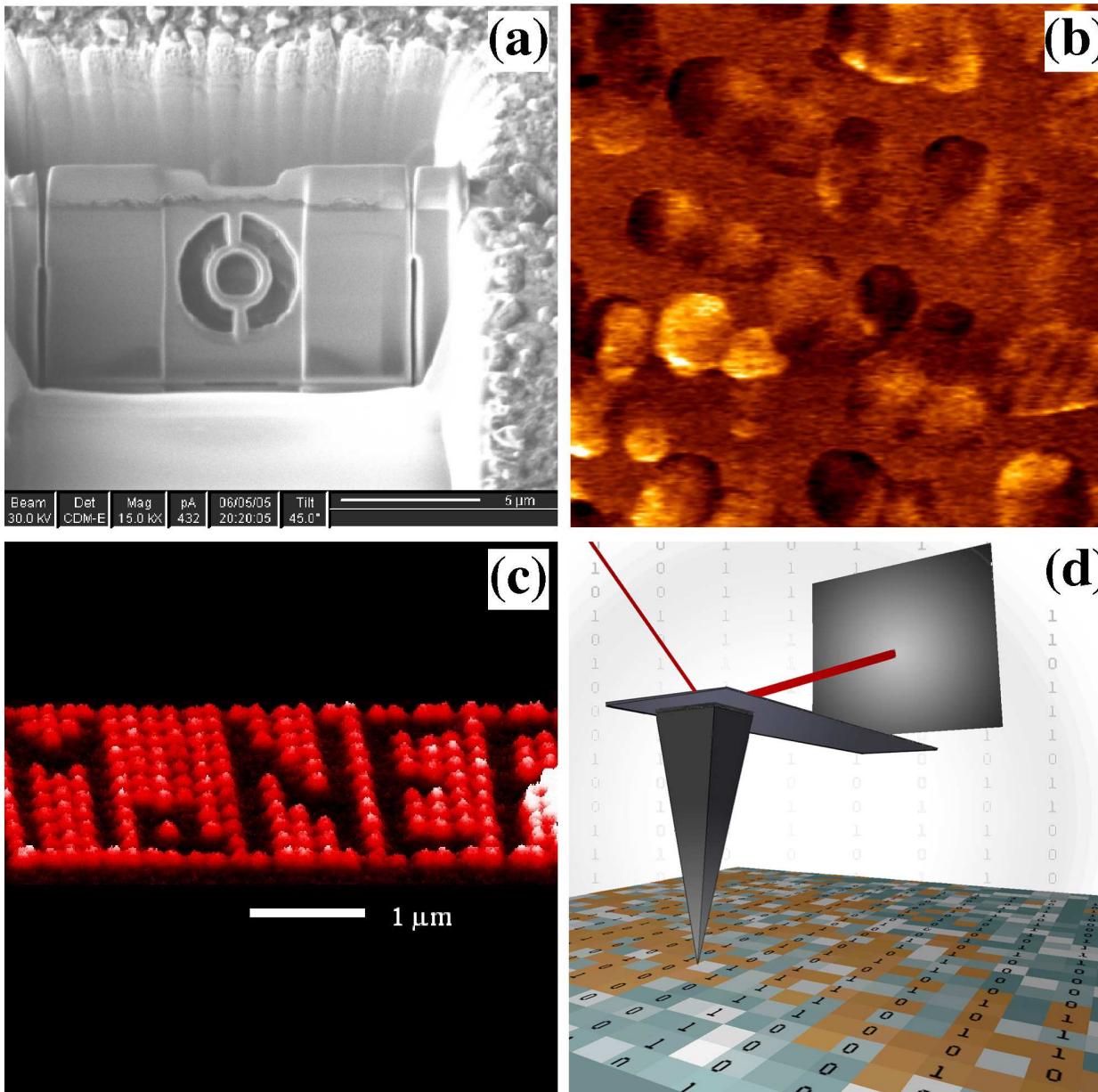
**“Scanning Probe Microscopy: Electrical and Electromechanical Phenomena at the Nanoscale”, Springer, 2006.**

Review of Ferroelectric Domain Imaging by Piezoresponse Force Microscopy

A. L. KHOLKIN, S. V. KALININ, A. ROELOFS, AND A. GRUVERMAN

This chapter describes the principles, theoretical background, recent developments, and applications of a local probe-based technique for nondestructive high-resolution ferroelectric domain imaging and manipulation—piezoresponse force microscopy (PFM). This technique has proven to be a powerful tool for the characterization of ferroelectric thin films, ceramics, and single crystals. Recent advances in application of PFM for studying a mechanism of polarization reversal at the nanoscale, domain dynamics, degradation effects, and size-dependent phenomena in ferroelectrics are reviewed in detail. Examples of using PFM for the characterization of various polar materials such as ferroelectric films, piezoelectric semiconductors, and ferroelectric relaxors are given.

# *Special Issue on "Nanoscale Ferroelectrics", TUFFC, Dec. 06*



# *Actuation principles used in MEMS*

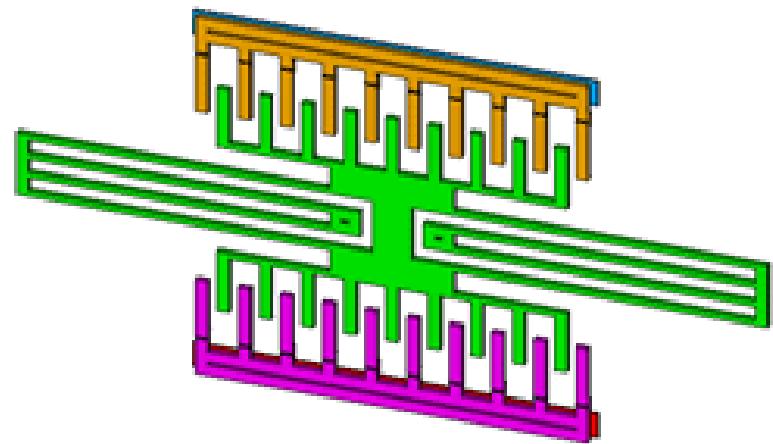
- Piezoelectric effect
  - strong force, low displacement, good scalability, hard to integrate, high efficiency, fast response
- Electrostatic effect
  - weak force, low to moderate displacement, non-linearity, most easy to integrate, large voltages even for small gaps
- Electromagnetic effect
  - high displacement, low to moderate force, limited by heating of coils, slow response, hard to integrate
- Thermal effect
  - strong force, large displacement, slow, hard to control, low efficiency

# *Other principles...*

- Shape memory effect  
(austenite/martensite phase transition)
  - 5% strain (NiTi alloy),
  - slow response
- Piezoresistivity
  - easily processed
  - sensing only
- Magnetostriction
  - processing can be complicated (Terfenol-D)
  - similar performance to piezoelectrics

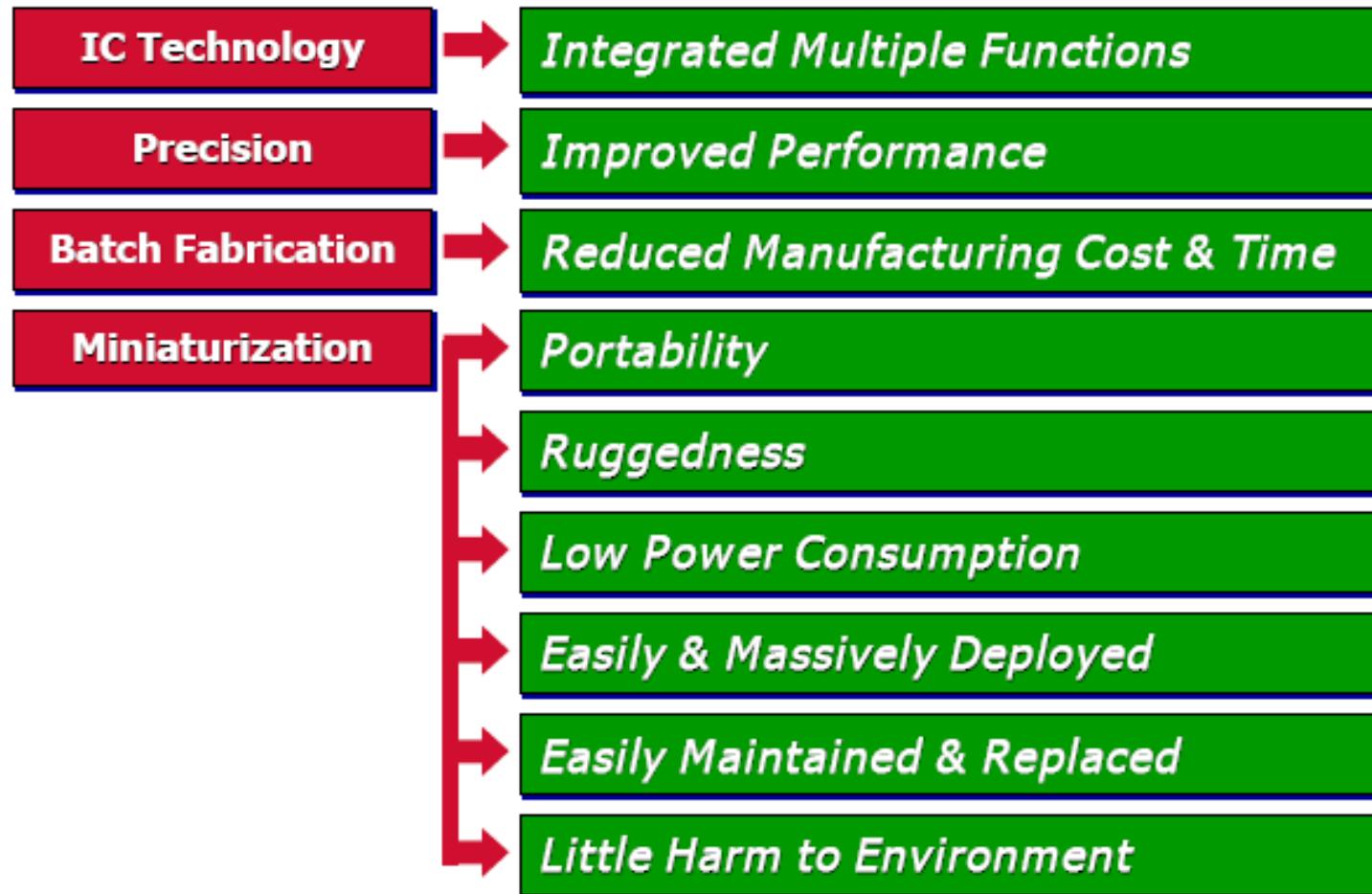
# *Possibilities provided by Si technology*

- Most common structures
  - Electrostatic drive
  - Capacitive sensor
- Complex geometries
  - Large strain with moderate force
  - High sensitivity-large areas
- Easy processing
  - 30+ years of Si processing technology
  - Surface micromachining
  - Bulk surface micromachining
  - LIGA process, etc...



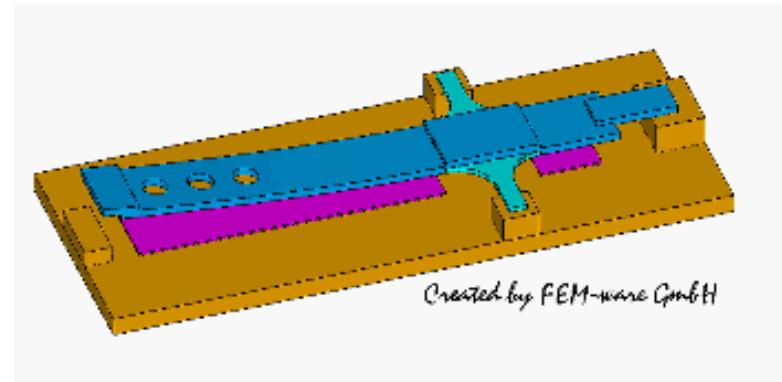
**Schematic of linear electrostatic drive**

## *Advantages of MEMS*

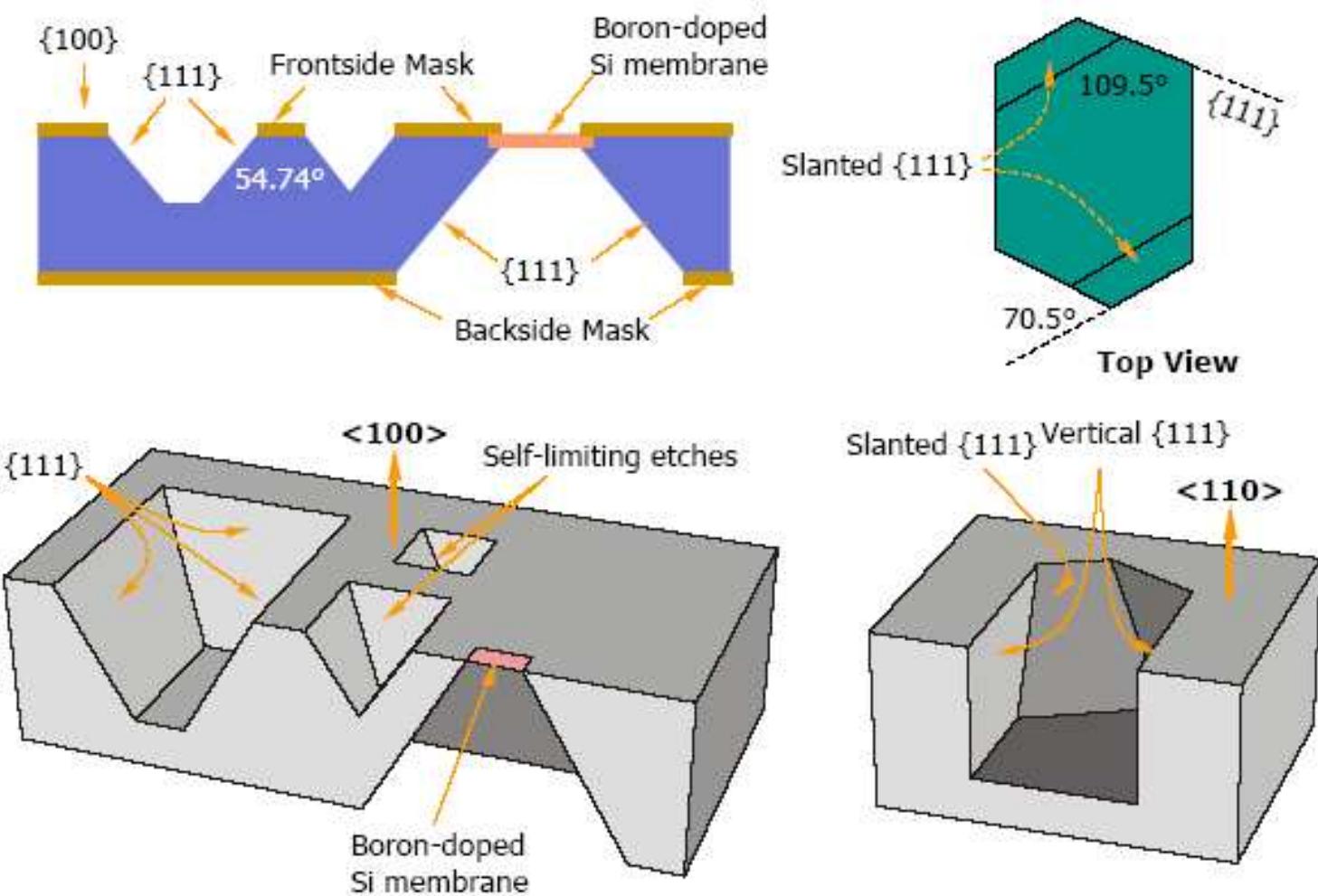


# *MEMS technologies*

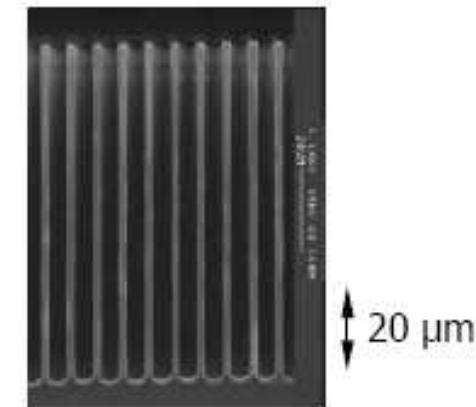
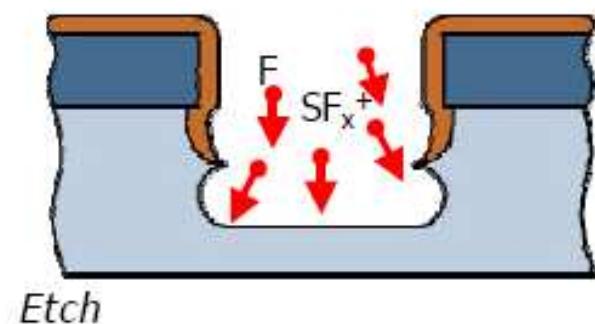
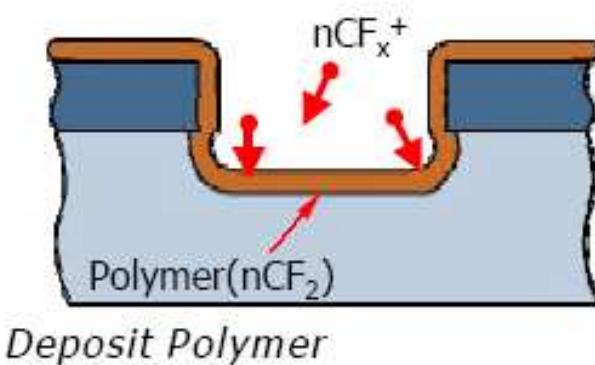
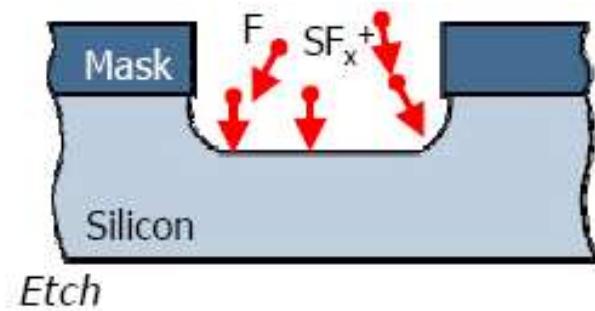
- Bulk micromachining
- Surface micromachining
- Wafer bonding
- LIGA and similar techniques
- Others
  - MicroEDM
  - 3D lithography
  - Laser micromachining



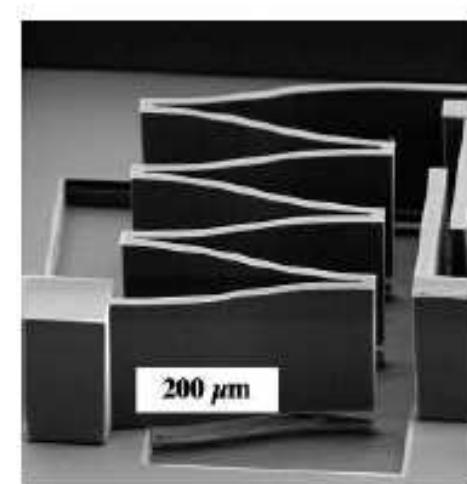
# *Wet etching of Si*



## Dry etching of Si

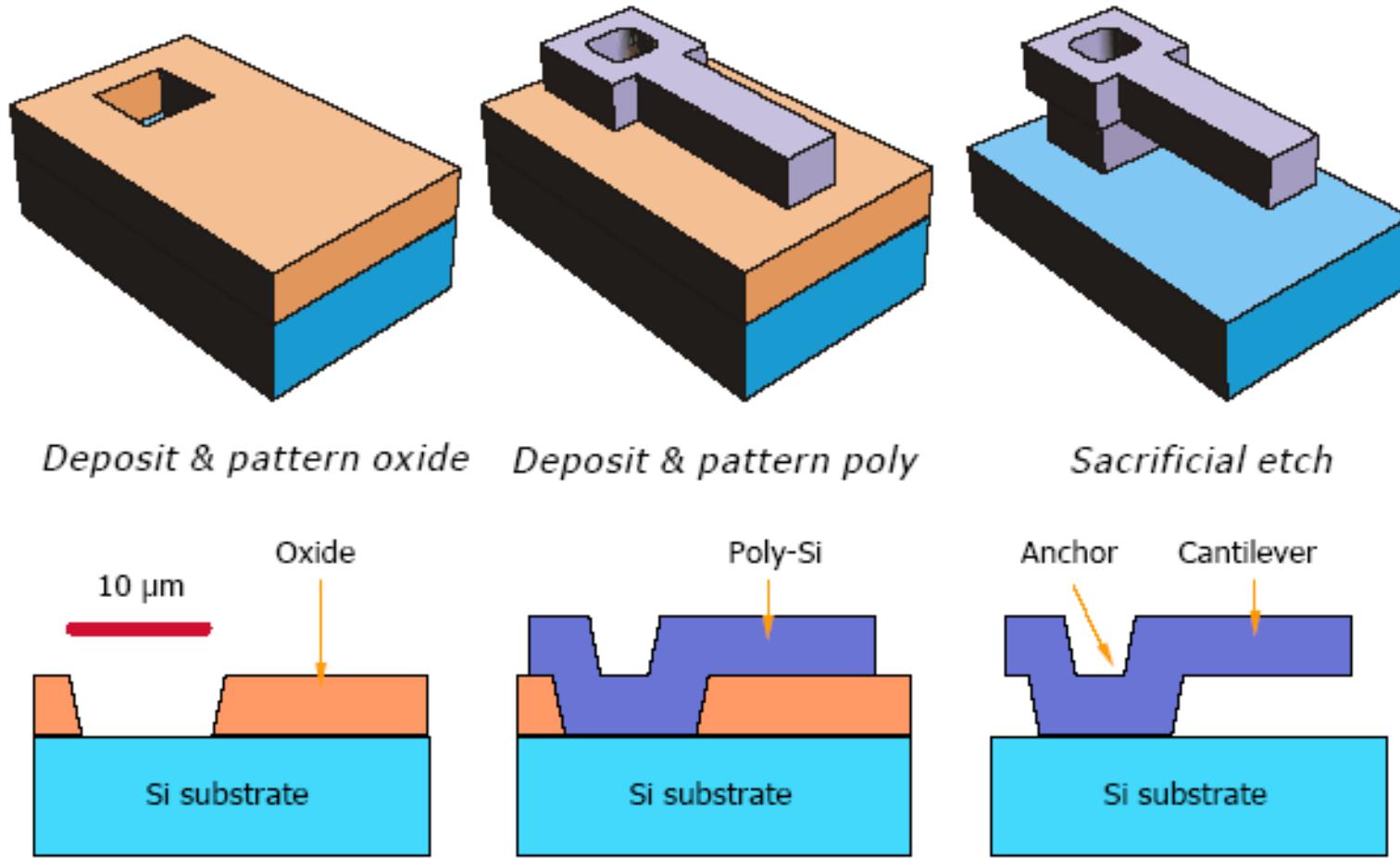


Trenches - Surface Technology Systems



Spring - Klaassen, et al, 1995

# *Surface micromachining*



# Piezoelectric micromotor (EPFL)

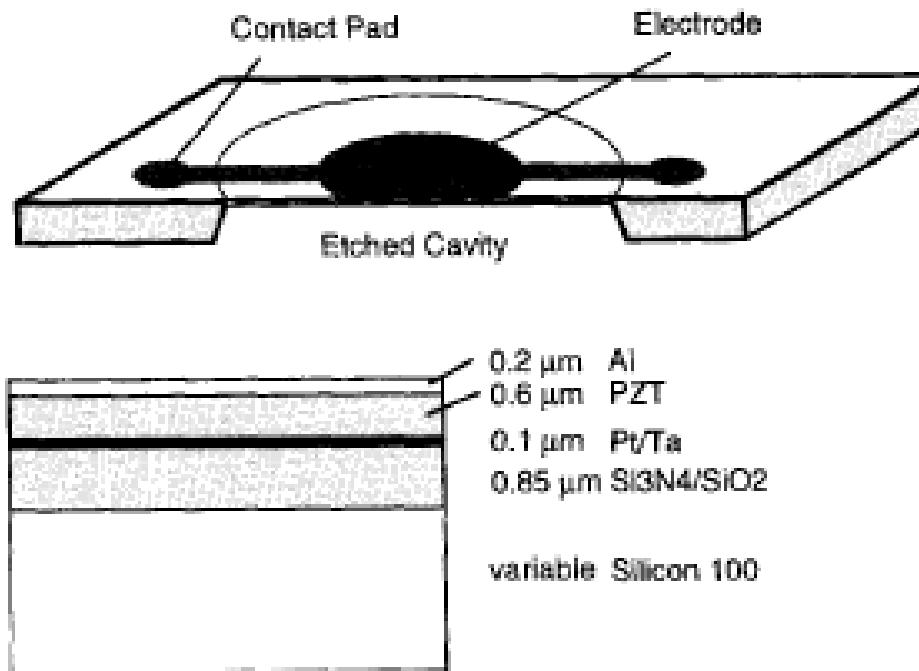


Fig. 1. Schematic drawing of the membrane with view onto the top side of the wafer and a schematic cross section through the films.

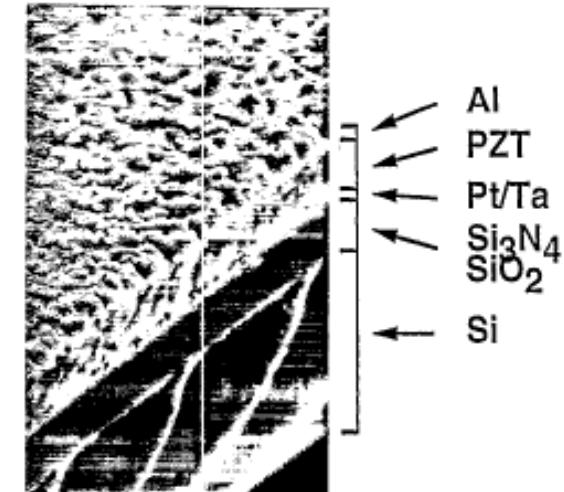
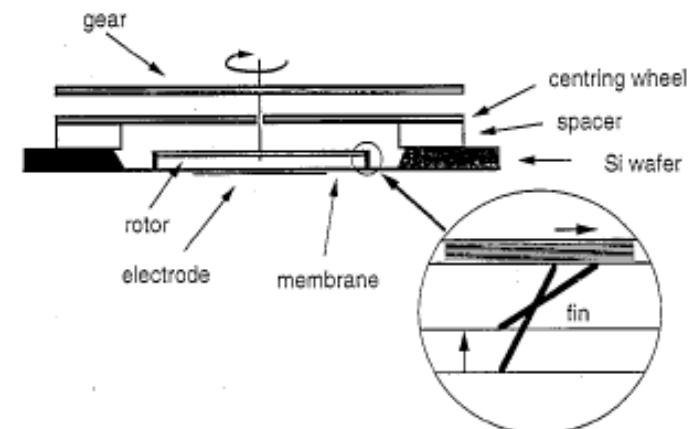
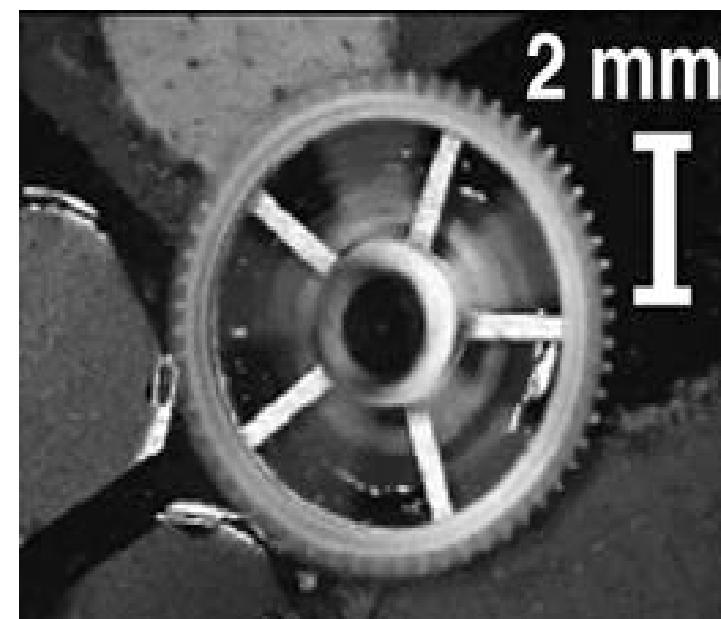
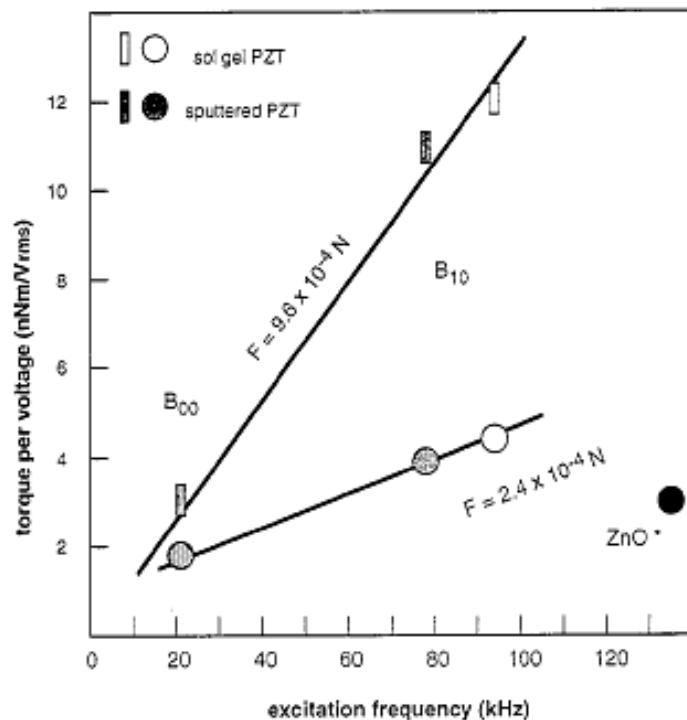


Fig. 2. Image of a cut through a membrane as obtained by scanning electron microscopy.

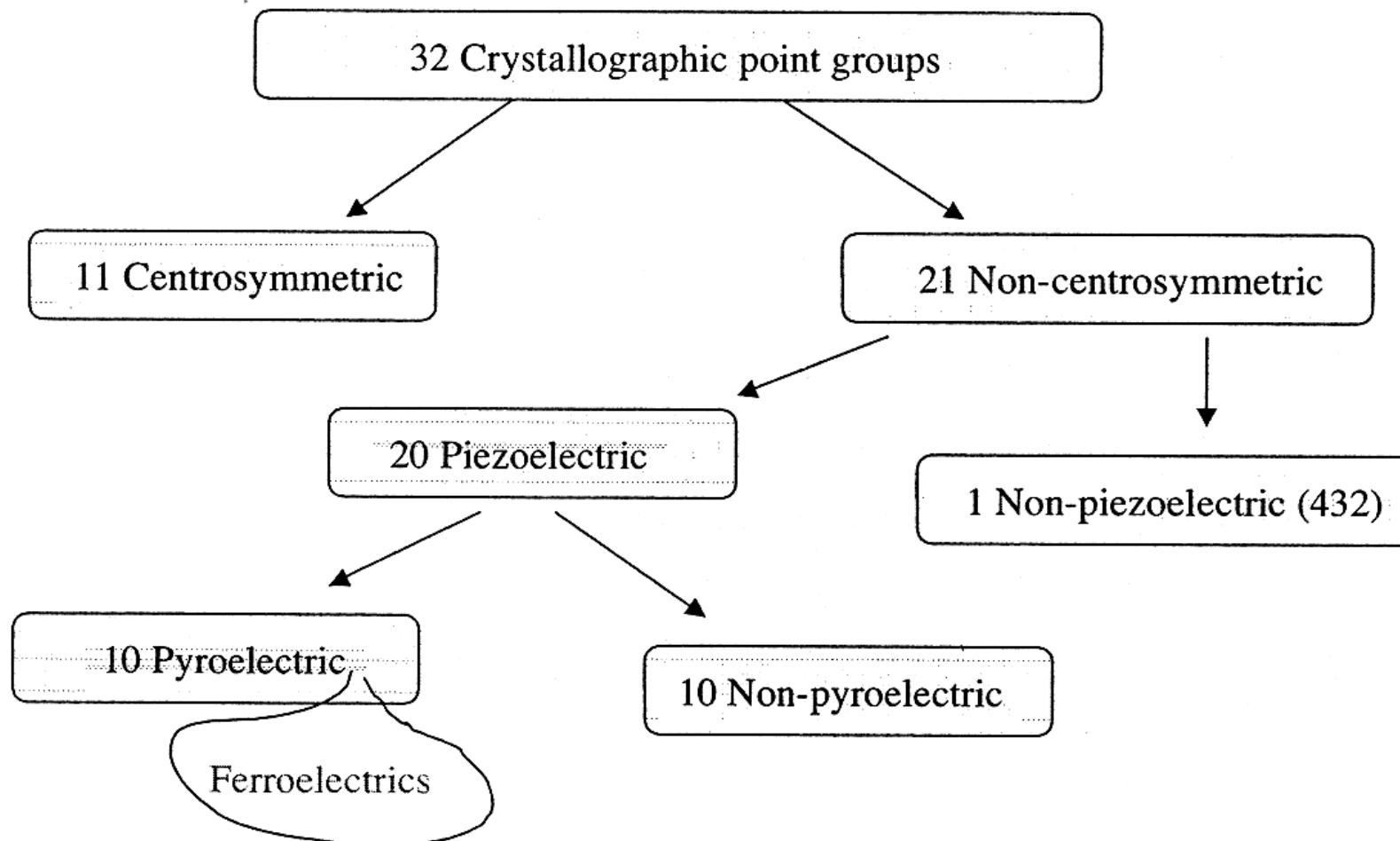


# *Motor performance vs. PZT routes*

Performance	Sol-gel no bias 94 kHz $B_{10}$	Sol-gel +2 V 94 kHz $B_{10}$	Sputter -2 V 21 kHz $B_{00}$	Sputter no bias 78 kHz $B_{10}$	Sputter -2 V 78 kHz $B_{10}$
Speed per a.c. voltage $F=2.4\times 10^{-4}$ N (rpm/V <sub>r.m.s.</sub> )	140	180	32	80	100
Torque per a.c. voltage $F=2.4\times 10^{-4}$ N (nN m/V <sub>r.m.s.</sub> )	4.3		1.8		3.9
Torque per a.c. voltage $F=9.8\times 10^{-4}$ N (nN m/V <sub>r.m.s.</sub> )	12		3.2		11



# Piezoelectricity: Coupling between mechanical and electrical energies



Direct piezoelectric effect:

Charge=d × force or

$$\mathbf{D} = \mathbf{d} \times \mathbf{X} \text{ (stress)}$$

Converse piezoelectric effect:

Dispacement=d × Voltage or

$$\mathbf{x} \text{ (strain)} = \mathbf{d} \times \mathbf{E}$$

or in general form

$$x_{ij} = \alpha_{ij}^{X,E} \Delta T + s_{ijkl}^{T,E} X_{kl} + d_{ijk}^{T,X} E_k$$

$$D_i = p_i^{X,E} \Delta T + d_{ikl}^{T,E} X_{kl} + \epsilon_{ij}^{T,X} E_i$$

(X, E) independent variables

Other choices of independent variables are possible:

(X,D), (x, E), and (x, D)

## Constitutive equations:

$$x_{ij} = s^E_{ijkl} X_{kl} + d^X_{ijk} E_k$$

$$D_i = d^E_{ikl} X_{kl} + \varepsilon^X_{ij} E_j$$

$$x = s^D X + g D$$

$$E = -g X + \beta^X D$$

$$X = c^E x - e E$$

$$D = e x + \varepsilon^x E$$

$$X = c^D x - h D$$

$$E = -h x + \beta^x D$$

$$d = \varepsilon^x g = e s^E$$

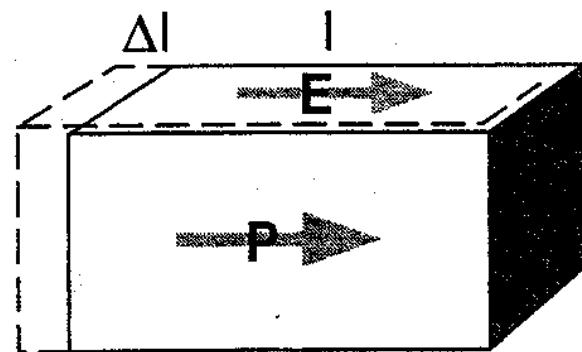
$$g = \beta^x d = h s^D$$

$$e = \varepsilon^x h = d c^E$$

$$h = \beta^s e = g c^D$$

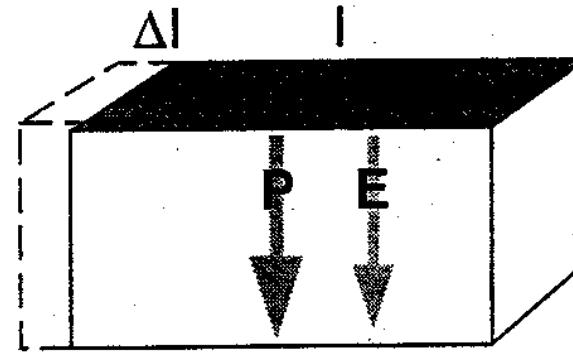
# *Major piezoelectric coefficients*

$$d_{33} = (\Delta l/l)_{||}/E$$



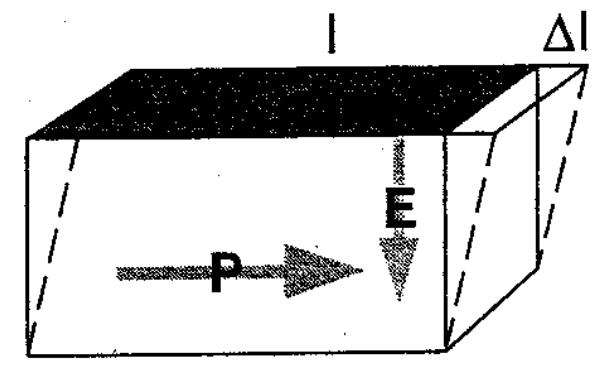
(a)

$$d_{13} = (\Delta l/l)_{\perp}/E$$



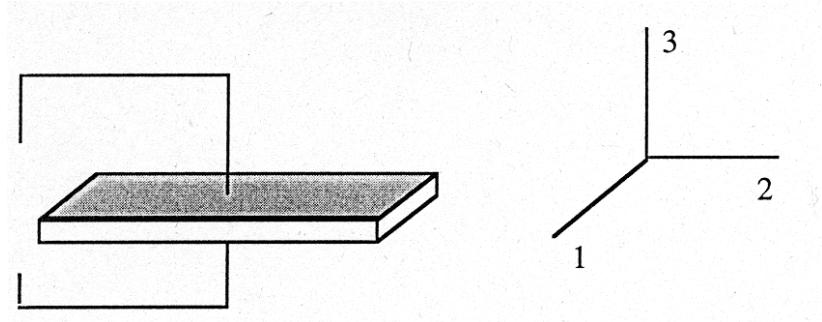
(b)

$$d_{15} = (\Delta l/l)_{\text{shear}}/E$$



(c)

The choice of equations depends on boundary conditions:



### 1D problem for a plate:

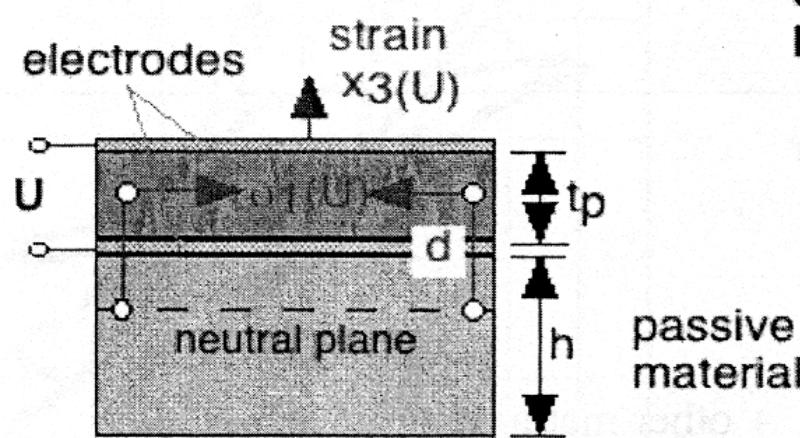
$E_1 = E_2 = 0$ ,  $E_3 \neq 0$  but spatially constant  $\rightarrow$  choose  $E$  as one independent variable.

$X_1 \neq 0$ ,  $X_2 = X_3 = X_4 = X_5 = X_6 = 0$ ,  $\rightarrow$  choose  $X$  as the other independent variable

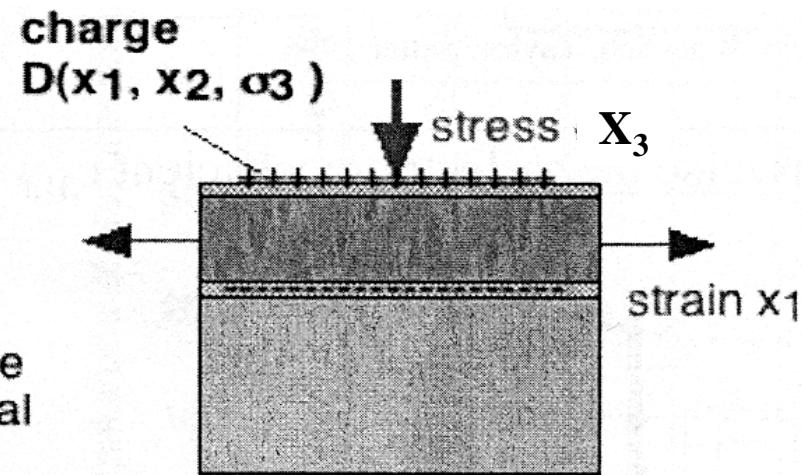
$$x_1 = s_{11}^E X_1 + d_{31} E_3$$

$$D_3 = d_{31} X_1 + \epsilon_{33}^X E_3$$

# *Clamping in thin films*



Converse effect



Direct effect

*For ideal clamping (thin film  
on thick rigid substrate) the equations:*

$$\frac{\sigma_1}{E_3} = e_{31,f} = \frac{d_{31}}{s_{11}^E + s_{12}^E} = e_{31} - \frac{c_{13}^E}{c_{33}^E} e_{33} \quad |e_{31,f}| > |e_{31}| \quad \text{Muralt 1996}$$

$$\frac{S_3}{E_3} = d_{33,f} = \frac{e_{33}}{c_{33}^E} = d_{33} - \frac{2s_{13}^E}{s_{11}^E + s_{12}^E} d_{31} \quad < d_{33} \quad \text{Royer & Kmetik 1992}$$

$$\frac{D_3}{E_3} = \varepsilon_{33,f} = \varepsilon_{33}^\sigma - \frac{2d_{31}^2}{s_{11}^E + s_{12}^E}$$

Actuator:

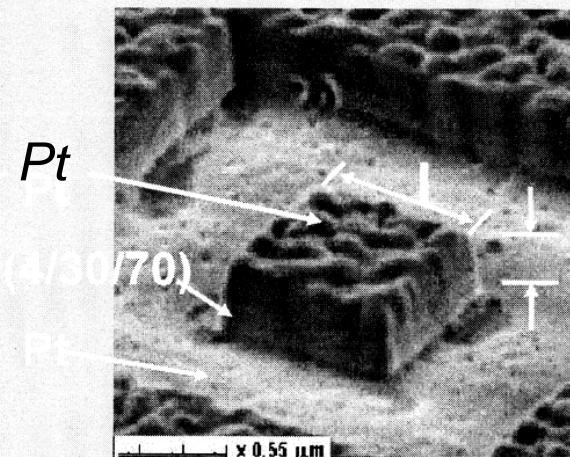
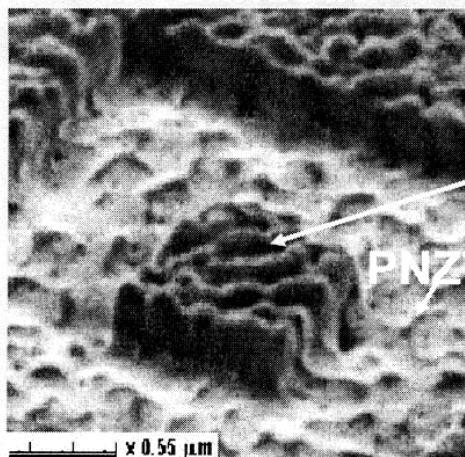
$$\sigma_{1,2} = -e_{31,f} * E_3$$

$$S_3 = d_{33,f} * E_3$$

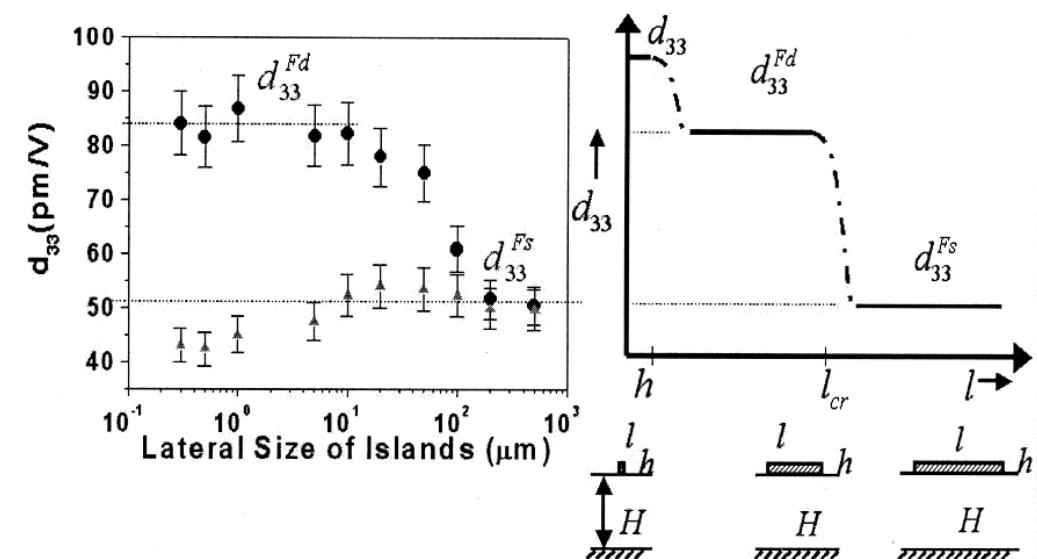
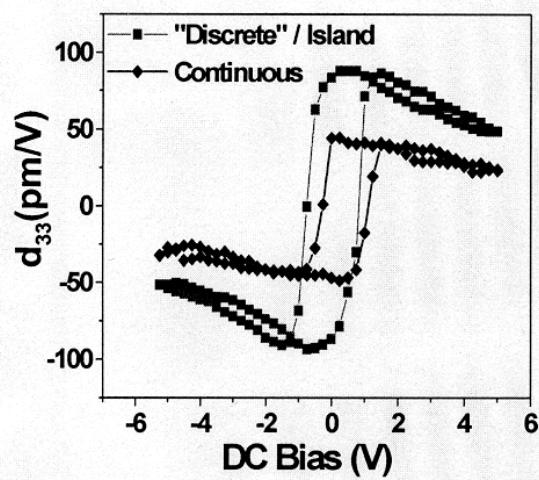
Sensor:

$$D = e_{31,f} * (S_1 + S_2) + d_{33,f} * \sigma_3$$

# De-clamping in nanostructures



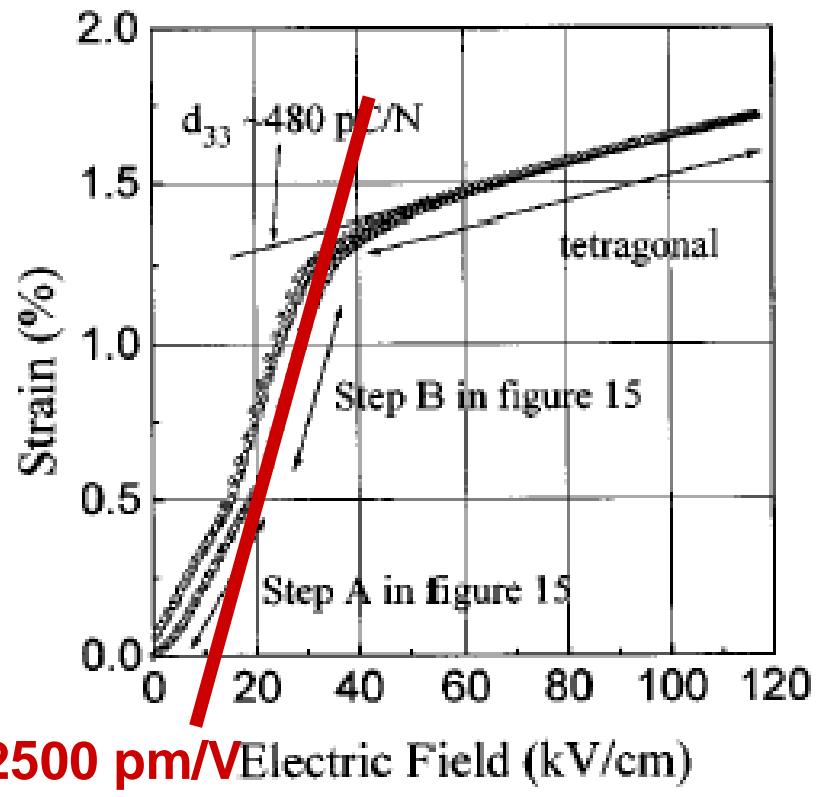
$$\Delta d = d_{33}^{Fd} - d_{33}^{Fs} = \frac{2S_{12}}{S_{11}^F + S_{12}^F} d_{13}$$



# *Advantages of piezoelectric actuation*

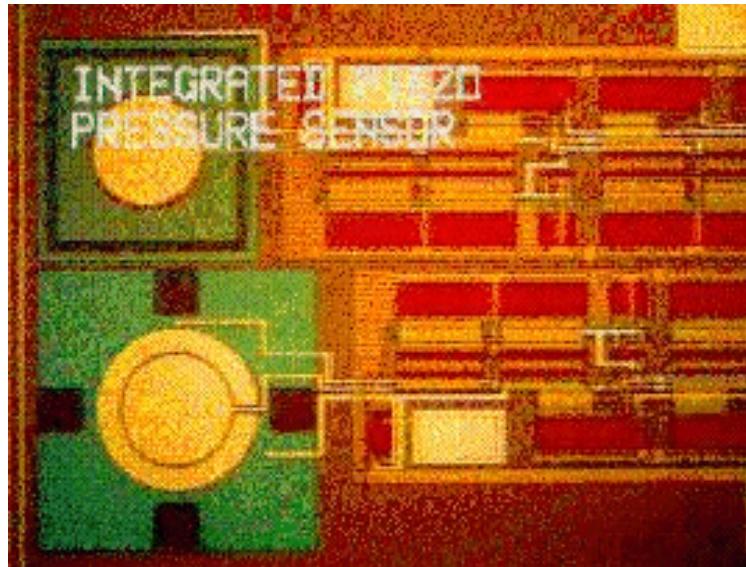
- High strain (up to 3-4 %)
- Unlimited resolution
- Rapid response ( $\mu$ s range)
- Good scalability
- High generative force
- Low power consumption
- Unlimited reliability
- *Problems with integration*

Giant strain in (001)-oriented  
 $\text{Pb}(\text{Zn},\text{Nb})\text{O}_3$ -8%PT

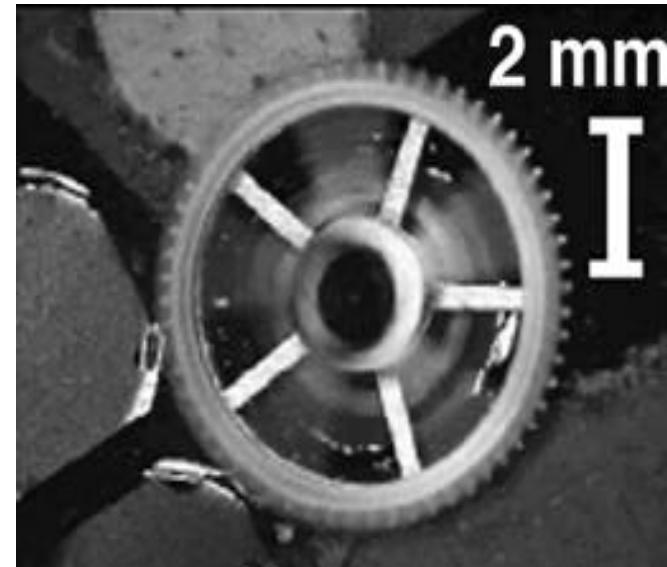


Park & Shroud, JAP (1997)

# *Examples of piezoelectric MEMS*

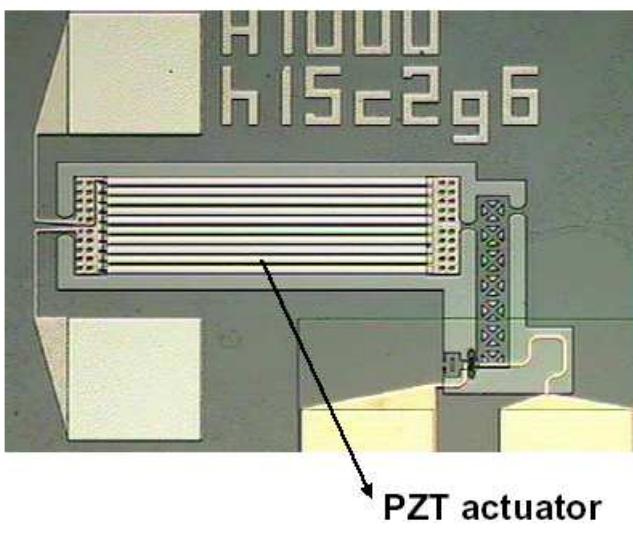
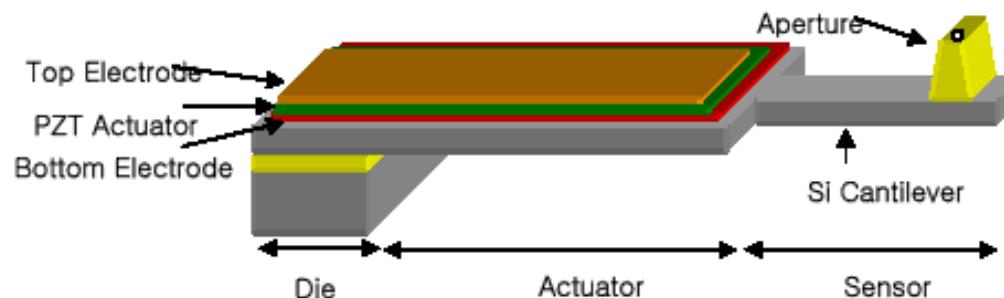


**Acoustic pressure sensor  
(U. of Minnesota)**

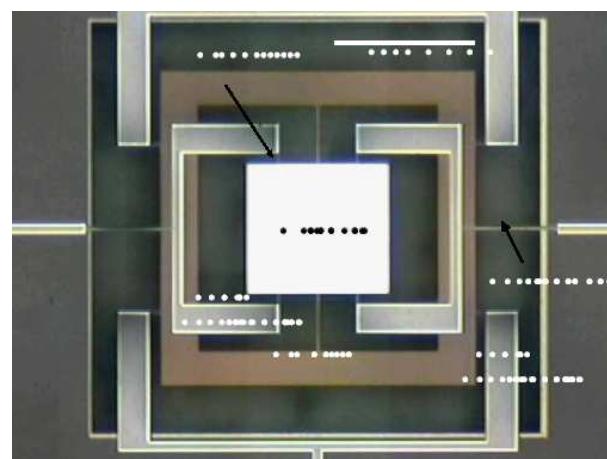


**Piezoelectric micromotor  
(EPFL, Switzerland)**

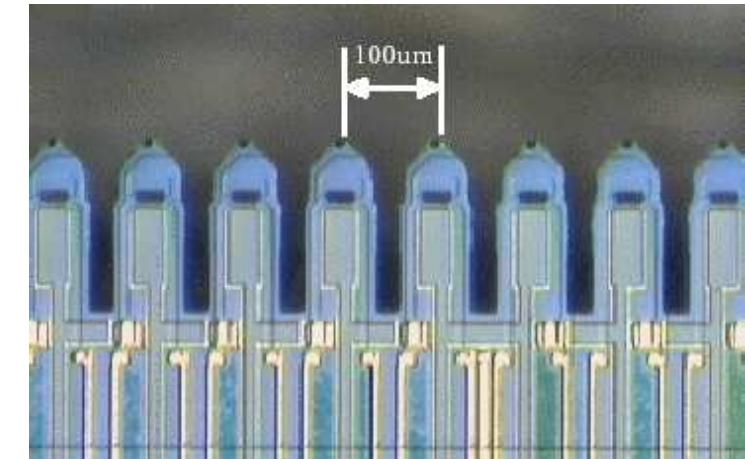
## *Examples of piezoelectric MEMS (cont.)*



Piezo microswitch



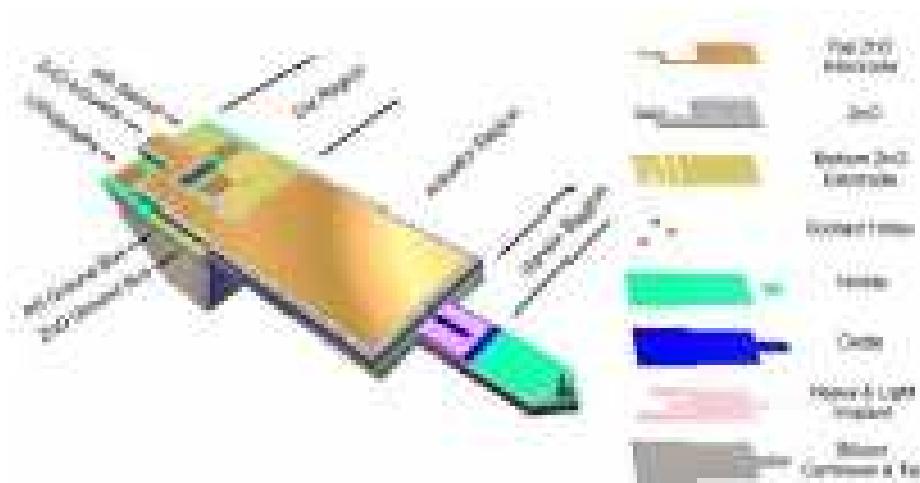
Tilting micromirror



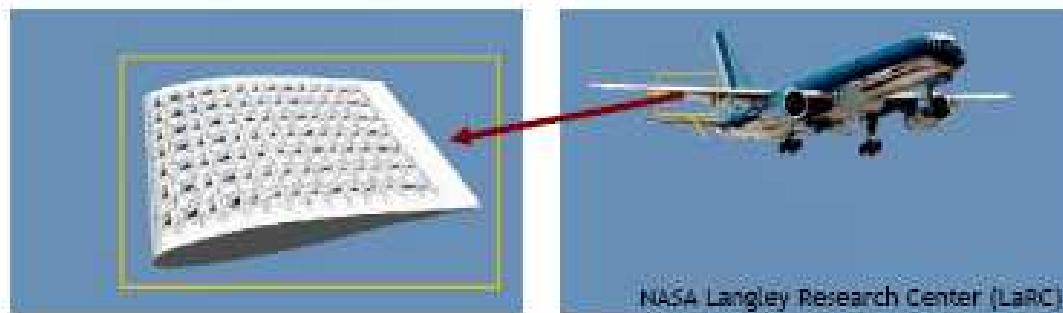
Cantilever array

Courtesy of Inostek, Inc. (S. Korea)

## *Examples of piezoelectric MEMS (cont.)*



Sensor/actuator for AFM



Smart wing for aircraft

# *Ferroelectric MEMS*

- General features

- Small strains (1% max, 0.2%-0.5% more realistic)

Typically designs rely on “d<sub>31</sub>” mode to increase strain:

diaphragms, bridges & cantilevers

- High sensitivity to stress (stress-compensating layers)
- Efficient electromechanical coupling (95% in some cases)
- Weak hysteresis (in some compositions)
- Excellent high-frequency potential

- Complicated processing

-two or more metal cations in compositions of interest (PZT, d<sub>33</sub>= 250-600 pm/V))

-Typically contain Pb (need of lead-free materials)

-Oxidizing process atmospheres

# *Non-ferroelectric MEMS*

- **Non-ferroelectric piezoelectric MEMS**

- ZnO  $d_{33} \sim 4 \div 10 \text{ pm/V}$  (2 times larger than quartz)

- AlN  $d_{33} \sim 2 \div 6 \text{ pC/N}$  (compatible with  $\text{Si}_3\text{N}_4$ )

- Modulation of  $n$  in ZnO-clad optical fiber

- Much easier processing**

- Relatively low-temperature growth

- **Comparatively small electromechanical response**

- **But must be deposited with specific crystallographic texture**

- i.e., piezoelectric axis normal to substrate

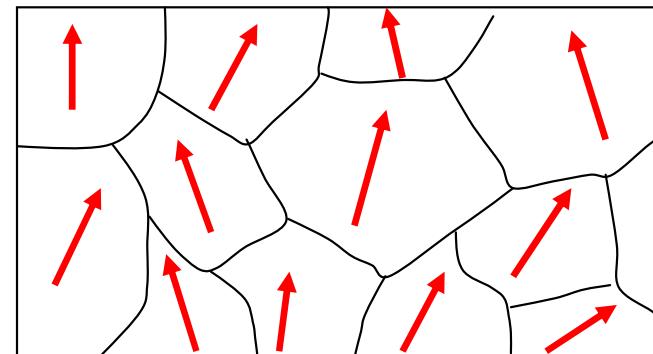
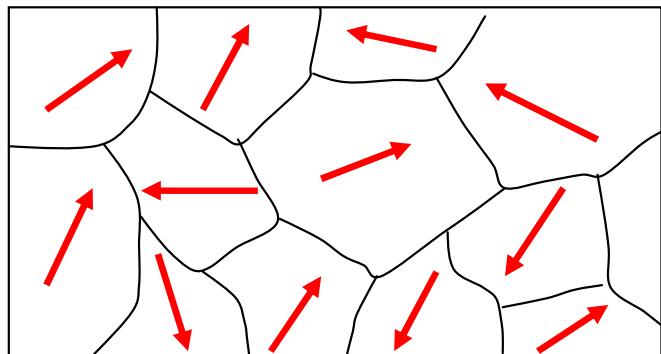
- $c$ -oriented film for both ZnO and AlN

- **Advantages of FE MEMS**

- Potentially much larger piezoelectric response

- $d_{33}$  values up to **2500 pm/V**

- as deposited, FE films are NOT piezoelectric and poling can produce piezoelectricity for any texture



- No crystallographic texture required
  - In some cases, random orientation gives best results

# *MEMS materials*

- Same materials as in bulk ceramics

$\text{Pb}(\text{Zr}, \text{Ti})\text{O}_3$  (PZT) solid solution – ferroelectric

$\text{Pb}(\text{Mn}_{1/3}\text{Nb}_{2/3})\text{O}_3$  – electrostrictors

Relaxor- $\text{PbTiO}_3$  solid solutions and new MPB systems

such as  $\text{BiScO}_3$ - $\text{PbTiO}_3$  (High  $T_c$ )

- Most research still on  $\text{Pb}(\text{Zr}_x\text{Ti}_{1-x})\text{O}_3$  solid solution

- $x \sim 0.52$  (morphotropic composition) ideal

- $d_{33}$  coefficients approach now above 200 pm/V (unclamping)

- Non-lead containing compositions

- $\text{BaTiO}_3$ ,  $(\text{Sr}_x\text{Ba}_{1-x})\text{NbO}_3$ , or  $(\text{K},\text{Na})\text{NbO}_3$ -based

- difficult to process

- Important observations

- bulk properties difficult to achieve in thin films

- scaling effects

- extrinsic vs. intrinsic contribution

# *Related issues*

## - Intrinsic contribution to piezoelectricity

$$d_{ijk} = 2Q_{ijkl}\epsilon_0\epsilon_{kl}P_l$$

large permittivity

large remanent polarization

## - Extrinsic contribution to piezoelectricity

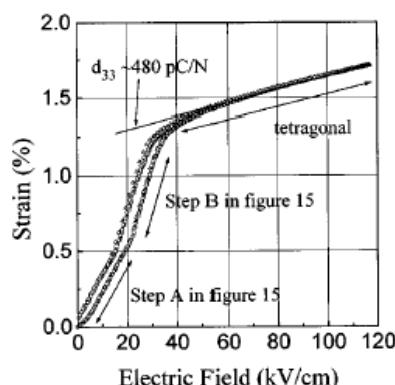
existence of non-180° domains

large spontaneous strain

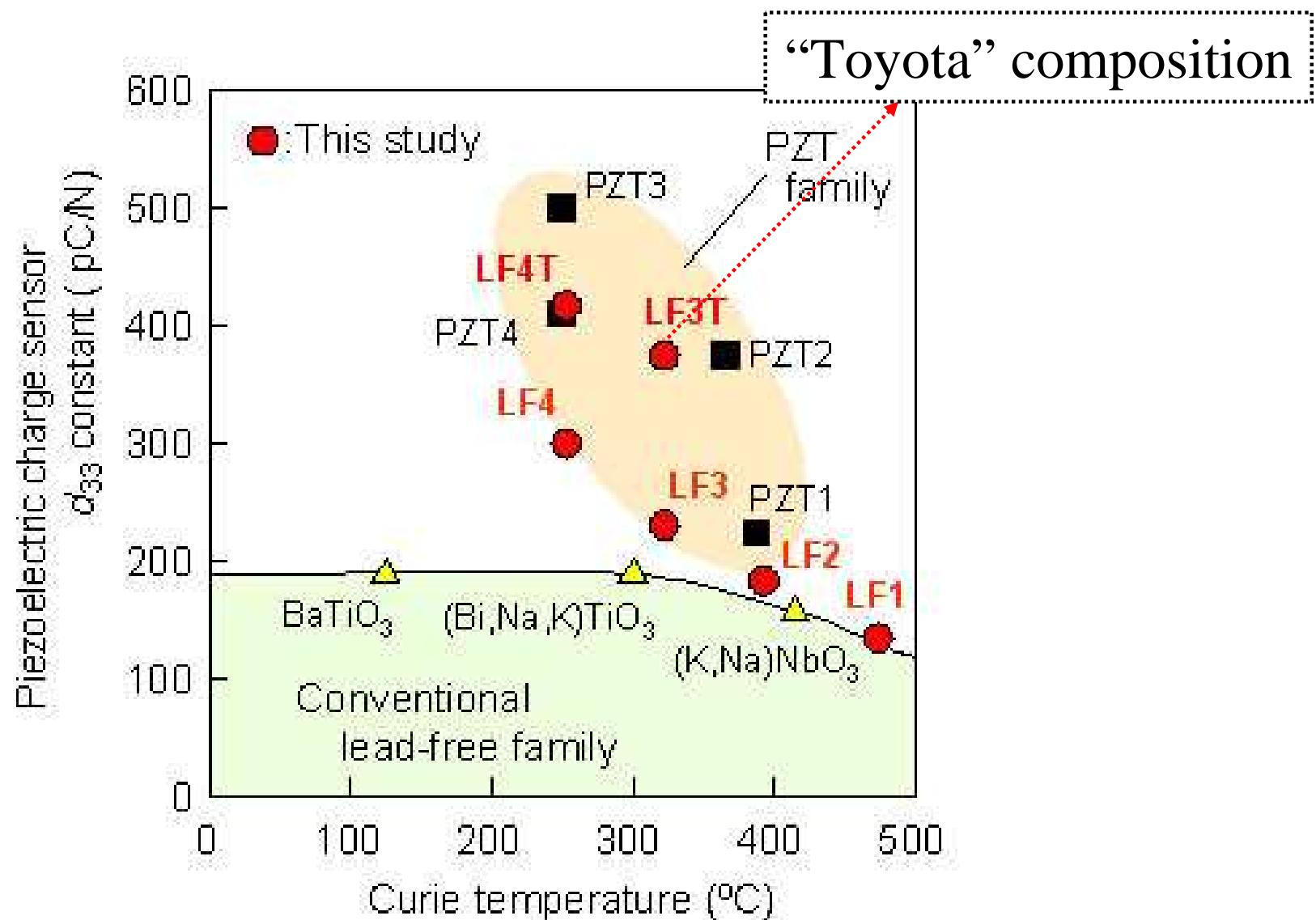
multiple available domain variants

morphotropic compositions desirable

composition-induced symmetry transition

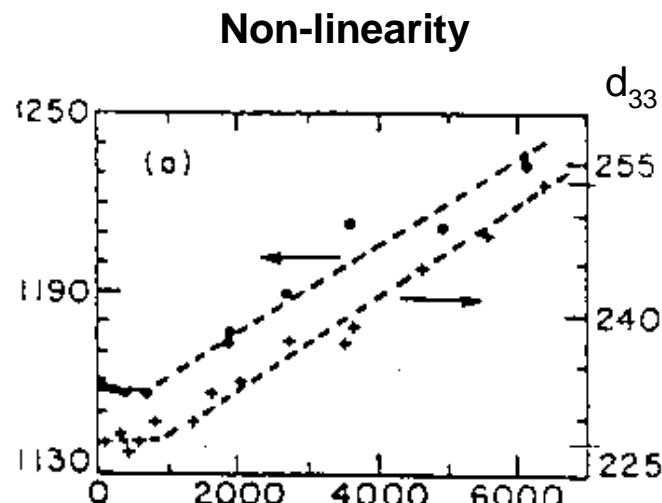


# PZT vs. lead-free materials

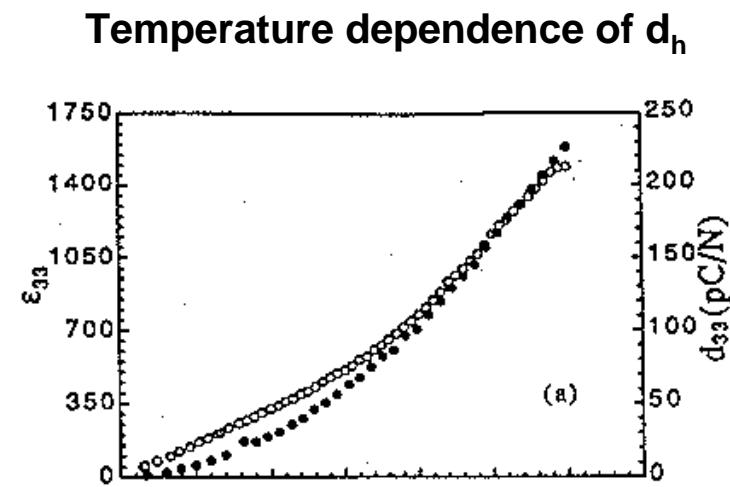


Y. Saito, et al. "Lead-free piezoceramics", *Nature*, **432**, Nov.4, 84-87 (2004).

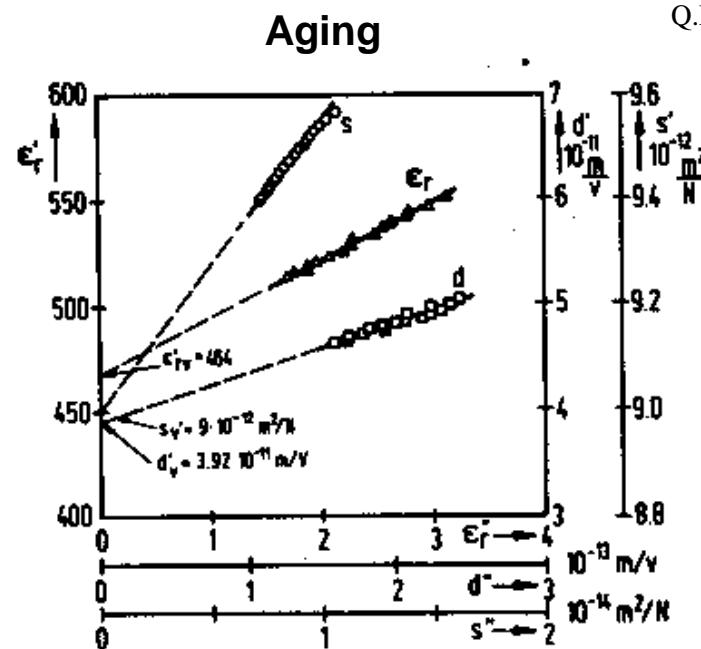
# *Experimental evidence of domain wall contributions*



Q.M.Zhang et al, JAP **64**, 6445 (1988)

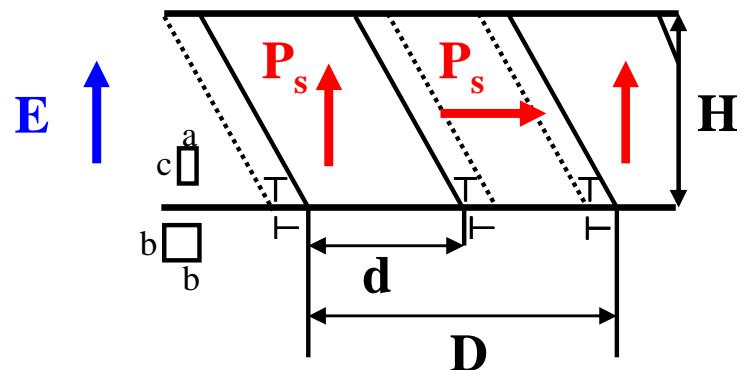


Q.M.Zhang et al, JAP **75**, 454 (1993)



R.Herbiet et al, Ferroelectrics **76**, 319 (1987)

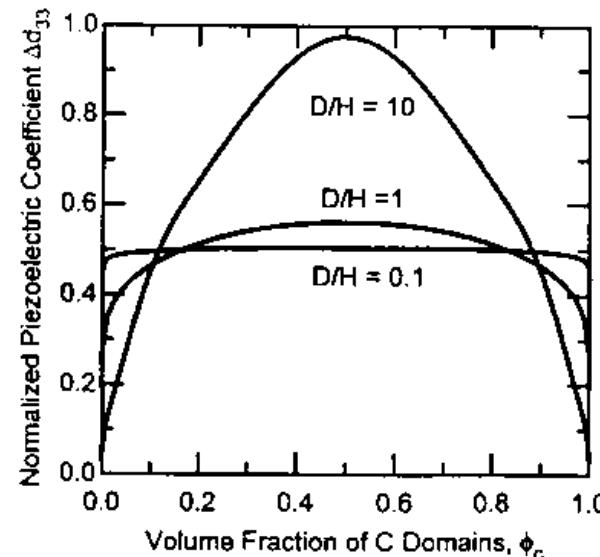
# Theoretical predictions for extrinsic effect



$$d_{ext} = 2\sqrt{2}\left(1 - \frac{\nu}{1-\nu}\right)(Q_{11} - Q_{12}) \frac{P_s^3}{kD}$$

for  $D/H \ll 1$

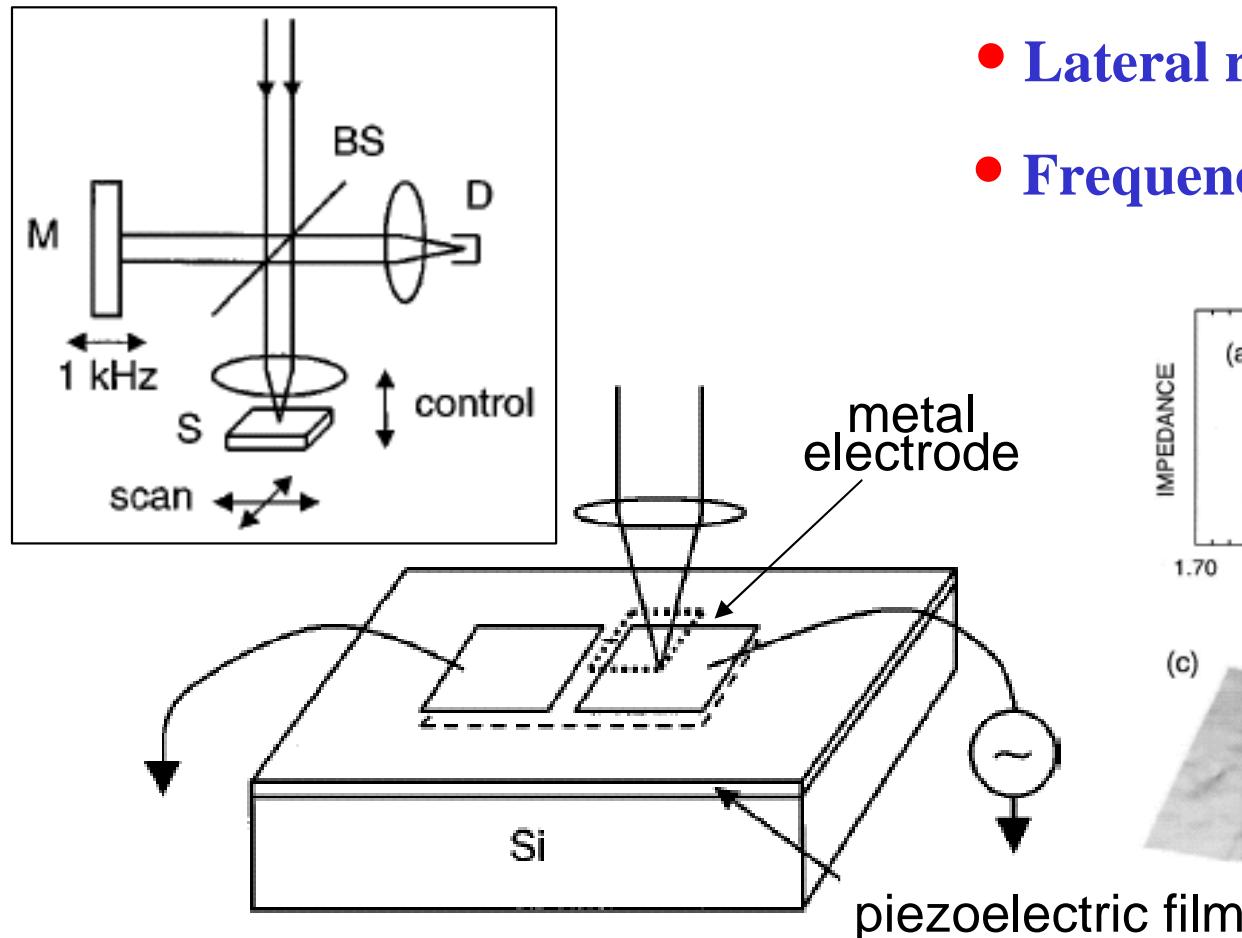
N. Pertsev & A. Emelyanov, APL (1997)



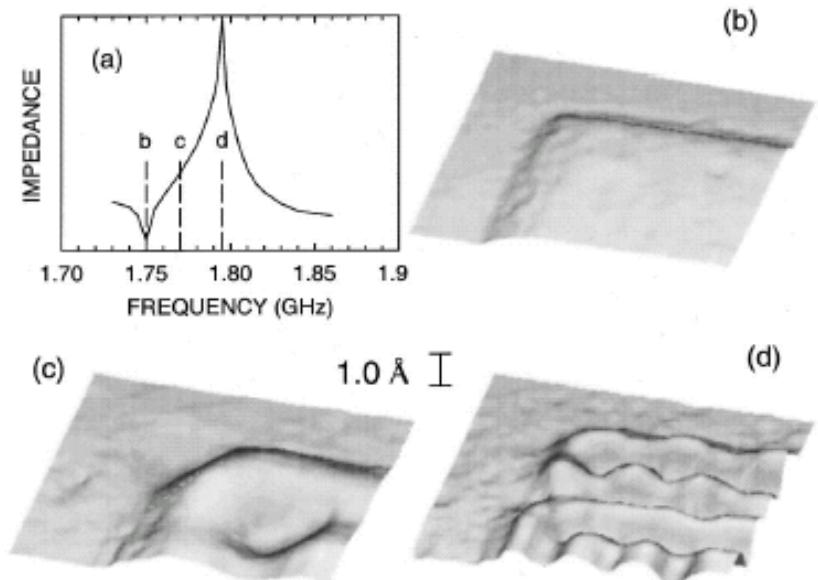
$d_{ext} = 70 \text{ pm/V}$      $\text{PbTiO}_3$   
 $d_{ext} = 100 \text{ pm/V}$      $\text{PZT}$

# *Macroscopic characterization: laser interferometry*

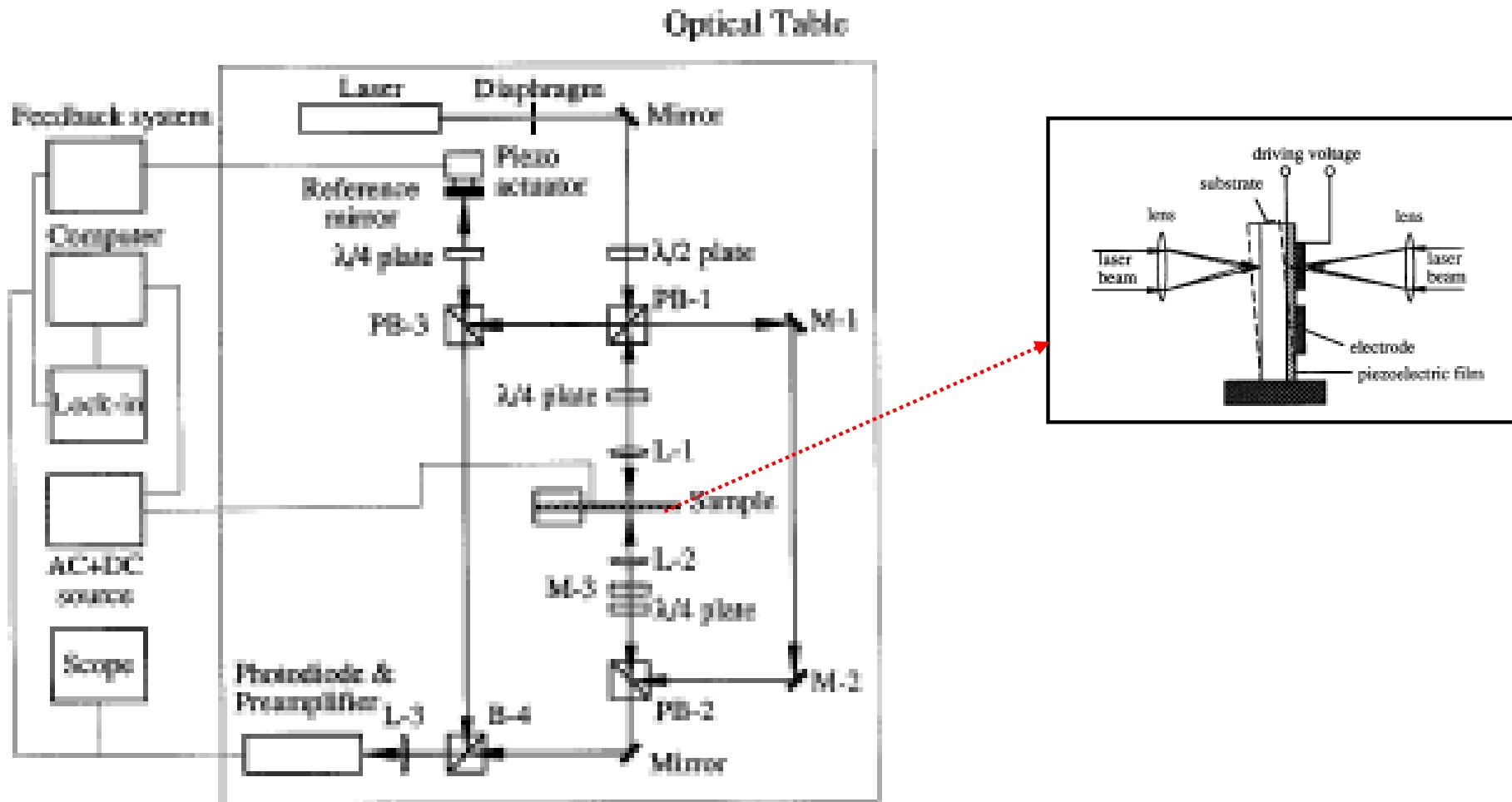
Michelson interferometer



- Vertical resolution –  $0.003 \text{ \AA}$
- Lateral resolution -  $0.6 \mu\text{m}$
- Frequency range  $< 6 \text{ GHz}$

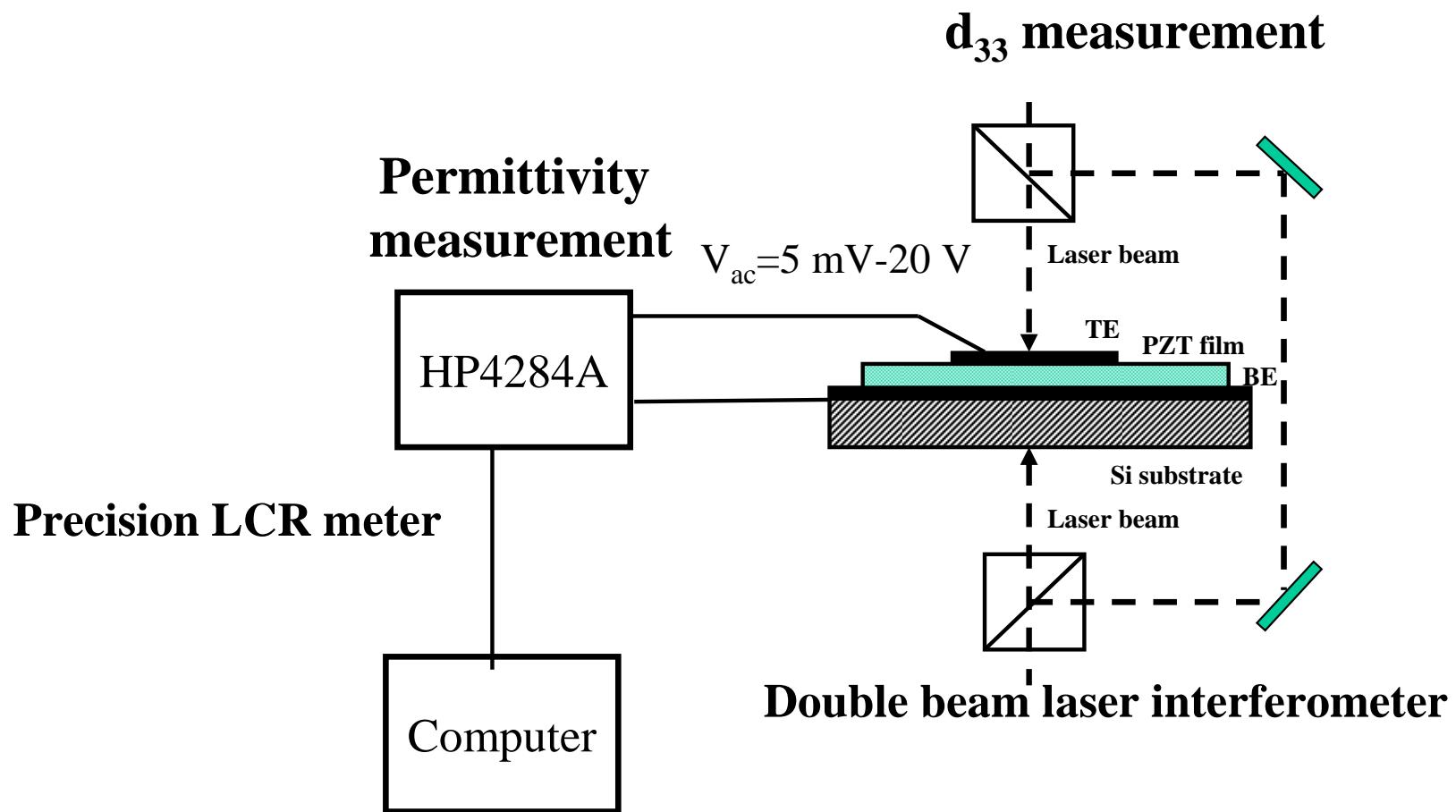


# Piezoelectric measurements in thin films

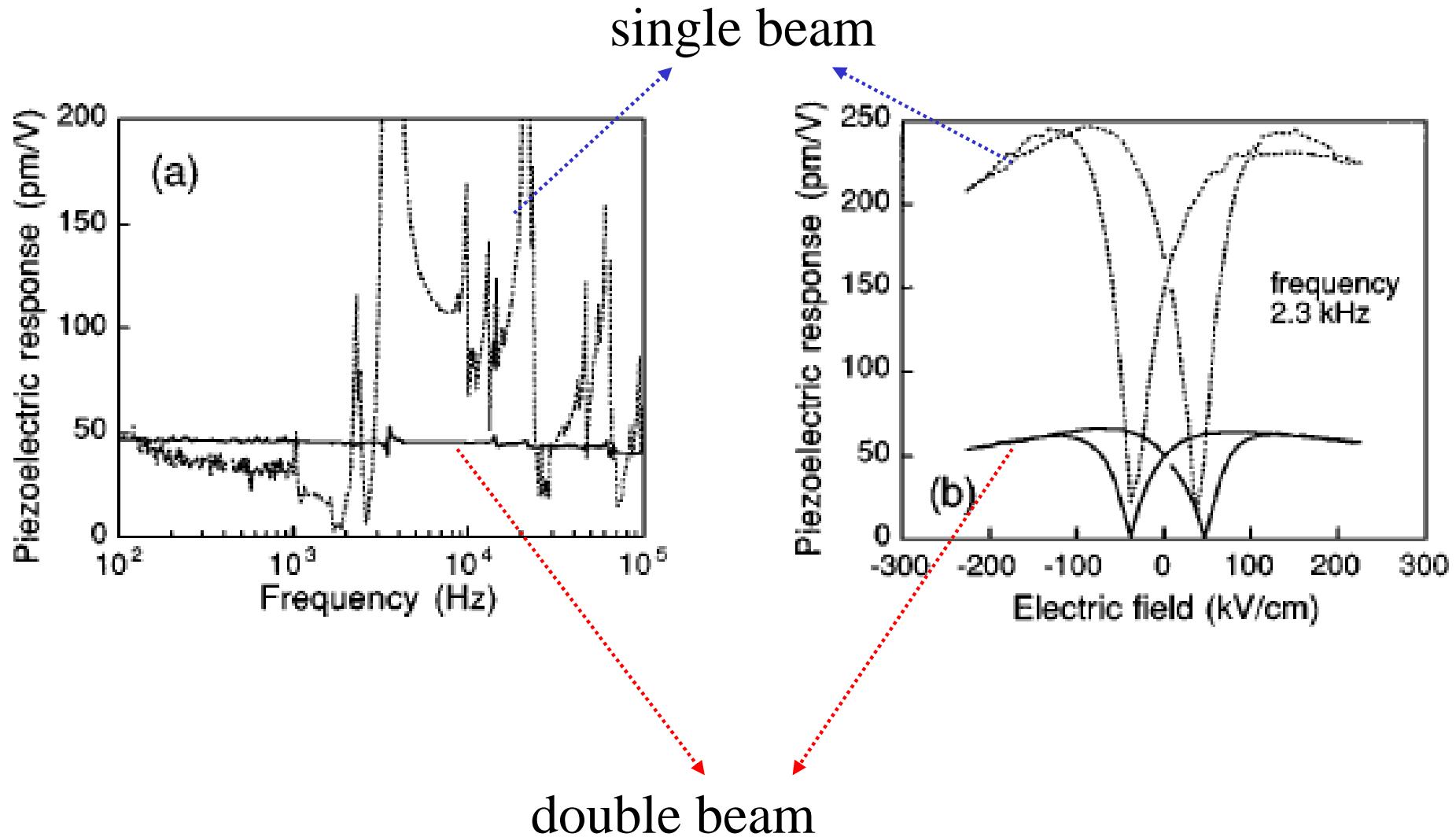


Kholkin et al., Rev. Sci. Instr. 1996

# *Simultaneous $d_{33}$ and $\epsilon$ measurements*



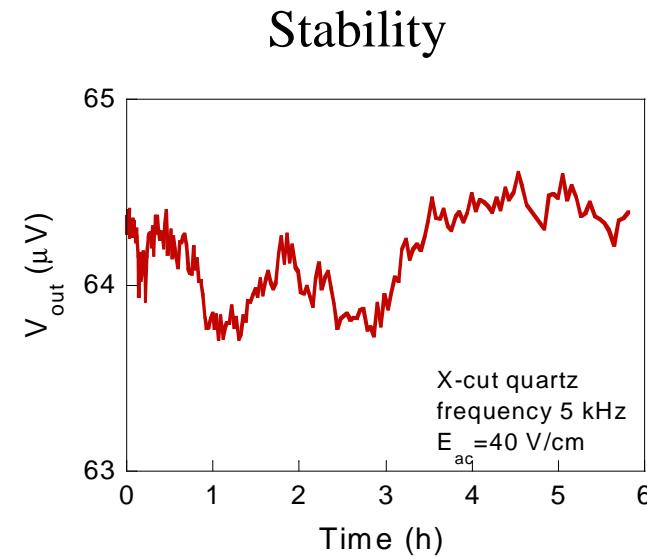
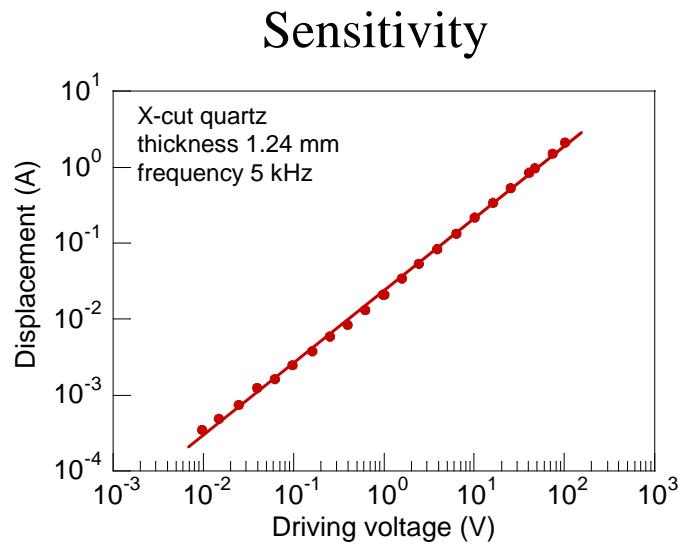
# *Substrate bending elimination*



# *Double-beam laser interferometer*

## Key issues

- Single-mode laser with improved stability
- Short optical path
- Reduced environmental noise
- Active stabilization of the optical path length difference



# Bending in single-beam interferometer

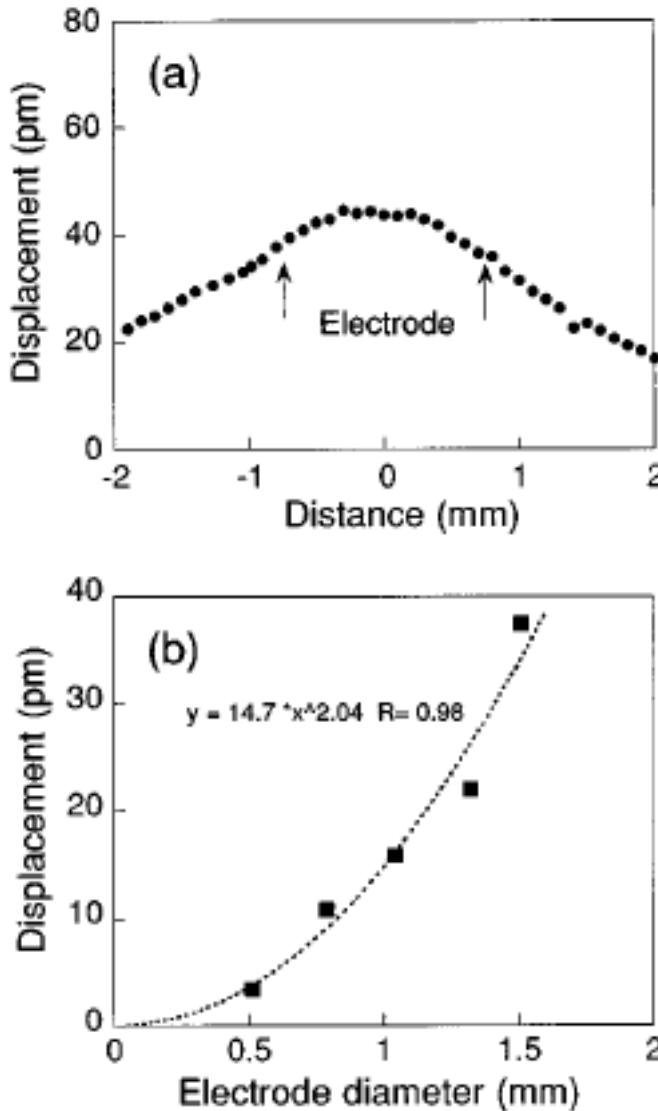


FIG. 2. (a) Deformation of the Si substrate induced by piezoelectric effect in the PZT film. Arrows indicate the edges of the top electrode. (b) Displacement of the top electrode as a function of the electrode diameter.

# *Typical values for PZT thin films*

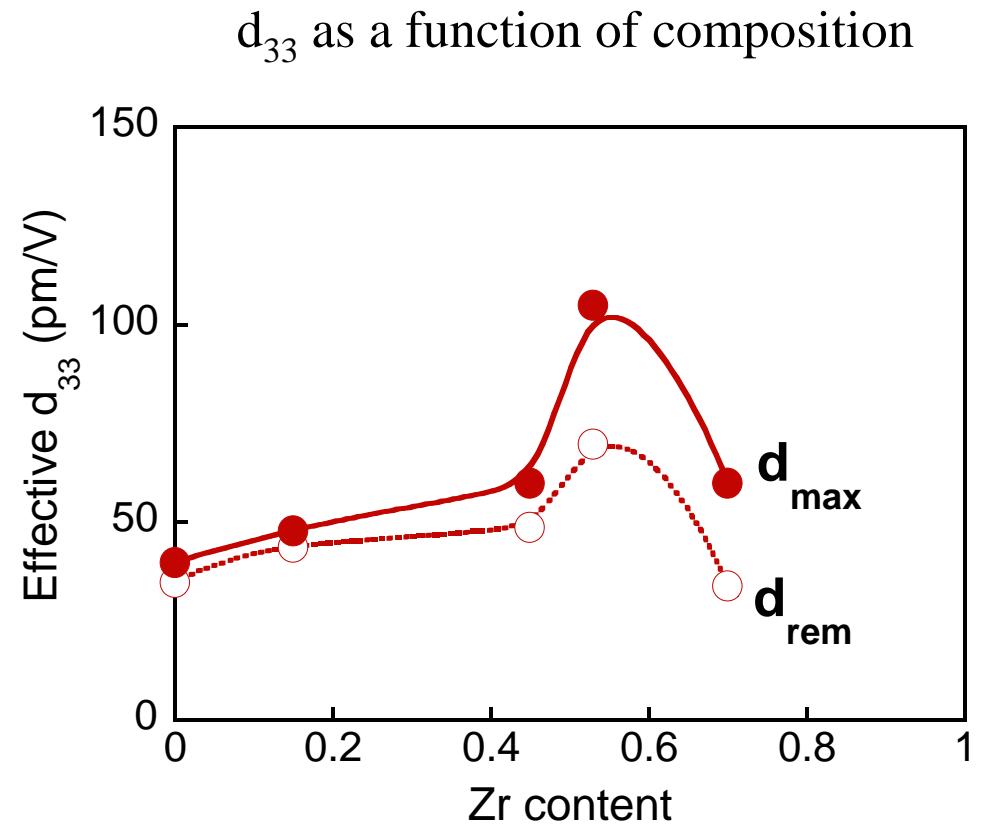
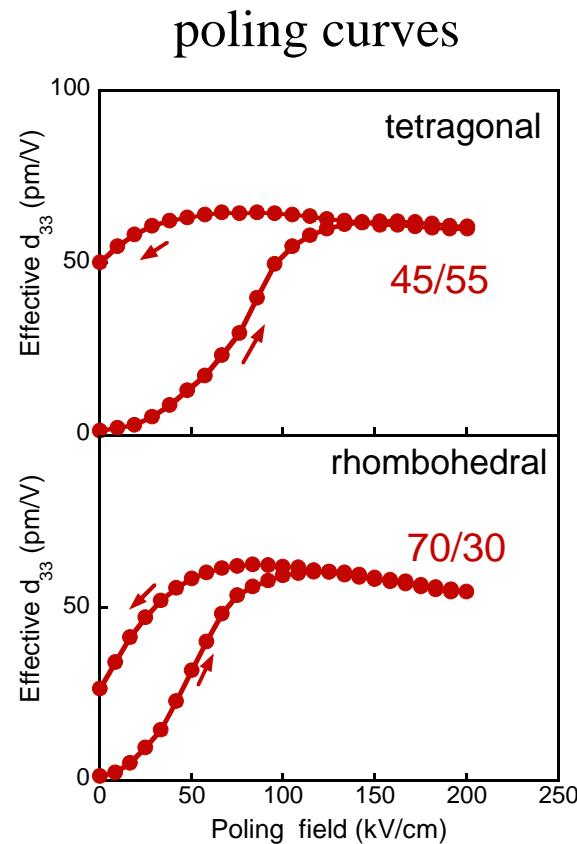
<b>Reference</b>	<b>Effective <math>d_{33}^*</math> (pm/V)</b>	<b>Meas. technique</b>
K.F.Etzold et al. (1990)	8	converse effect
Z.Surowiak et al. (1992)	70-120	direct effect
K.Lefki and G.J.M. Dormans (1994)	400	direct effect
J.-F.Li et al. (1994)	140	converse effect
H.D.Chen et al. (1995)	196	converse effect
A.L.Kholkin et al. (1996)	40-110	converse effect
F.Xu et al. (1997)	150	direct effect

\*PZT films of MPB composition and thickness  $\leq 1 \mu\text{m}$

For undoped PZT ceramics  $\mathbf{d_{33}\approx 230 \text{ pm/V}}$

# *Effect of the composition (sol-gel)*

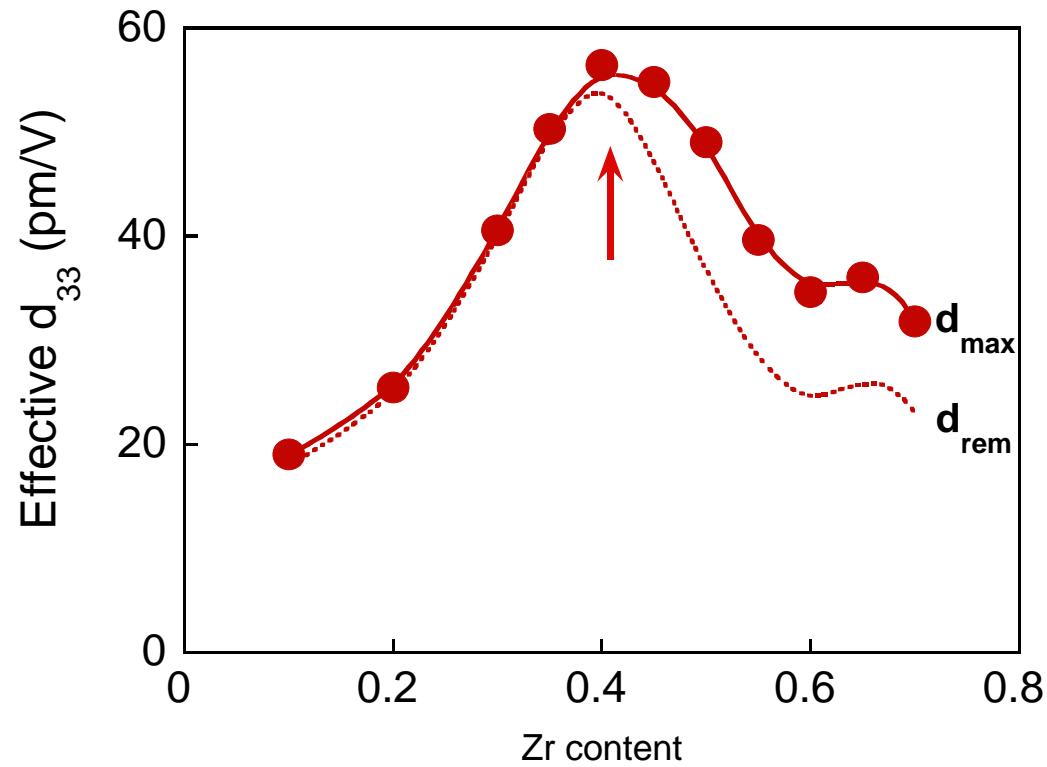
sol-gel PZT, thick. 0.3  $\mu\text{m}$ , (111) orientation



- significant depoling for MPB and rhombohedral compositions
- $d_{33}$  for MPB composition is about half of the value for undoped PZT ceramics

# *Effect of the composition (sputtering)*

sputtered PZT, thick. 0.3  $\mu\text{m}$ , (111) orientation

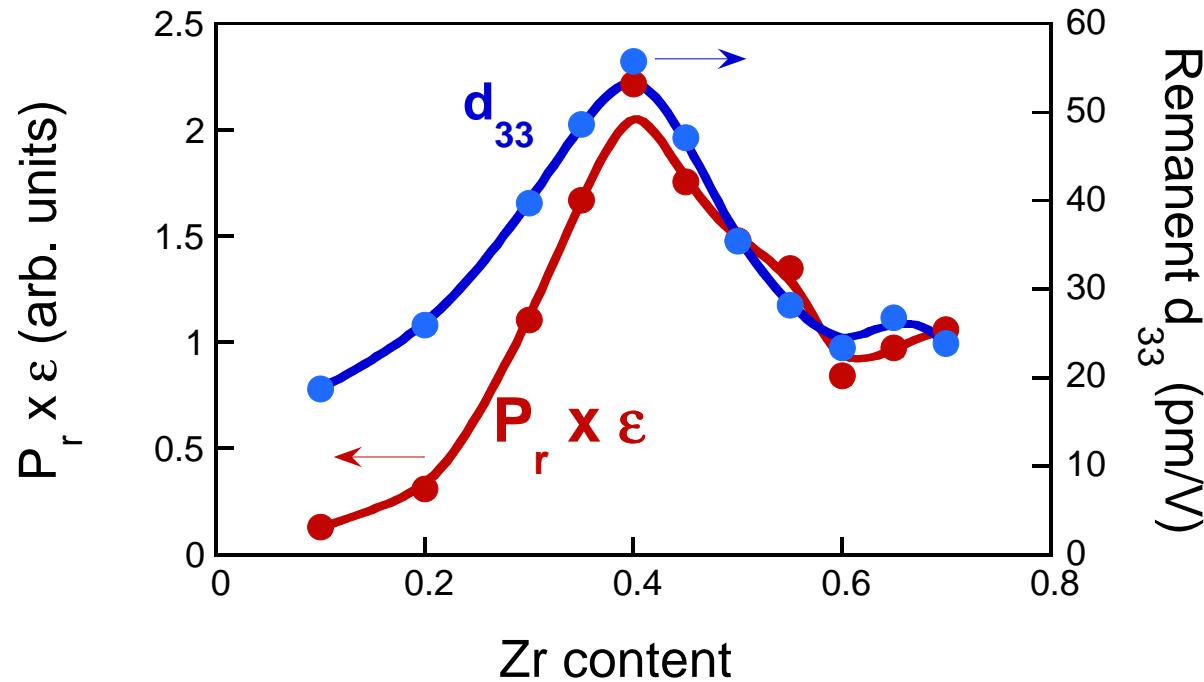


- In rhombohedral films non-180° domains may contribute to electromechanical strain under high ac-field

# *Intrinsic vs. Extrinsic effect*

Linearized electrostriction equation

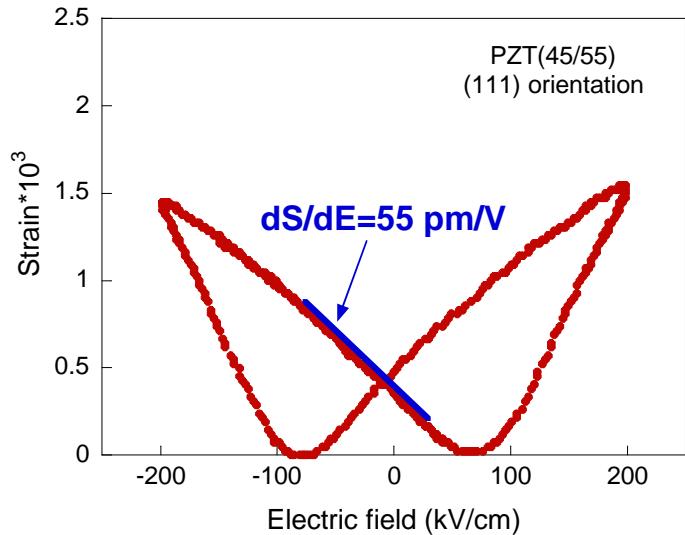
$$d_{33} = 2Q_{11}\epsilon_0 \epsilon P_r$$



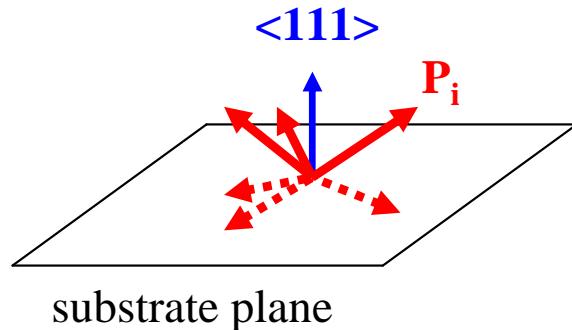
- $Q \approx 0.015-0.02 \text{ m}^4/\text{C}^2$  for rhombohedral compositions
- $d_{33}$  is relatively high in tetragonal films due to non-switchable self-polarization

# *Effect of orientation*

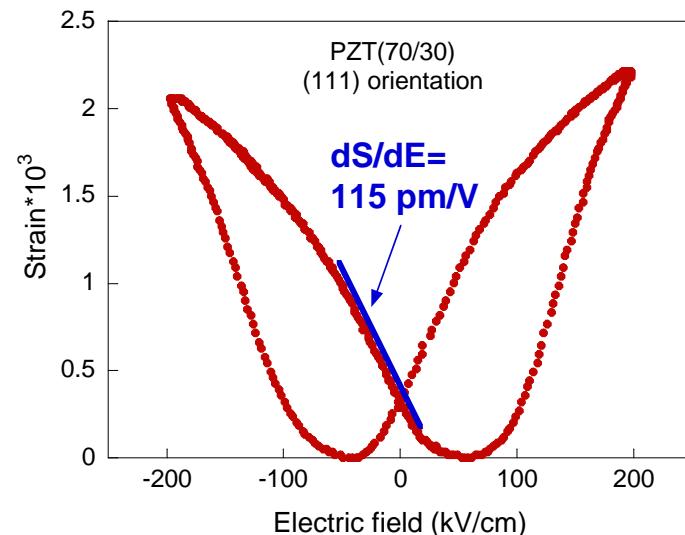
tetragonal (45/55)



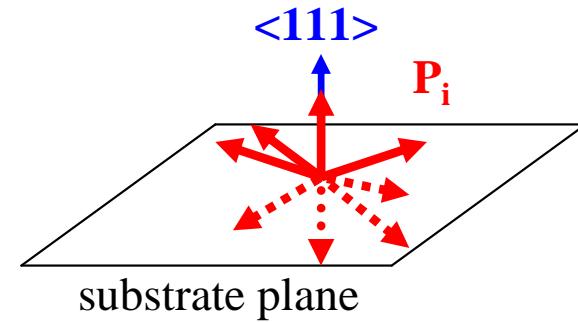
$$dS/dE \approx d_{33}$$



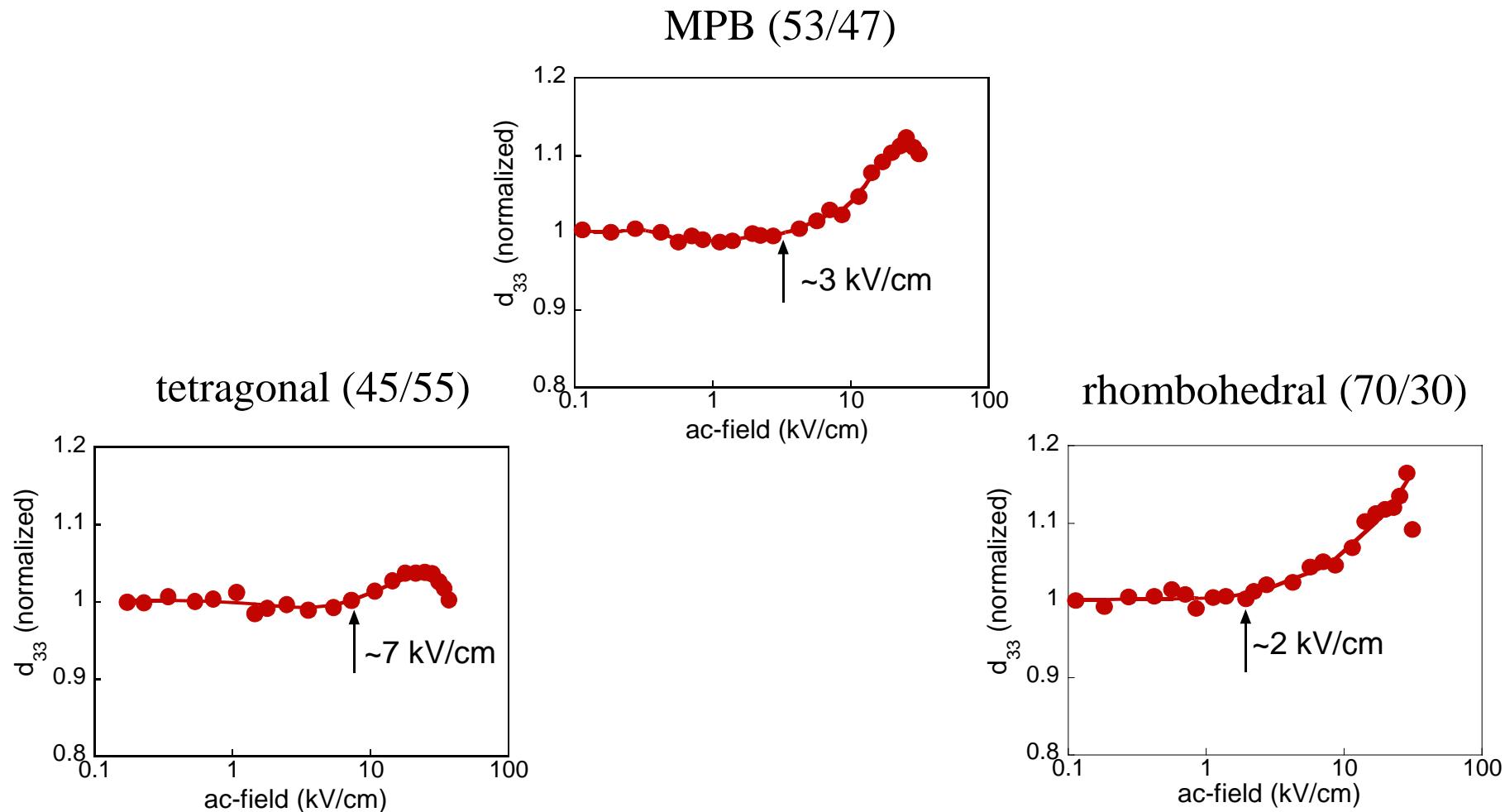
rhombohedral (70/30)



$$dS/dE \gg d_{33}$$

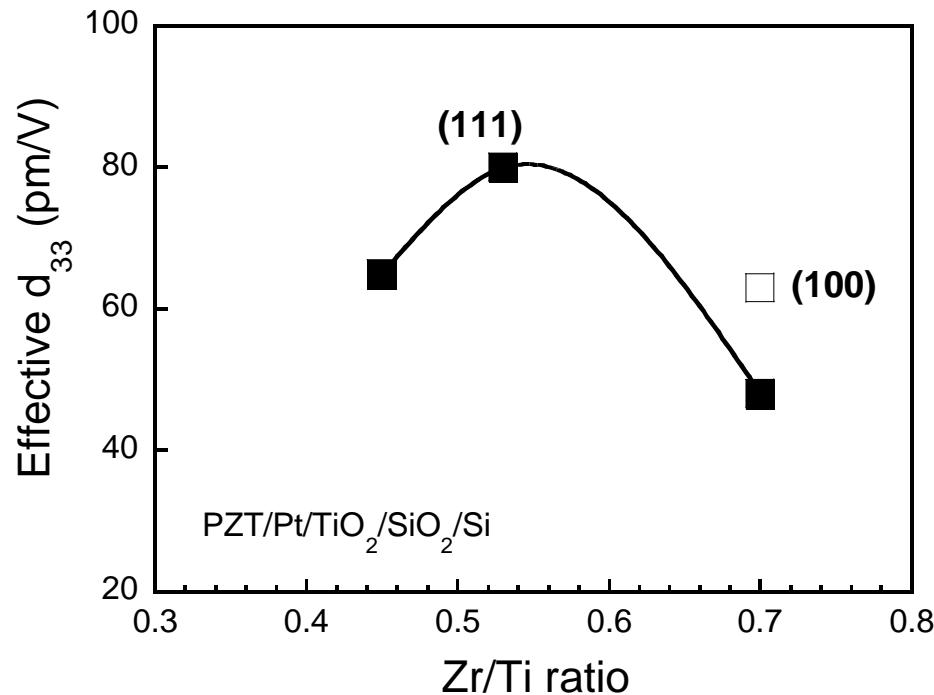


# Piezoelectric non-linearity



- $90^\circ$  domains are almost inactive in tetragonal films
- domain wall contribution is evident in MPB and rhombohedral films

# *Effect of orientation*



# Comparison with calculations

## Tetragonal composition

$$d_{31} = d_{32} = 2\epsilon_o \epsilon_{33} Q_{12} P_3$$

$$d_{33} = 2\epsilon_o \epsilon_{33} Q_{11} P_3$$

$$d_{15} = d_{24} = 2\epsilon_o \epsilon_{11} Q_{44} P_3$$

## Rhombohedral composition

$$d_{11} = d_{22} = d_{33} = 2\epsilon_o (\epsilon_{33} Q_{11} + 2Q_{12} \epsilon_{12}) P_3$$

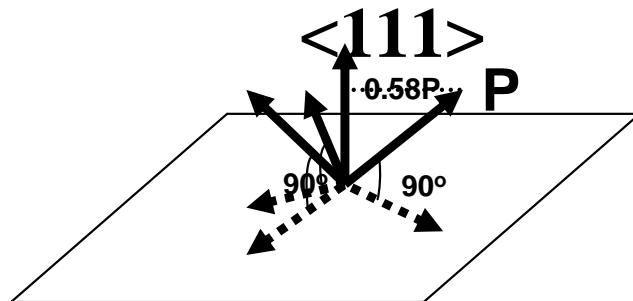
$$d_{12} = d_{13} = d_{23} = 2\epsilon_o [\epsilon_{11} Q_{12} + \epsilon_{12} (Q_{11} + Q_{12})] P_3$$

$$d_{14} = d_{25} = d_{36} = 2\epsilon_o \epsilon_{12} Q_{44} P_3$$

$$d_{15} = d_{16} = d_{24} = d_{26} = d_{34} = d_{35} = \epsilon_o (\epsilon_{11} + \epsilon_{12}) Q_{44} P_3$$

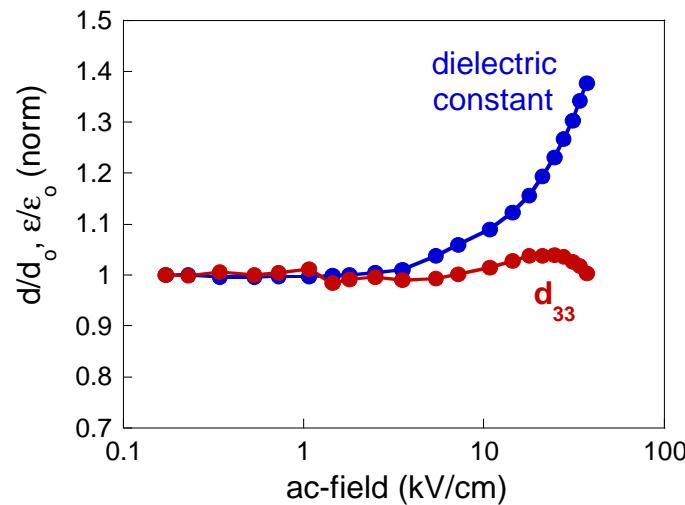
\*All thermodynamical coefficients are taken from Haun et al., Ferroelectrics **99**, 63 (1989)

Symmetry	Composition	Orientation	$d_{33}$ (free)	$d_{33}$ (clamped)	$\alpha$	$d_{33}$ (exp)
tetragonal	45/55	[111]	161	<b>58</b>	0.64	<b>62</b>
rhombohedral	70/30	[111]	63	<b>47</b>	0.25	<b>60</b>
		[001]	107	71	0.34	78
		[011]	88	42	0.52	—



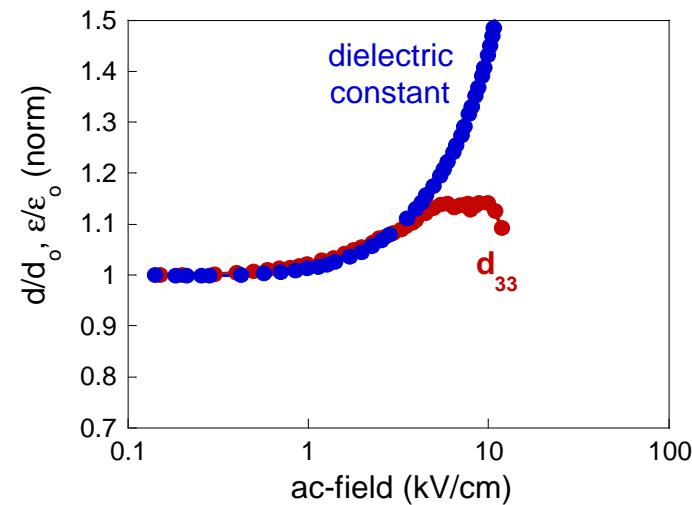
# Dielectric vs. Piezoelectric nonlinearity

tetragonal (45/55)  
<111> orientation



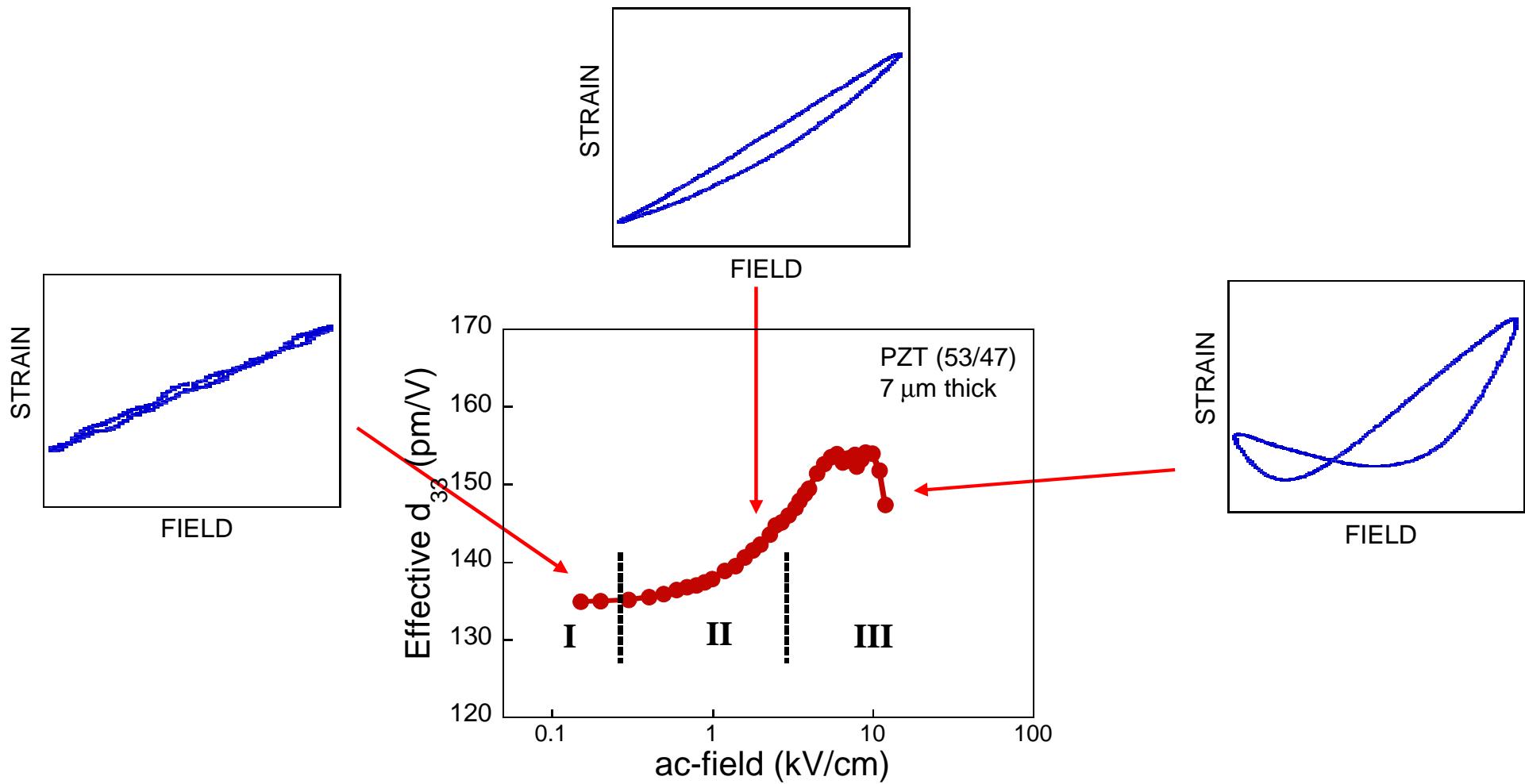
90° domains contribute  
only to dielectric constant

MPB (53/47)  
random orientation



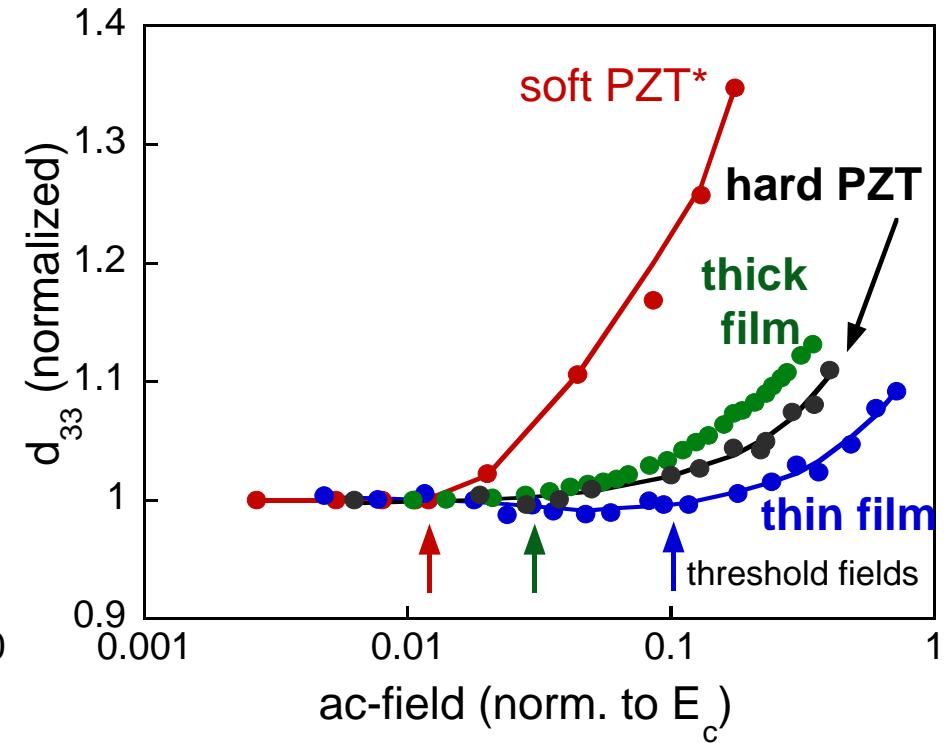
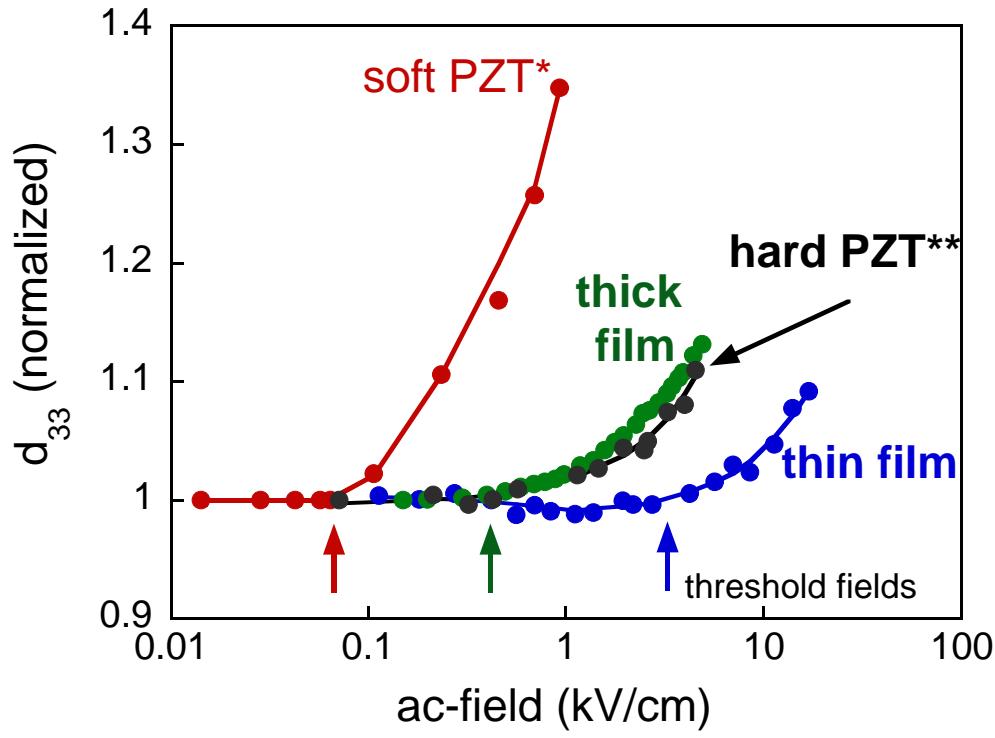
non 180° domains contribute  
to both dielectric constant and  $d_{33}$

# *Evidence of domain wall contribution*



- I - reversible domain wall motion
- II - irreversible domain motion
- III - depoling and switching

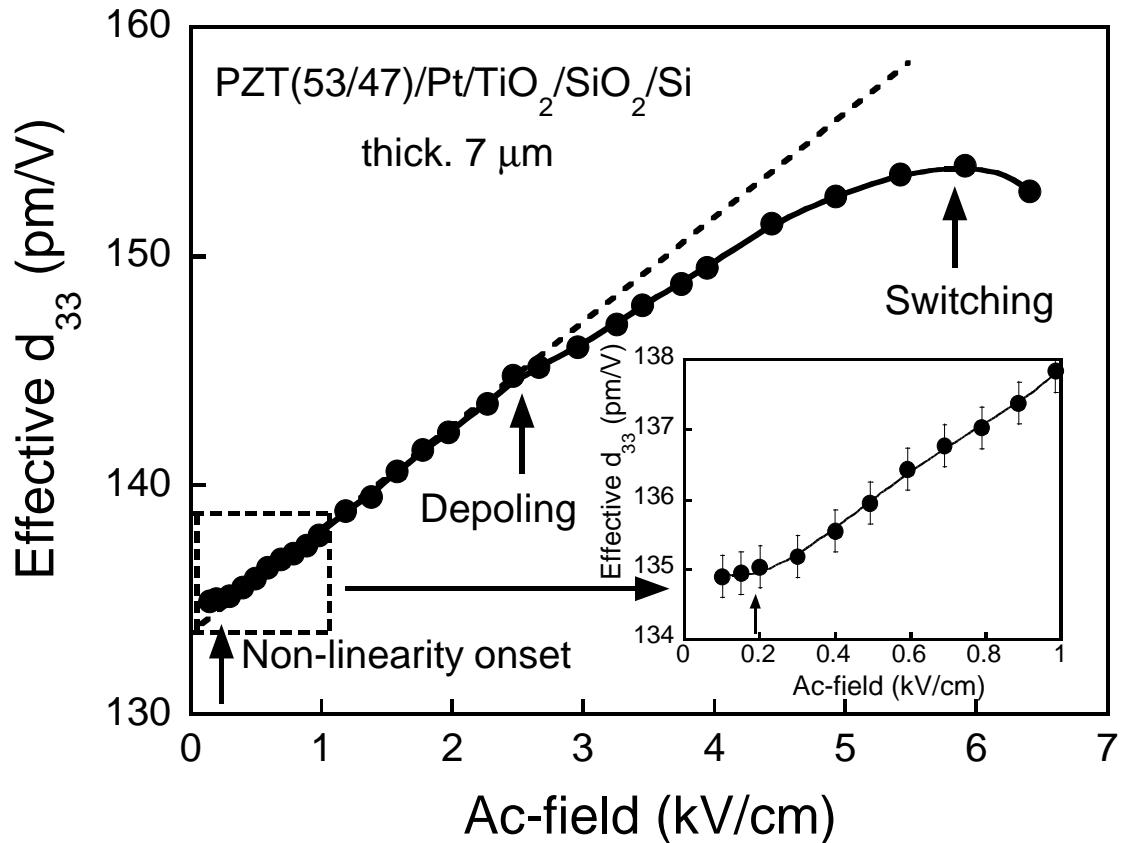
# Bulk ceramics vs. thin/thick films



\*Q.M.Zhang et al. JAP **64**, 6445 (1988), \*\*V.Mueller and Q.M.Zhang, JAP **83**, 3754 (1998)

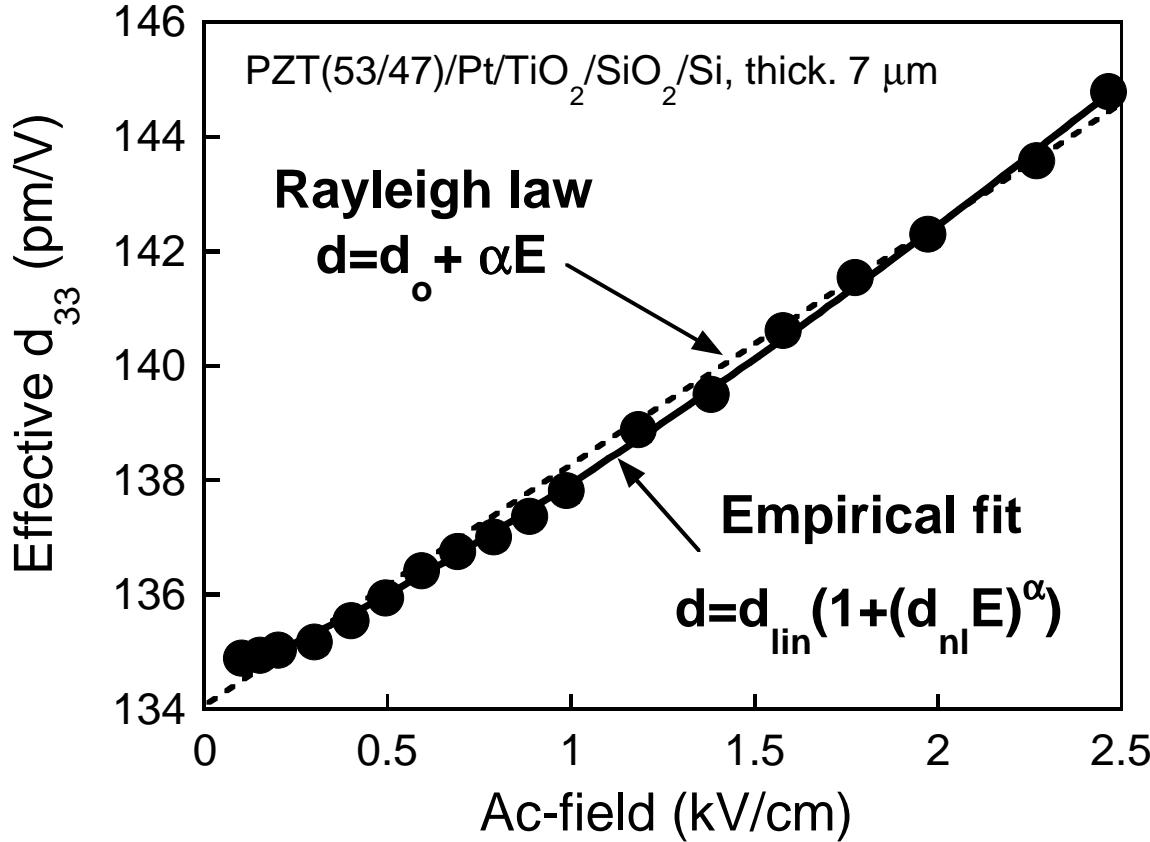
- non-linear piezoelectric properties are similar in thick films and hard PZT
- threshold field for irreversible domain wall motion is very high in thin films ( $\approx 0.1E_c$ )

# *Non-linearity in thick films*



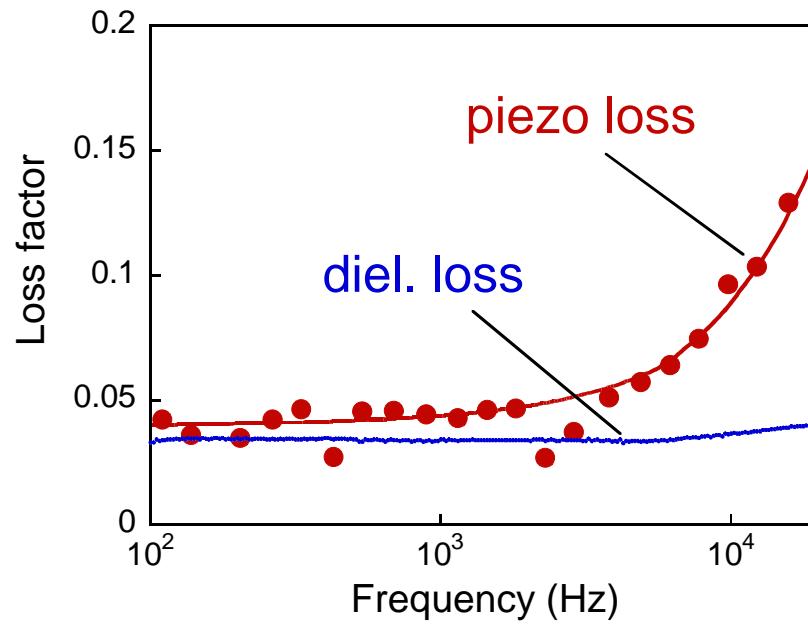
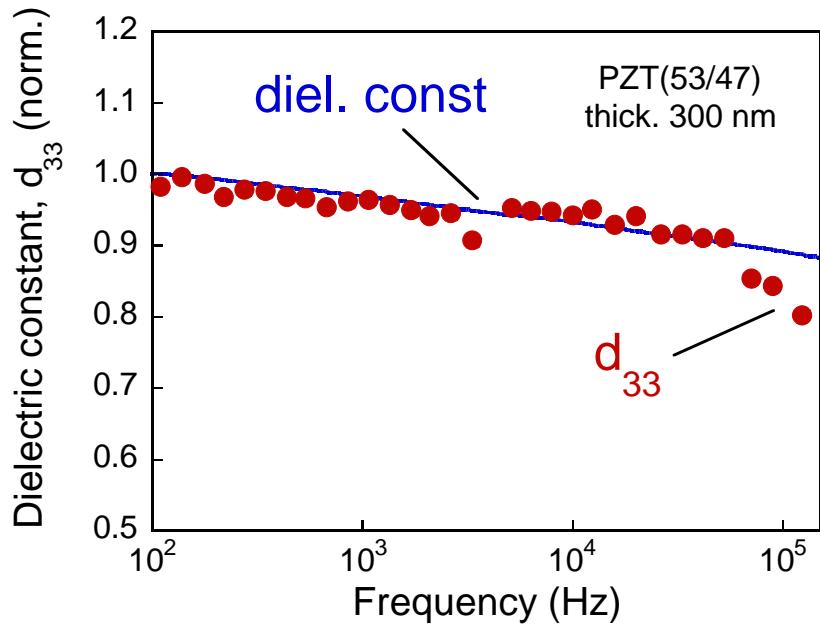
- Onset of non-linearity at ~0.2-0.3 kV/cm
- Quasi-linear response at higher fields
- Sub-linear behaviour at E>2.5 kV/cm (depoling)
- Switching at E>6 kV/cm

# Modelling of piezoresponse



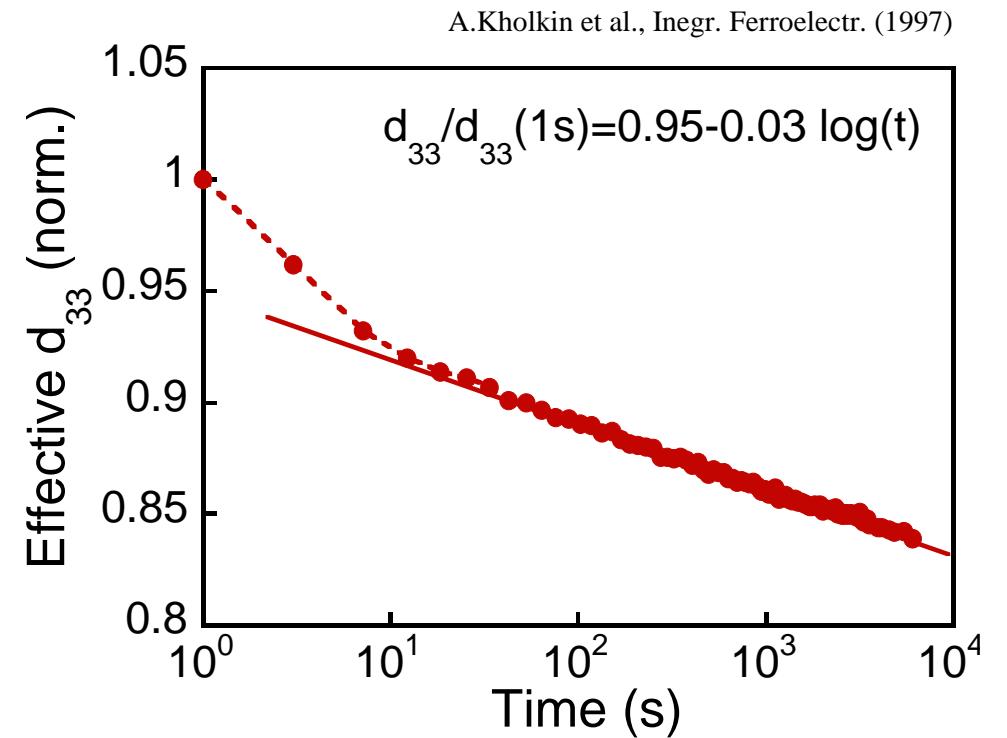
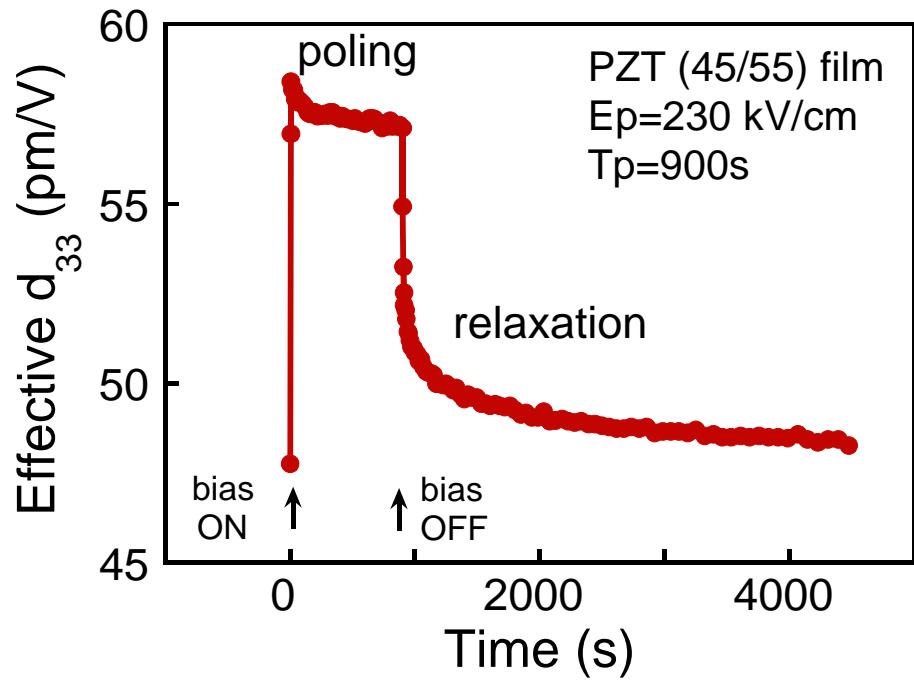
- Poor fit for Rayleigh law
- Non-linearity is better described by the empirical fit with  $\alpha=1.2$

# Frequency dispersion



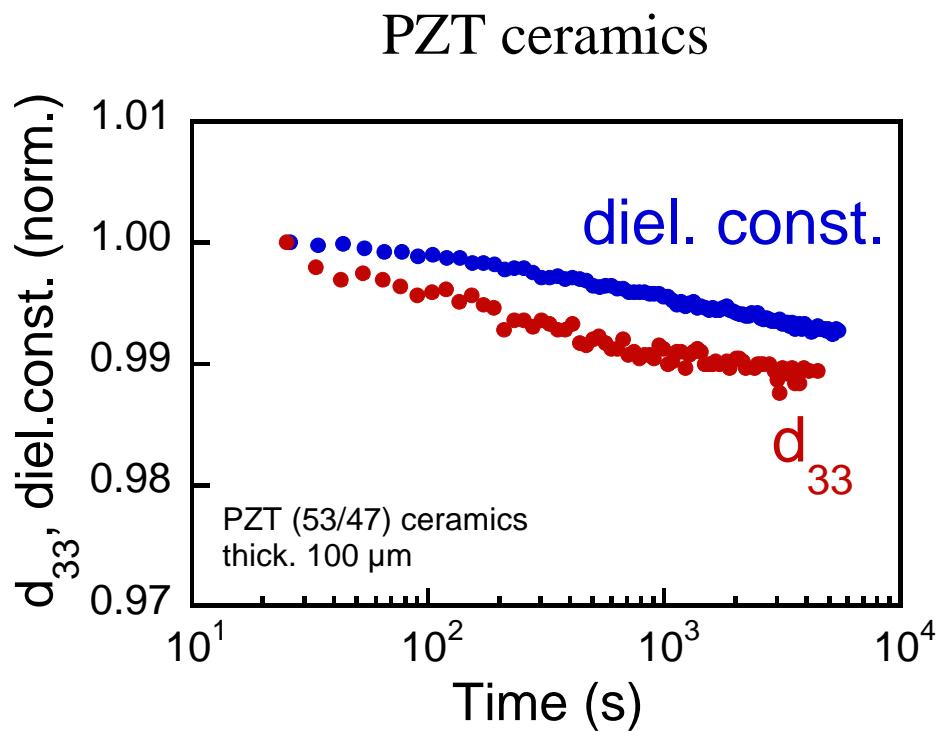
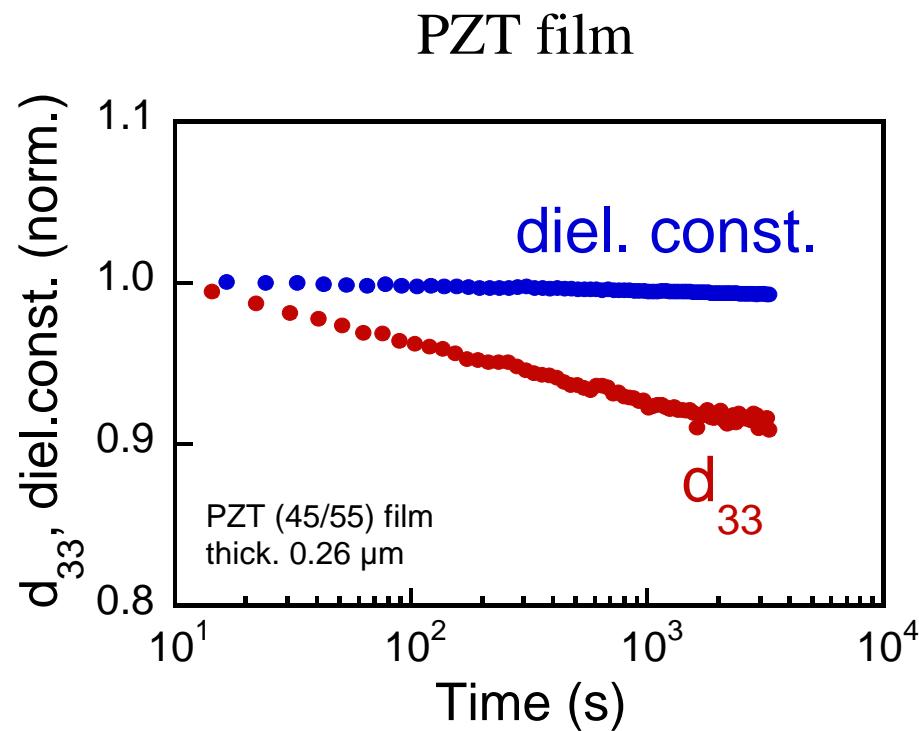
- frequency dispersion of  $d_{33}$  is close to dispersion of dielectric constant
- piezoelectric loss is close to dielectric loss

# *Aging after poling*



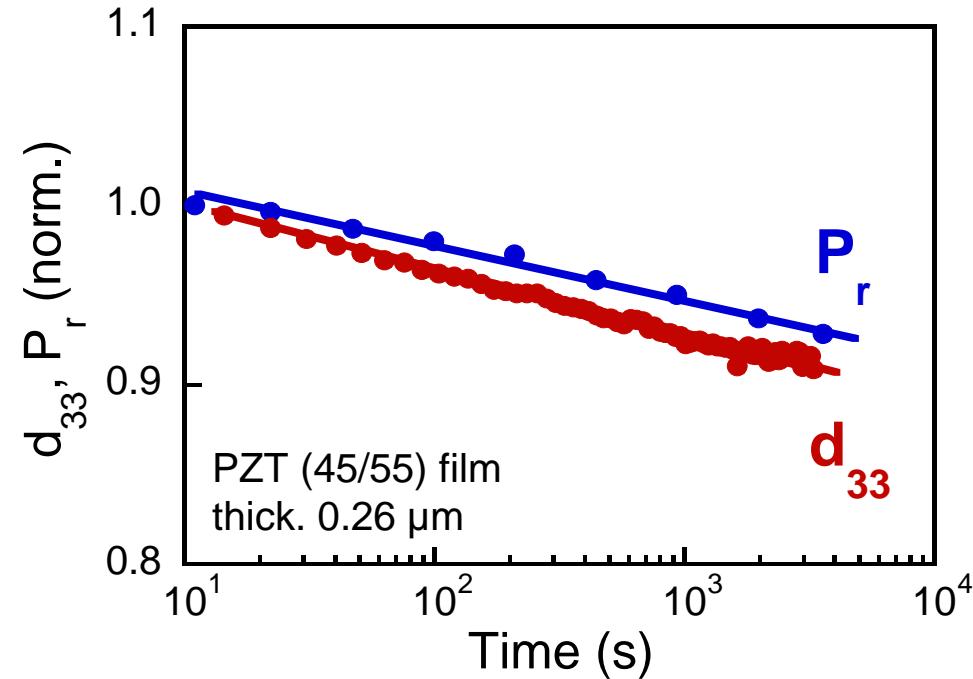
- two stages of  $d_{33}$  relaxation: fast depoling and slow logarithmic decrease
- aging rates  $\approx 2 \div 5$  %/dec depending on poling time

# *Thin films vs. Bulk ceramics*



- $d_{33}$  decays much faster than dielectric constant in PZT films
- aging of dielectric constant is comparable in films and ceramics

# *$d_{33}$ aging vs. polarization retention*

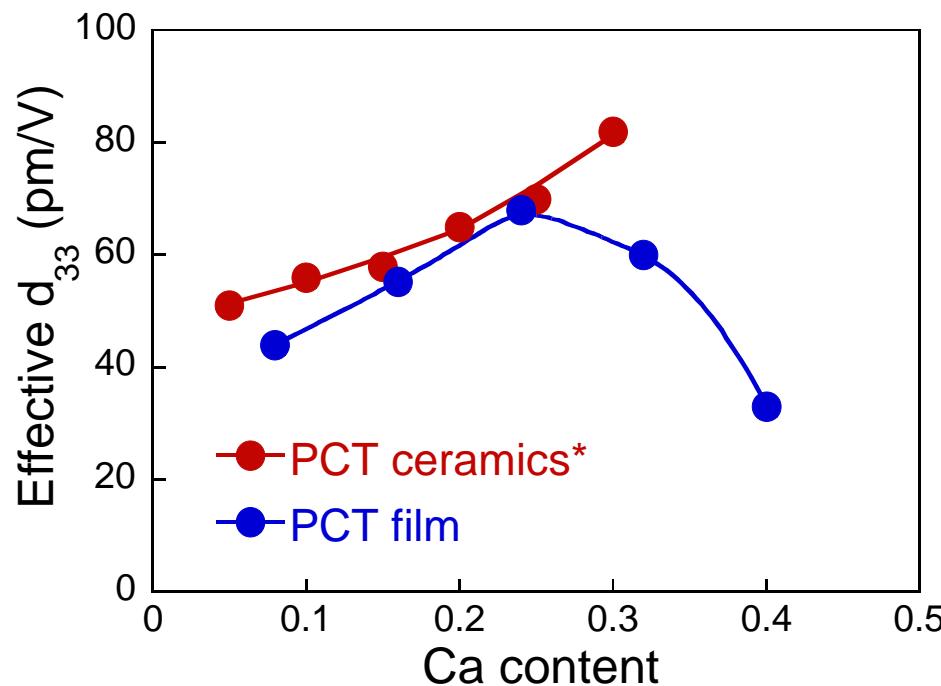


$$d_{33}(t) = 2Q_{11}\epsilon_0\epsilon(t)P_r(t)$$

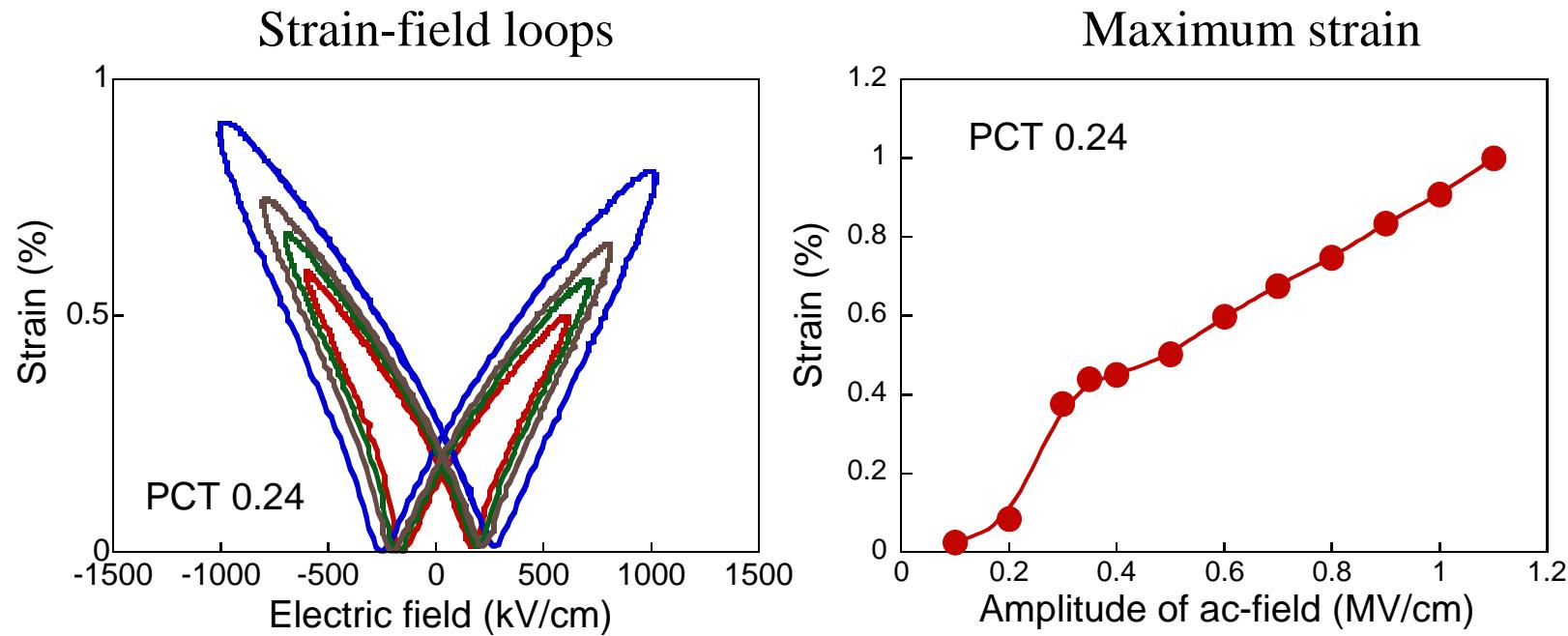
- aging of  $d_{33}$  is close to the reduction of  $P_r$  with time (retention)
- aging of  $d_{33}$  is due to depoling of PZT films rather than to the reduction of domain wall contribution

# *Case study: $PbTiO_3:Ca$*

- simple tetragonal structure
- decrease of tetragonality with Ca addition
- high  $g_{33}$  coefficients
- high piezoelectric anisotropy for low Ca content

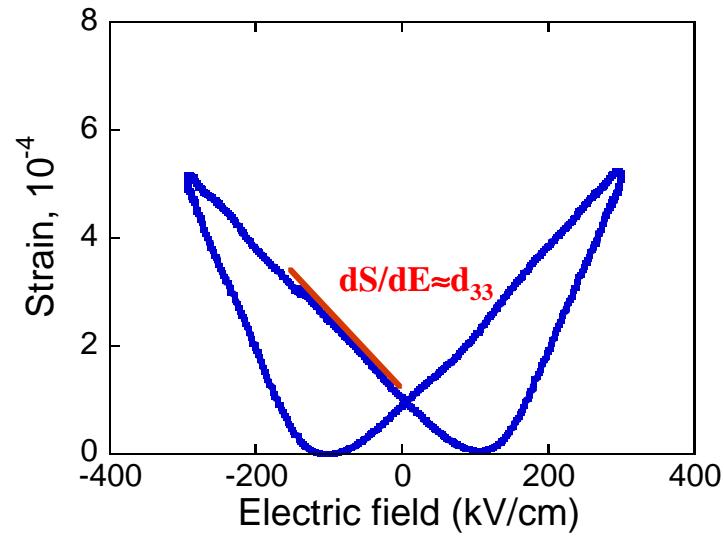
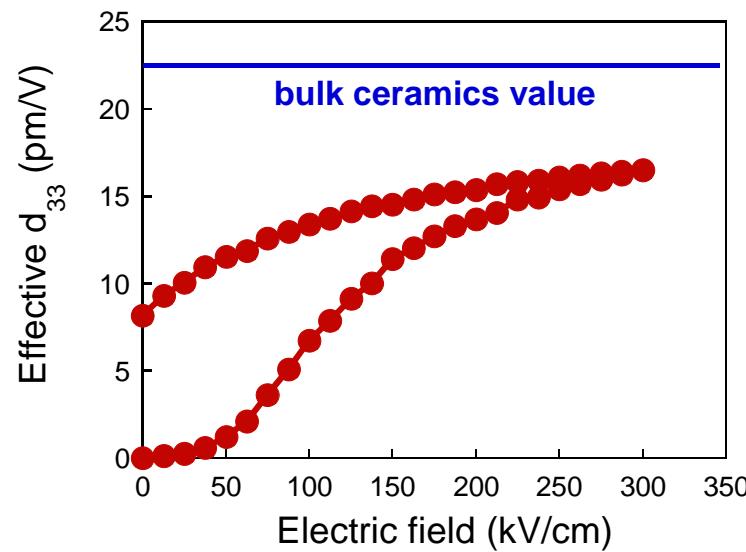


# *Giant strain in PCT films*



- High breakdown field results in a giant strain at 1 MV/cm ( $\approx 1\%$ )
- At high field strain is linear with the amplitude of ac-field

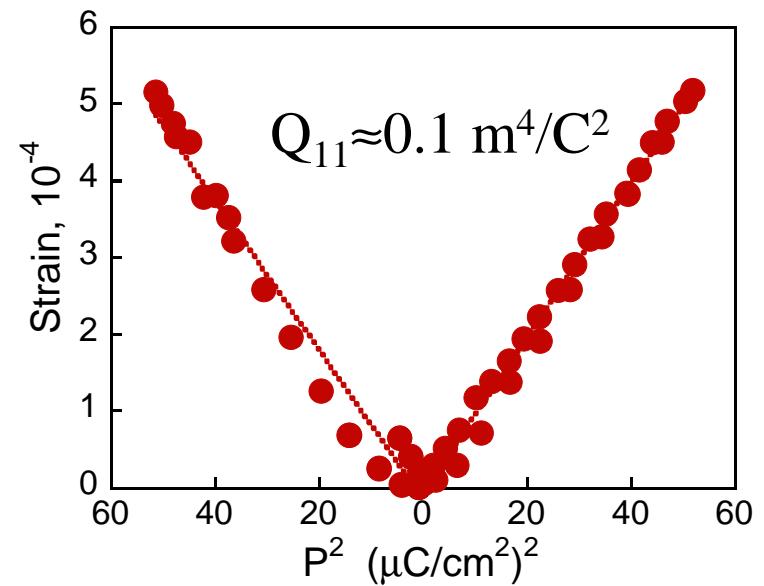
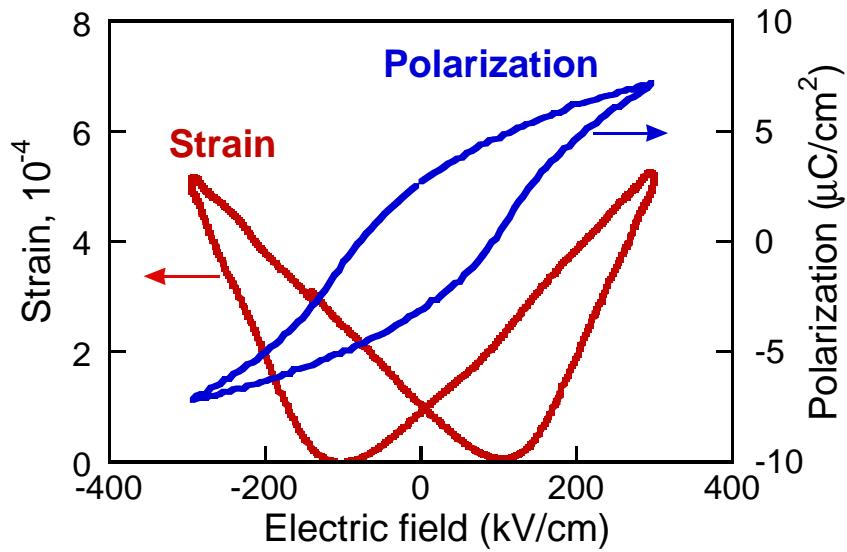
# *SBT films: absence of extrinsic strain*



- $(d_{33})_{\text{max}}$  is about 75% of bulk ceramics value
- $d_{33} \approx dS/dE$  in contrast to PZT films

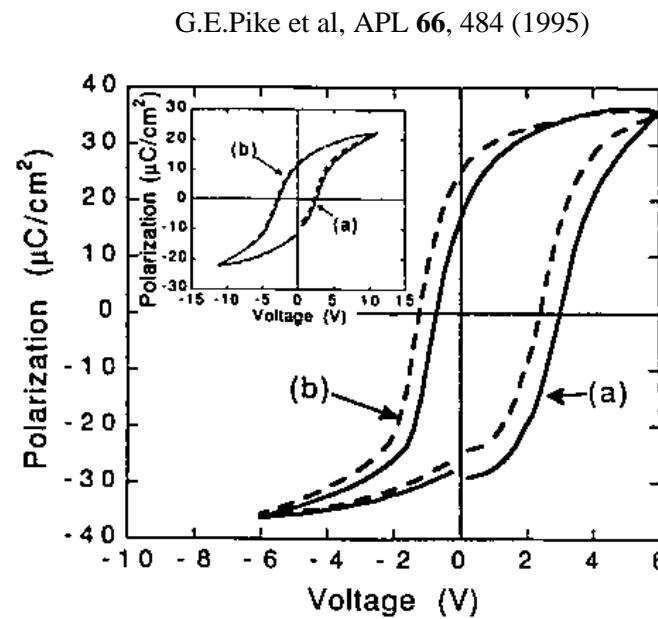
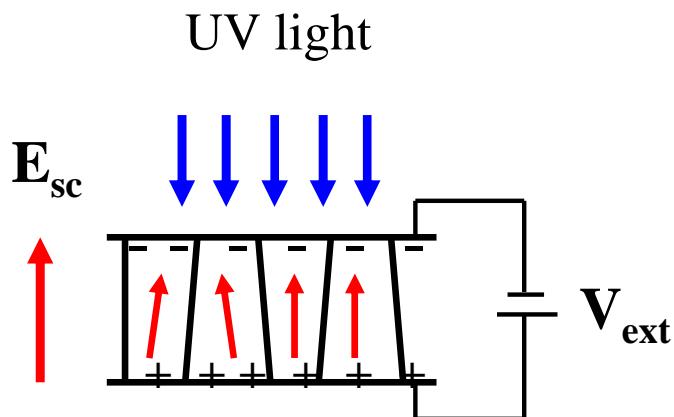
# *SBT films: absence of extrinsic strain*

$$\mathbf{S} = Q_{11} \mathbf{P}^2$$



- pure electrostrictive strain behavior at  $T \ll T_c$  ( $Q_{11} \approx 0.1 \text{ m}^4/\text{C}^2$ )

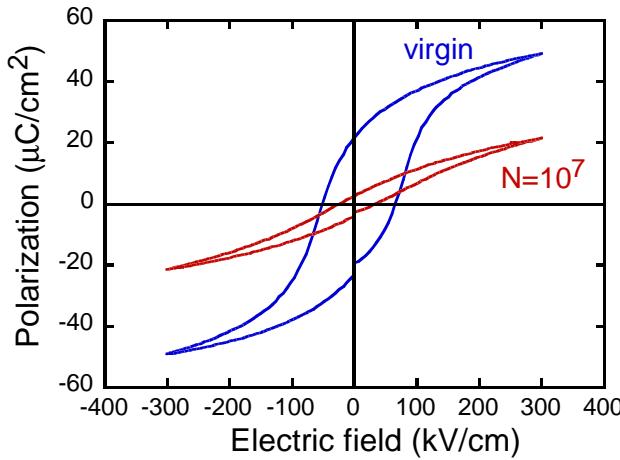
# *Domain pinning and voltage offset*



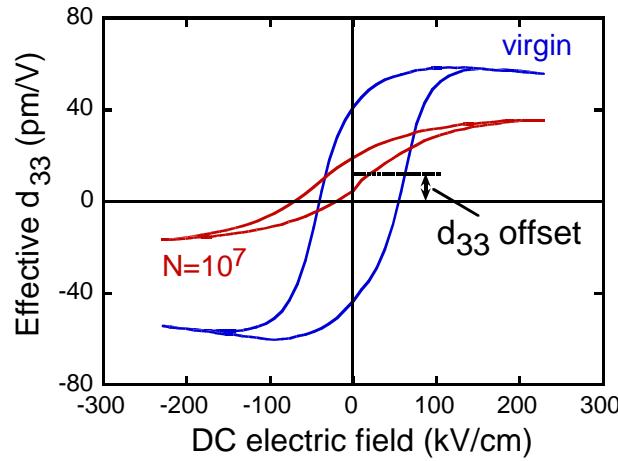
- electrons generated by UV light compensate polarization charge at the film-electrode interfaces
- voltage offset is due to the space-charge internal field

# Piezoelectric fatigue: evidence of domain pinning

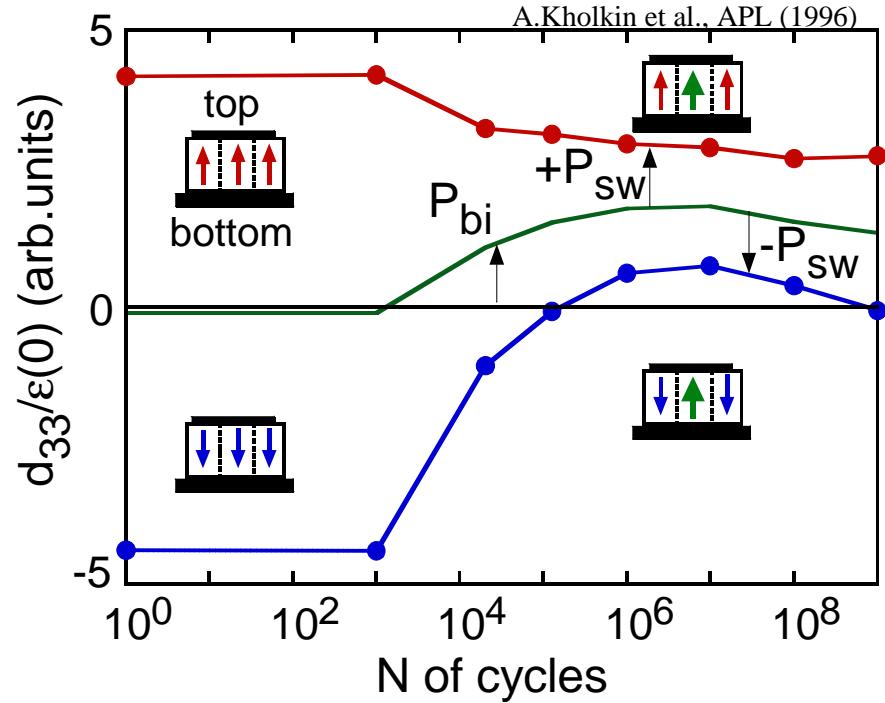
fatigue of polarization



fatigue of  $d_{33}$



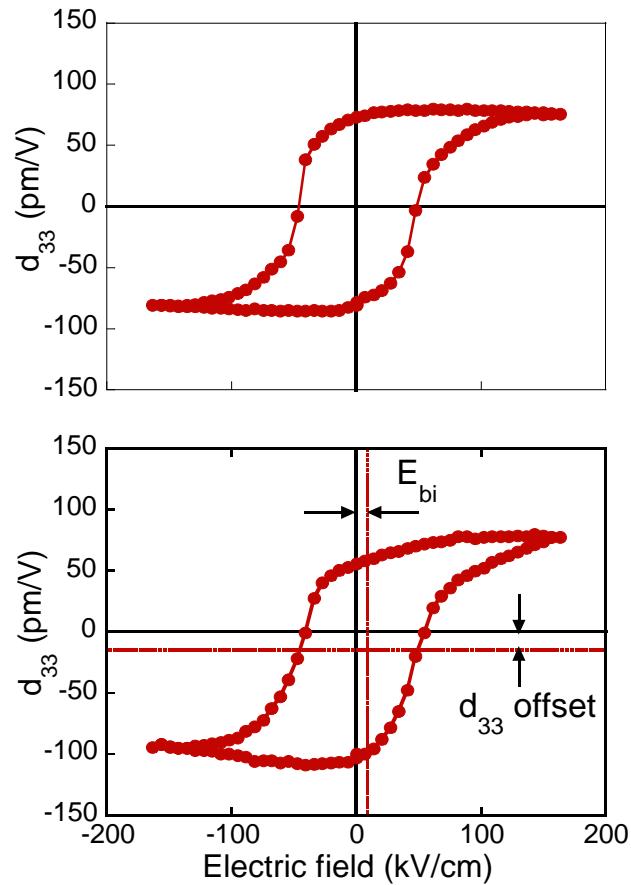
Polarization offset as a function of fatigue



$$P_{bi} = d_{33} / (2Q_{eff}\epsilon_0 \epsilon)$$

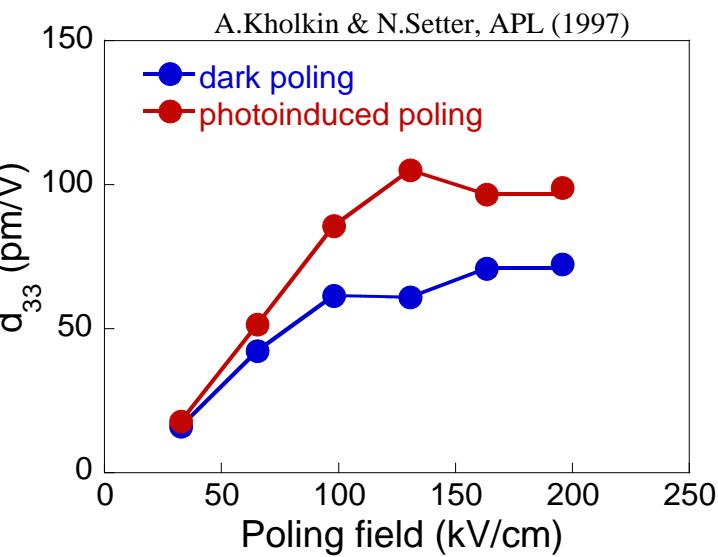
# *UV-induced poling*

$d_{33}$  loop of virgin film



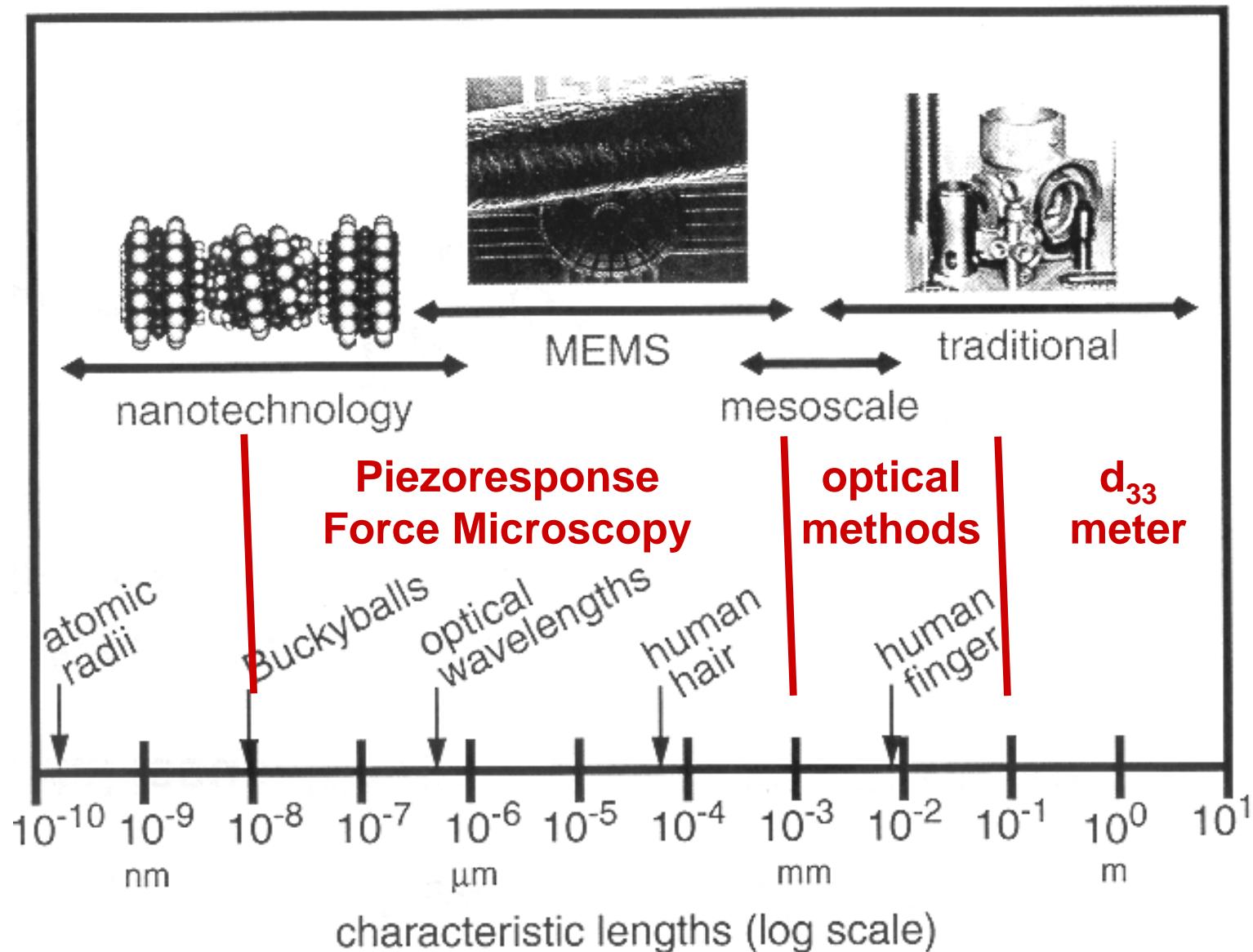
$d_{33}$  loop of illuminated film

Photoinduced poling effect

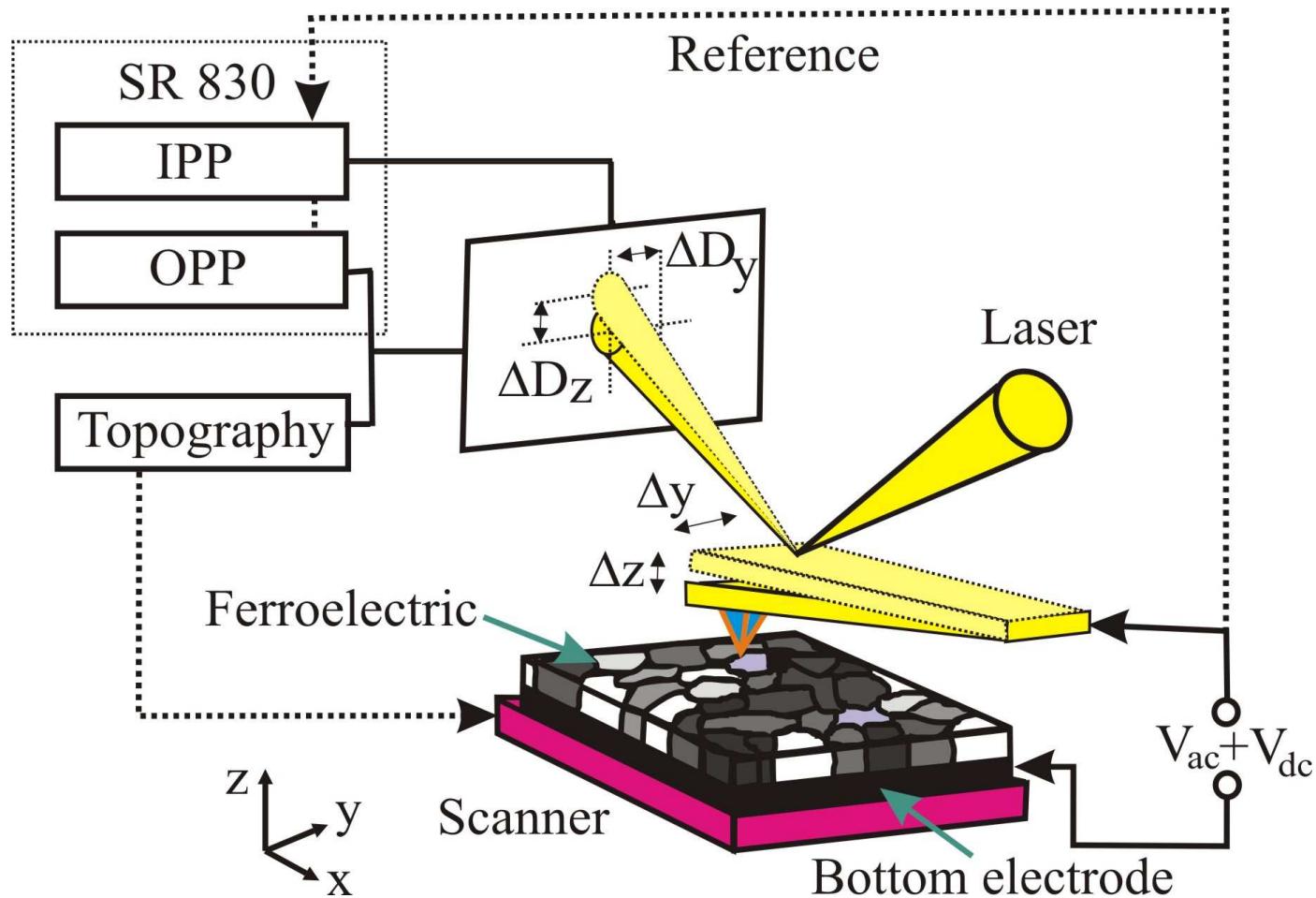


$$d_{33} = 2Q_{\text{eff}}\epsilon_0 \epsilon P_{\text{bi}}$$

# *Downscaling of actuators*

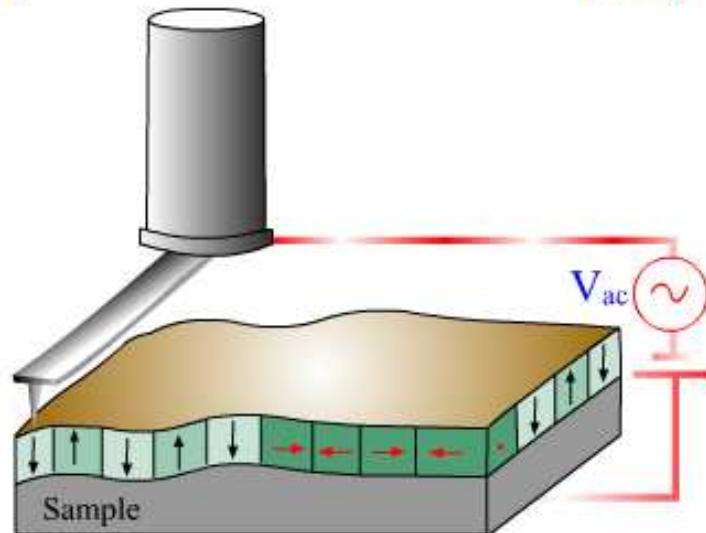


# *Local characterization: Piezoresponse Force Microscopy*

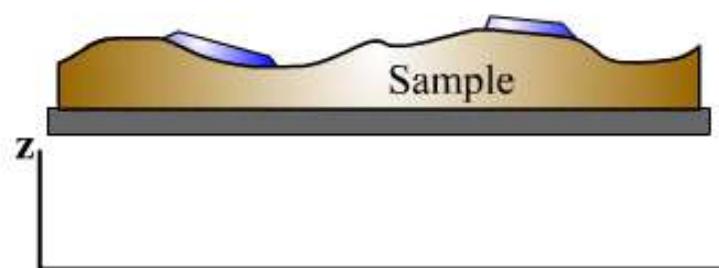
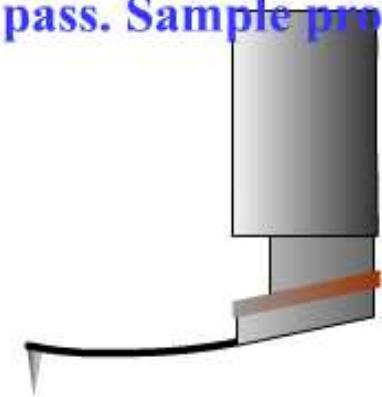


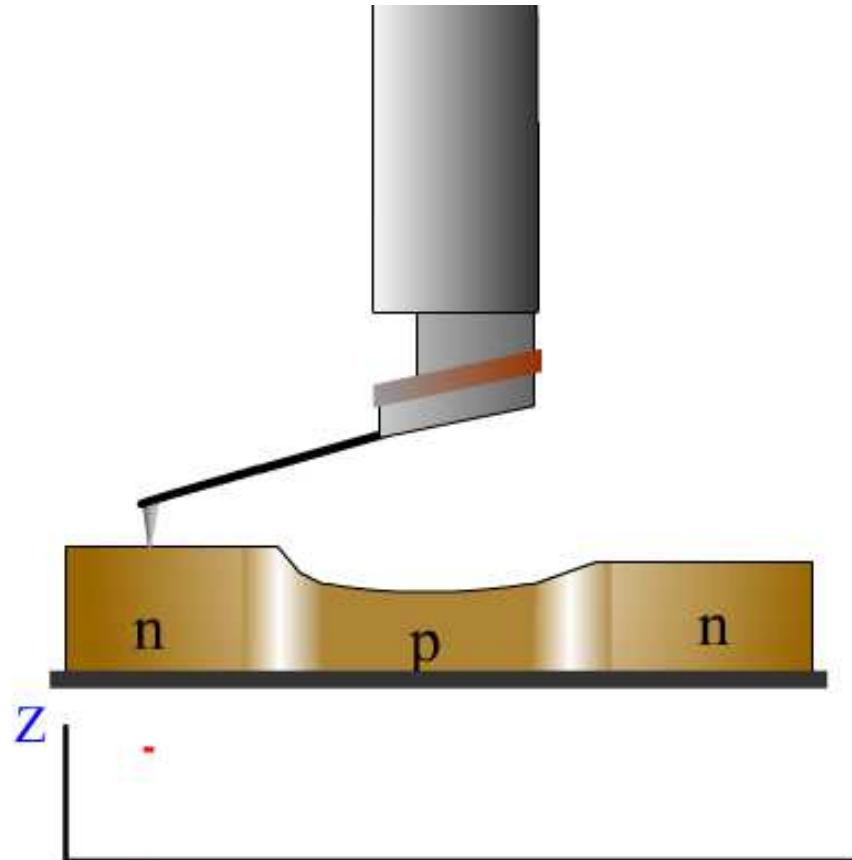
INTRO

INTRO2



## 1-st pass. Sample profile acquisition.



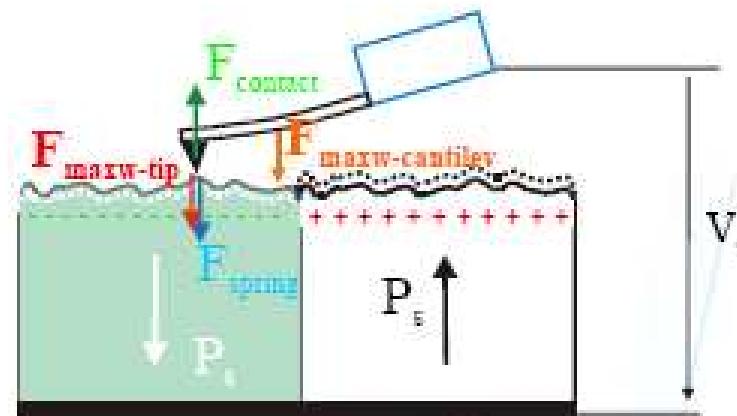


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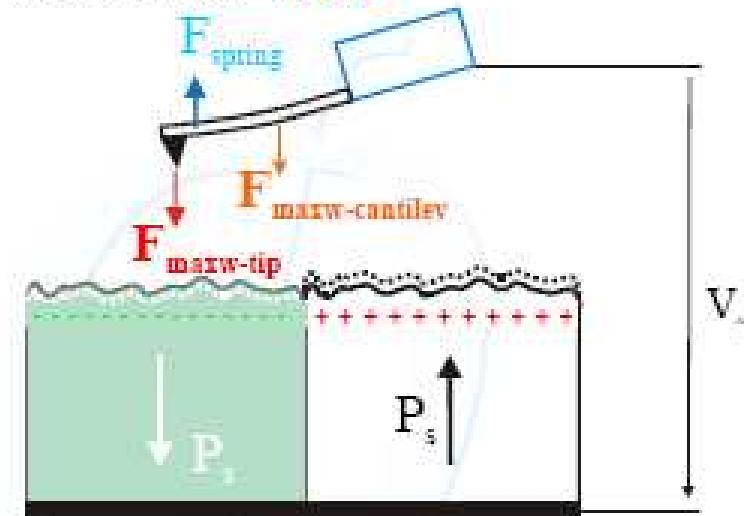
[www.ntmdt.com](http://www.ntmdt.com)

# PFM vs. EFM

Contact - Piezoresponse



Non-Contact - EFM

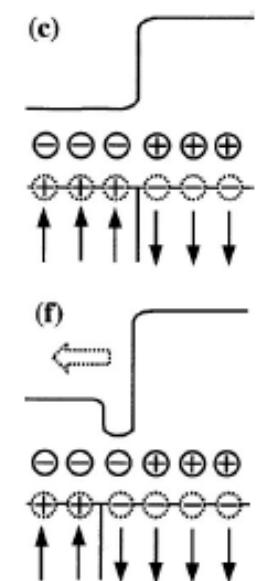
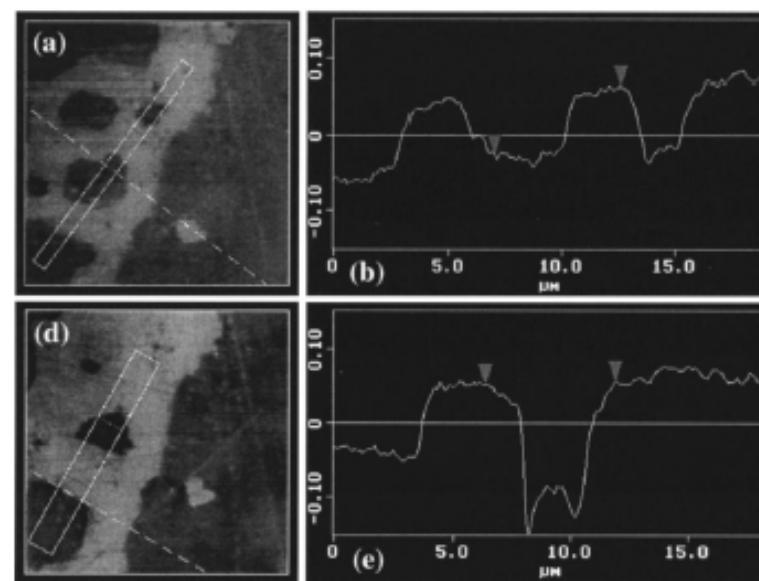
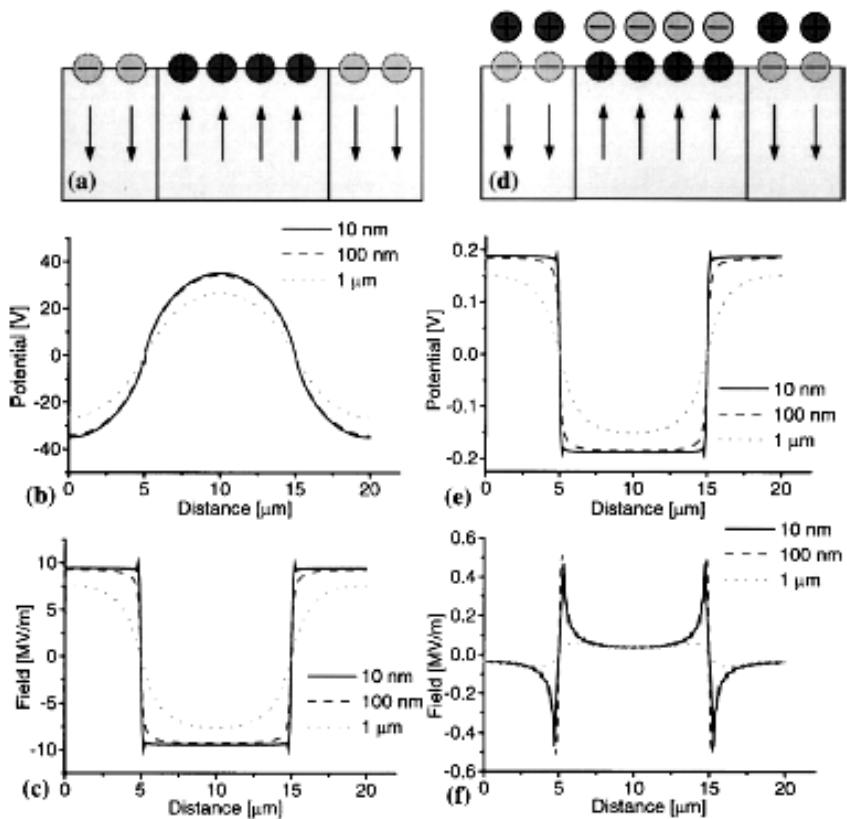


Topography

z - Piezoresponse

EFM-Surface charge

# *KFM and EFM for domain imaging*

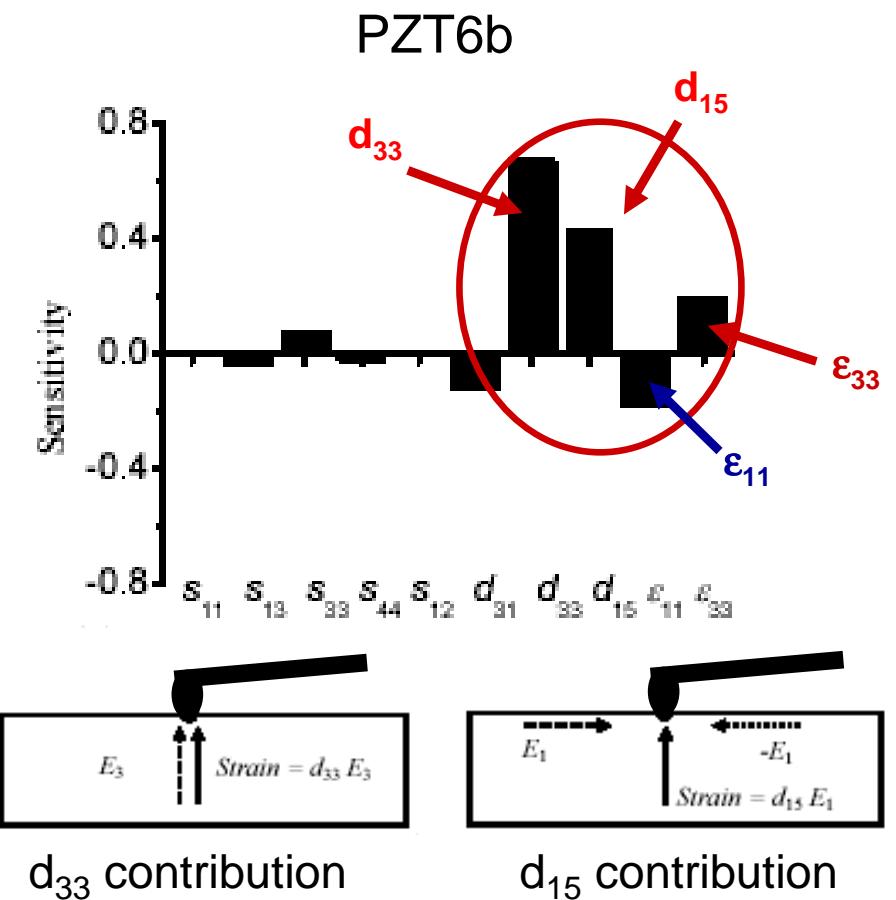
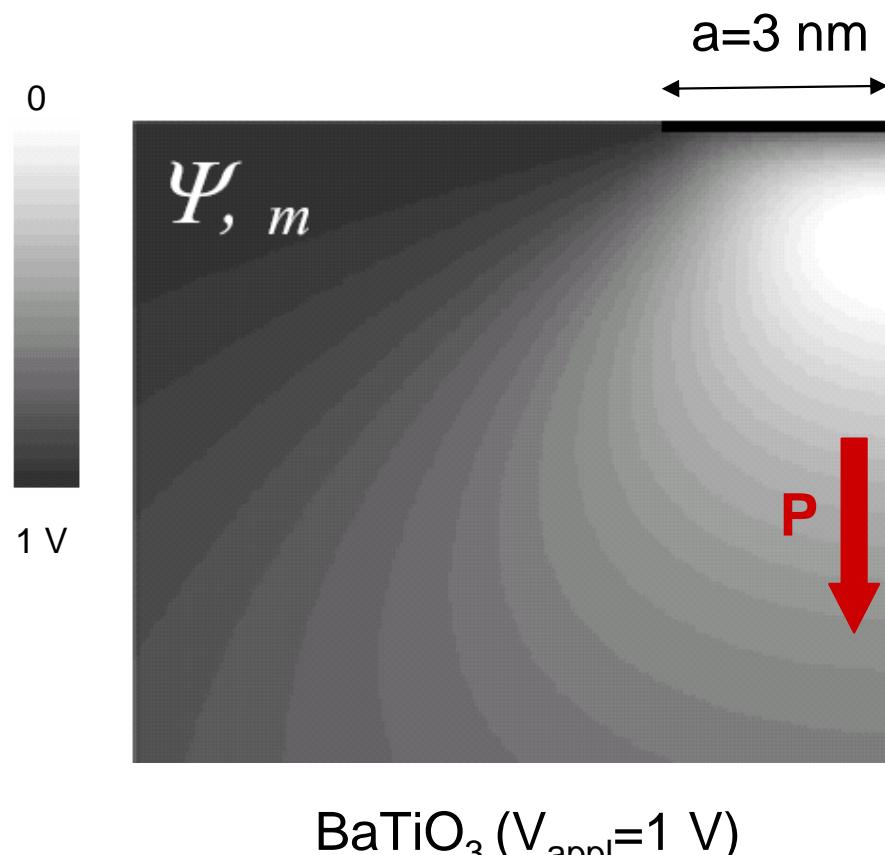


Kalinin & Bonnell PRB (2001)

# *Advantages/disadvantages of PFM*

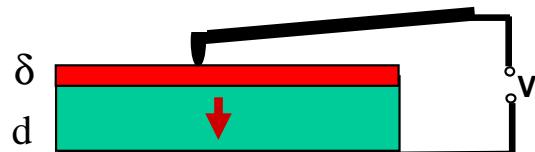
- Much higher resolution: limited by tip diameter
- Control of polar state: new applications of ferroelectrics:  
e.g., nanolithography and self-assembly routines
- Local domain engineering and interaction with defects
- Visualization of small domains impossible in the past
- Multiferroics and magnetoelectric interaction on the nanoscale
- Slow acquisition
- Difficult to get reliable absolute values of  $d_{33}$
- Not easy interpretation: electrostatics

# Potential distribution and effective $d_{33}$



Kalinin et al, Phys. Rev. B (2004)

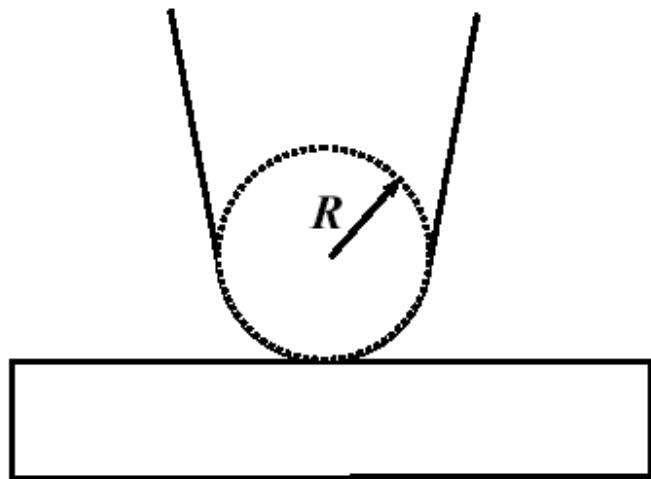
Influence of interface layer



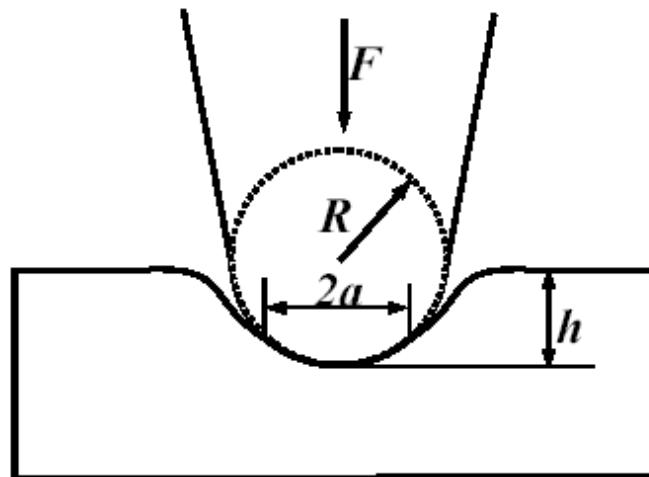
$$(d_{33})_{\text{eff}} \propto \frac{A_{1\omega}}{V_{\text{ac}}} = \frac{A_{1\omega}}{1 + \frac{\delta \epsilon}{d \epsilon_\delta}}$$

# *Effect of tip-surface interaction*

Weak indentation



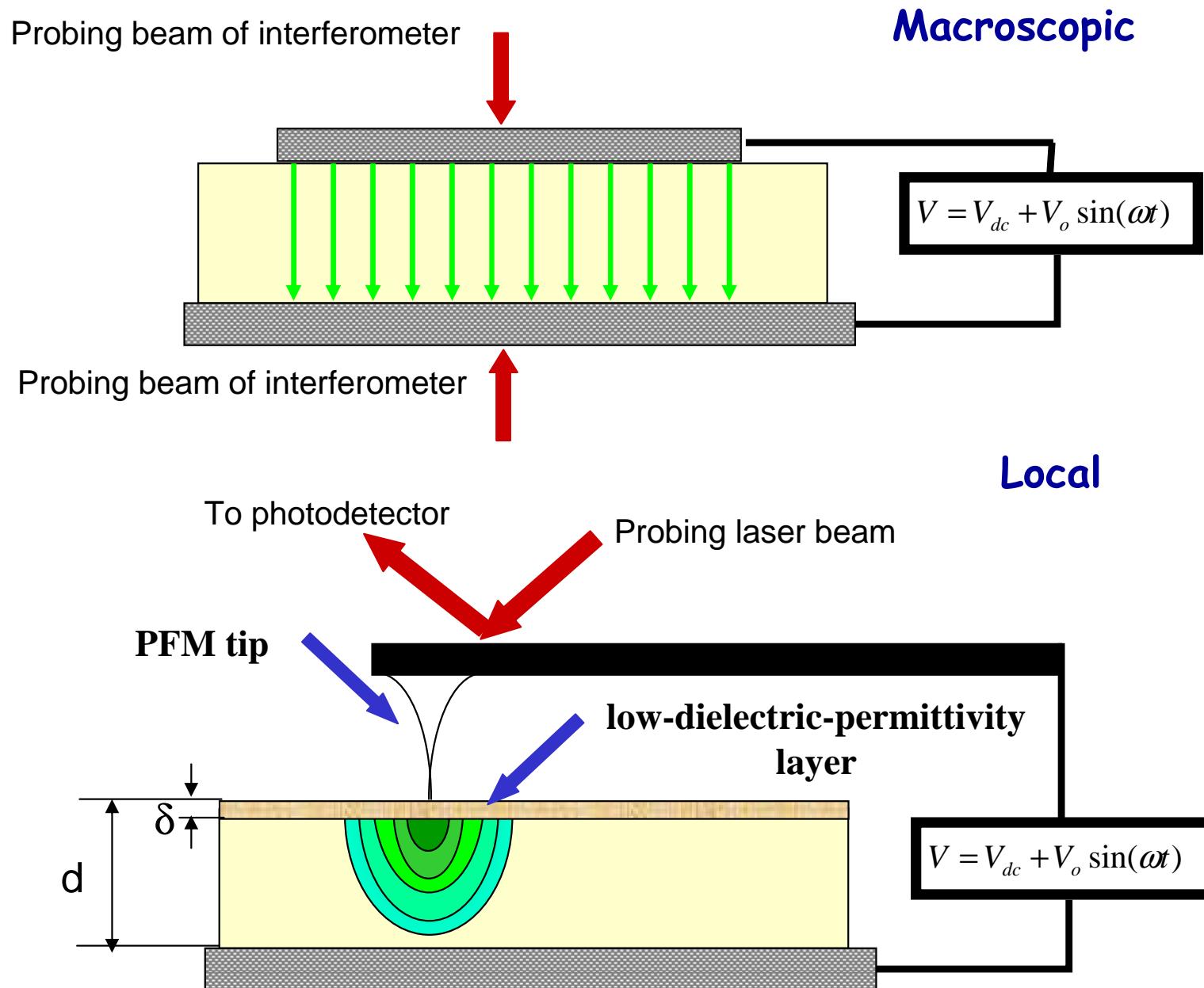
Strong indentation



- Small contact area  $\ll R$ .
- Electrostatic signal dominates electromechanical one.  $F = \frac{V^2}{2} \frac{\partial C}{\partial z}$
- Distribution of E-field can be calculated using image charge models.

- Medium contact area  $a \leq R$ .
- Electromechanical signal dominates electrostatic (Maxwell stress).
- Coupled **electrostatic + elastic + electroelastic** problems should be solved!

# *Local vs. macroscopic measurements*



# Piezoelectric hysteresis

Strain in monodomain ferroelectric with centrosymmetric paraelectric phase

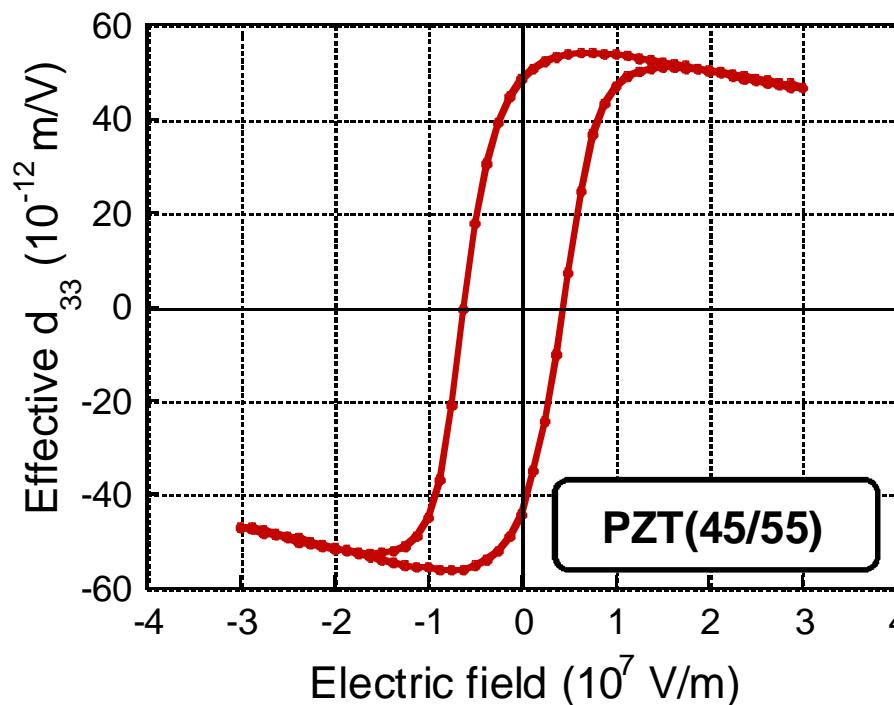
$$S = QP^2(E) = Q[P_s + P_{ind}(E)]^2 = QP_s^2 + 2Q\epsilon_o \epsilon P_s E + QP_{ind}^2$$

For small ac-field  $E = E_{dc} + E_o \sin(\omega t)$

1st harmonics of ac-strain  $S_{1\omega} = \underbrace{2Q\epsilon_o \epsilon P_s E_o}_{d_{33}=\text{const}} \sin(\omega t)$

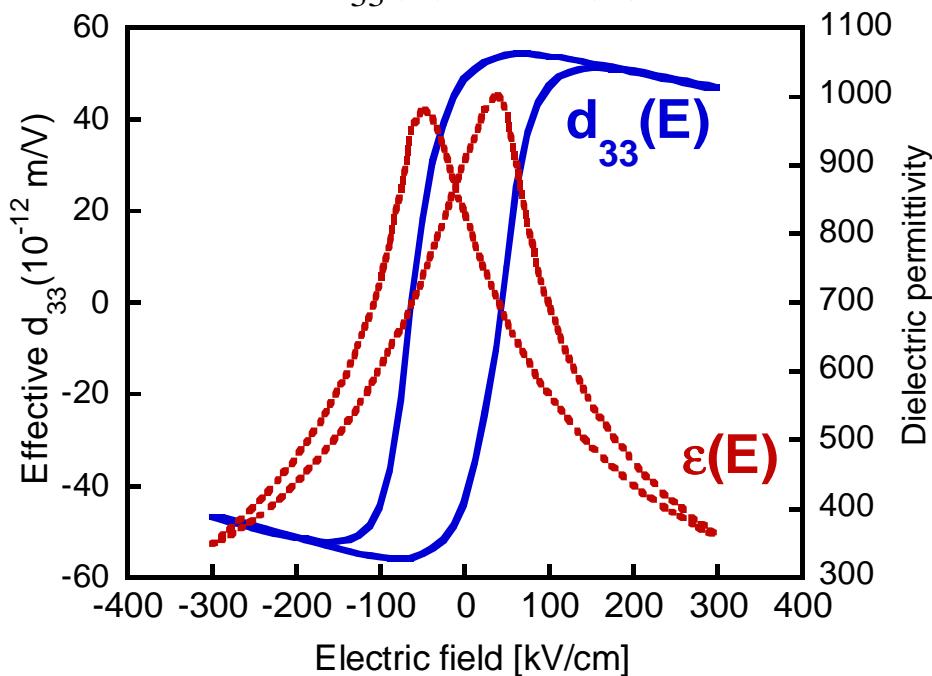
If polarization is changed by  $E_{dc}$ ,  $d_{33}=f(E_{dc})$  and can be formally rewritten as

$$d_{33}(E_{dc}) = 2Q\epsilon_o \epsilon(E_{dc}) P(E_{dc})$$

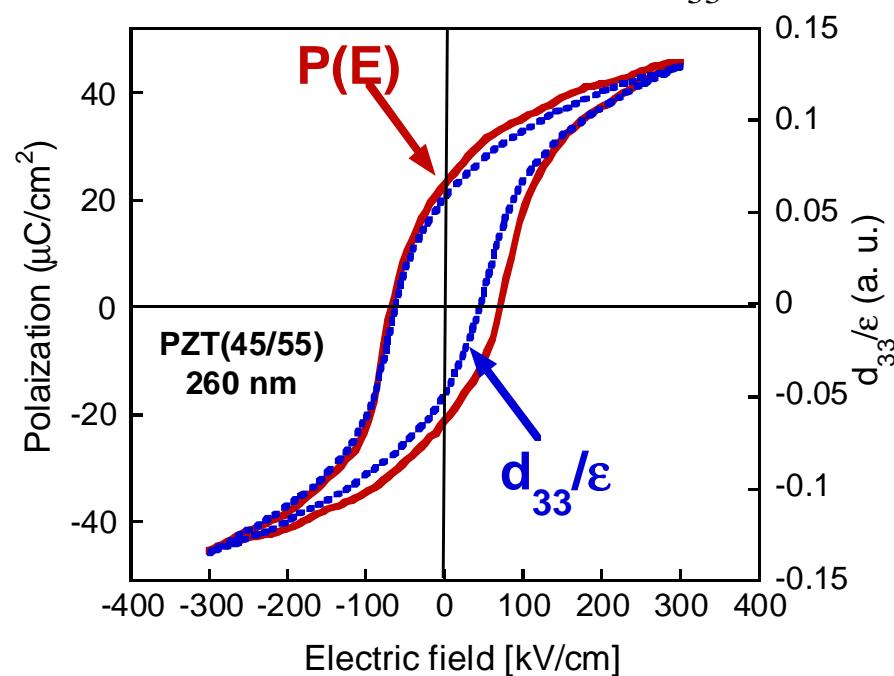


# *P(E) reconstruction*

Simultaneous measurements  
of  $d_{33}(E)$  and  $\epsilon(E)$



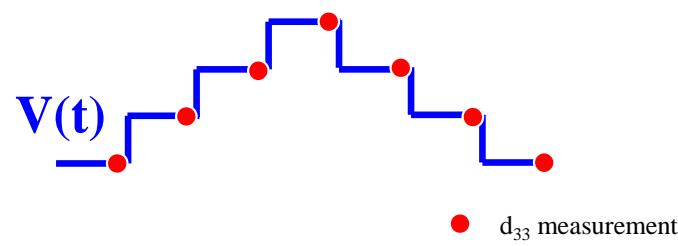
Comparison of  $P(E)$   
and reconstructed from  $d_{33}/\epsilon$



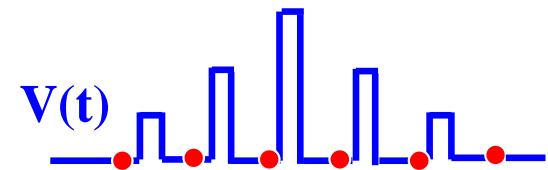
- Polarization hysteresis is close to  $d_{33}/\epsilon(E)$  hysteresis
- $d_{33}(E)$  shape is determined by both  $\epsilon(E)$  and  $P(E)$

# Influence of electrostatics

Two modes of voltage application



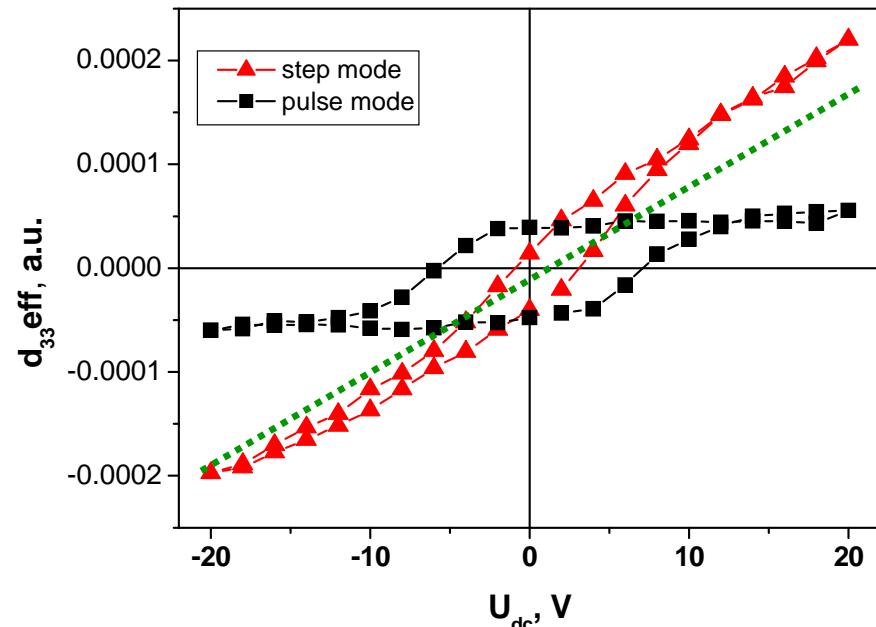
Step mode

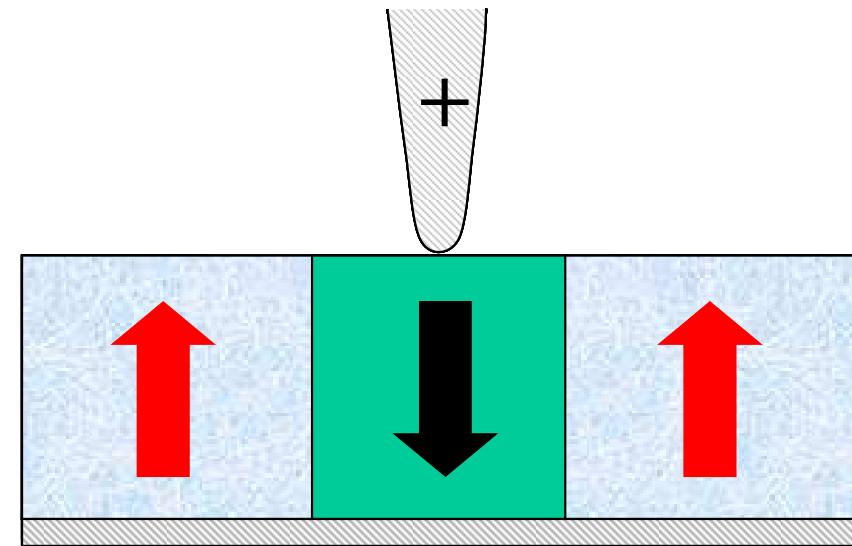
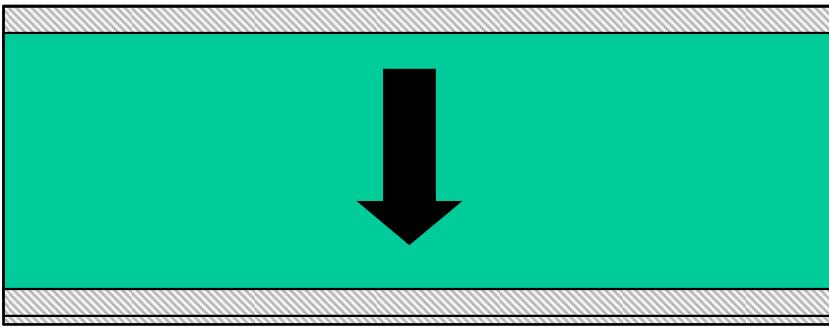


Pulse mode

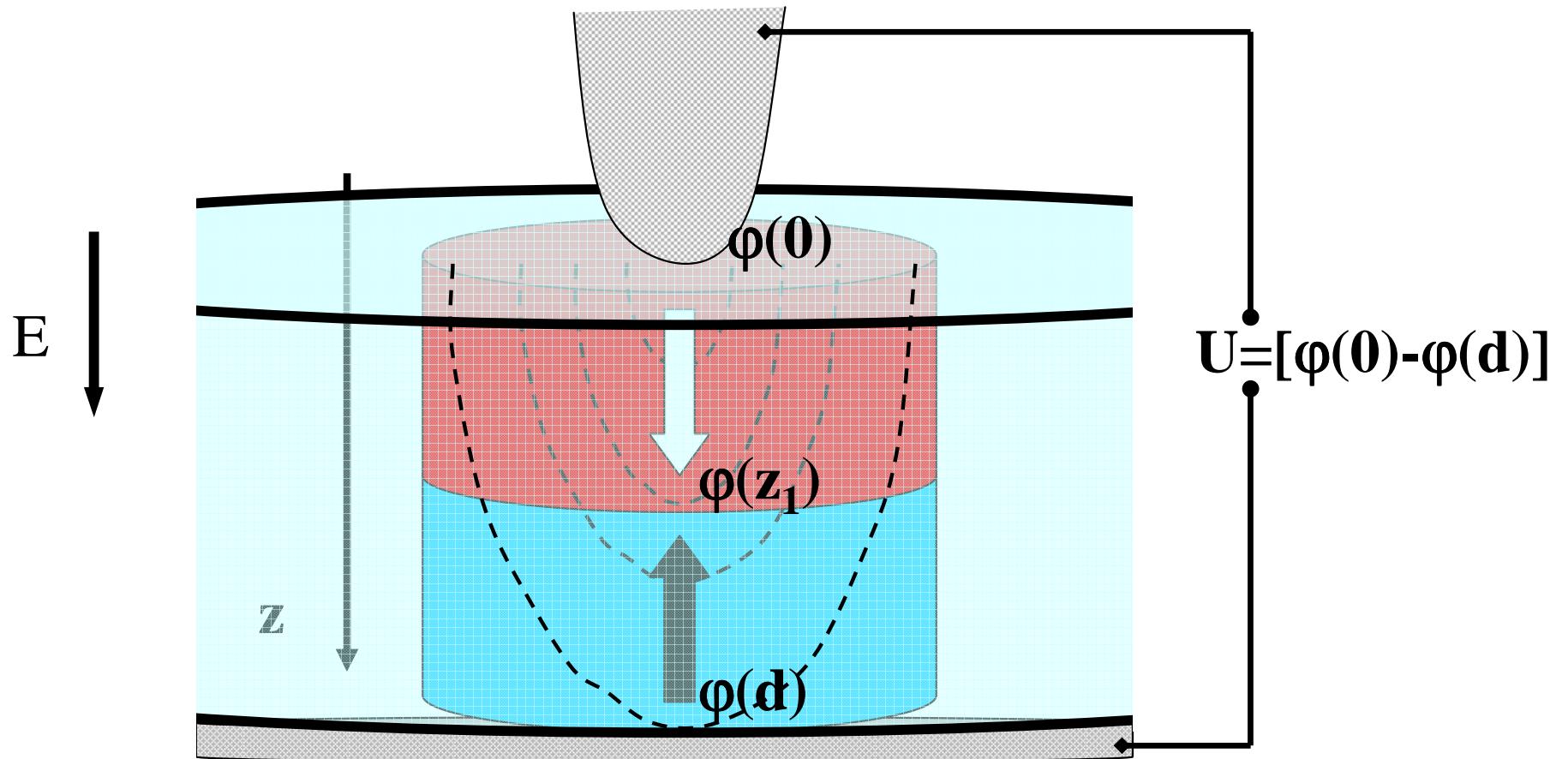
$$A \cos \varphi = -\underbrace{\frac{K}{k_{lever}} \frac{\partial C}{\partial z} (V_{dc} + V_k)}_{\text{Electrostatic term}} \pm d_{33} V_o \quad \text{Hong et al., JAP (2001)}$$

- Electrostatic interaction distorts the shape of the hysteresis in step mode and simulates polarization offset if  $V_k \neq 0$





# Origin of piezosignal



$$\begin{aligned}\Delta l &= \int_0^d d_{33}(z)E(z)dz = d_{33} \int_0^{z_1} E(z)dz - d_{33} \int_{z_1}^d E(z)dz = \\ &= d_{33}[\varphi(0) - 2\varphi(z_1) + \varphi(d)] = d_{33}\{U - 2[\varphi(z_1) - \varphi(d)]\}\end{aligned}$$

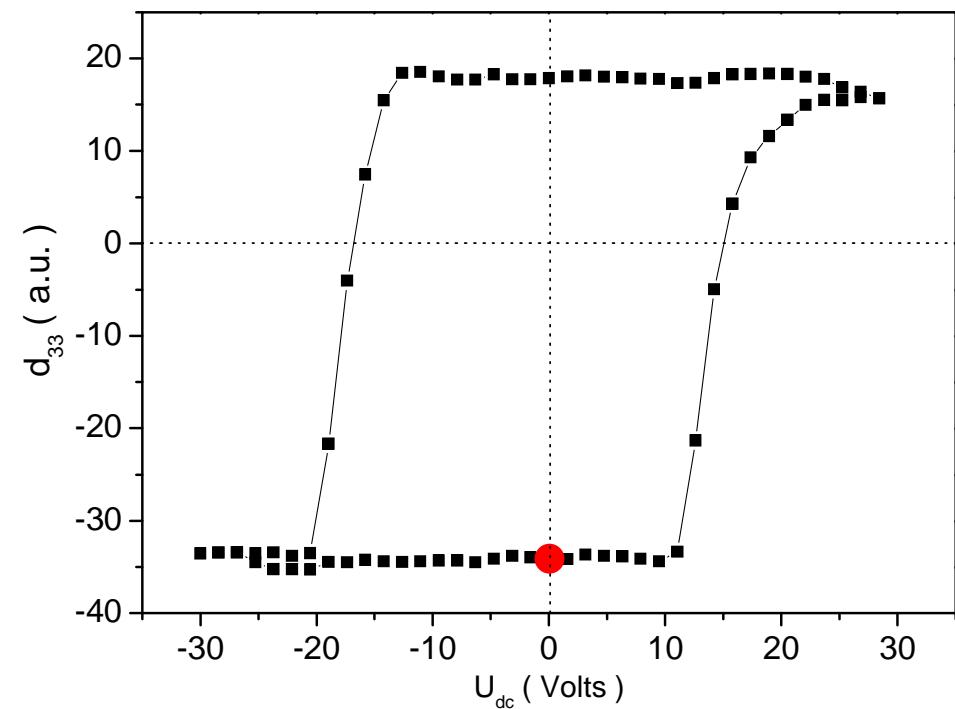
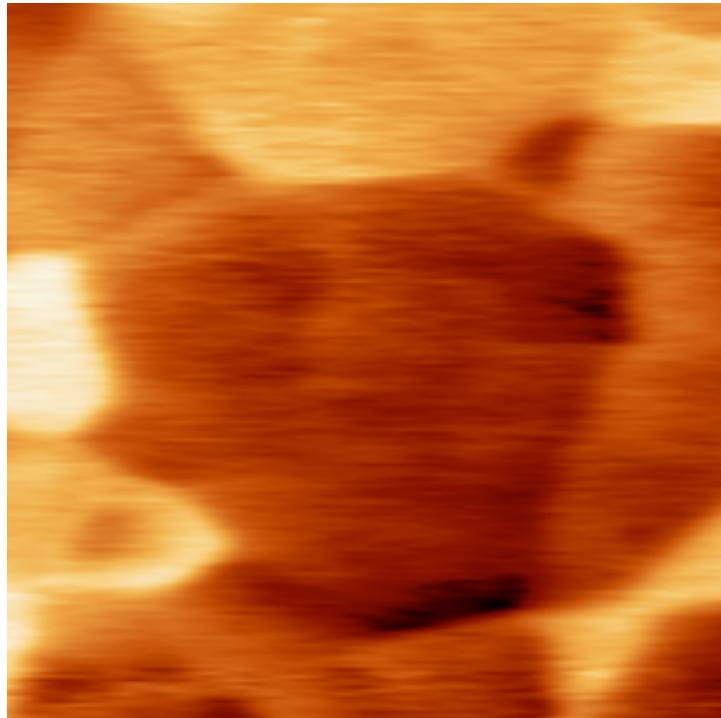
$$\Delta l = (d_{33})_{eff} U \Rightarrow (d_{33})_{eff} = d_{33} \left\{ 1 - \frac{2[\varphi(z_1) - \varphi(d)]}{U} \right\}$$

Wu et al., Nanotechnology 16 (2005)

$$E_c \text{ at } \varphi(z_1) = \frac{\varphi(0) + \varphi(d)}{2}$$

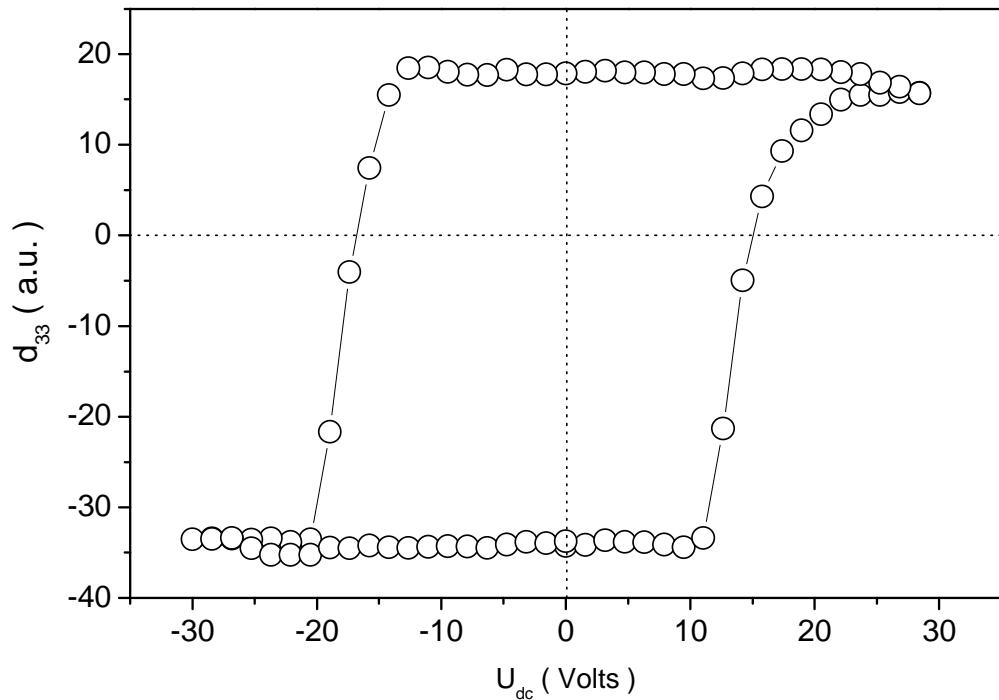
$$d_{eff} \propto \alpha \gamma d_{33} \left[ 1 - \frac{2R}{R + \frac{P_s U_{dc} \ln(\epsilon)}{\epsilon \sigma_w}} \right]$$

# *Hysteresis in PZT films*

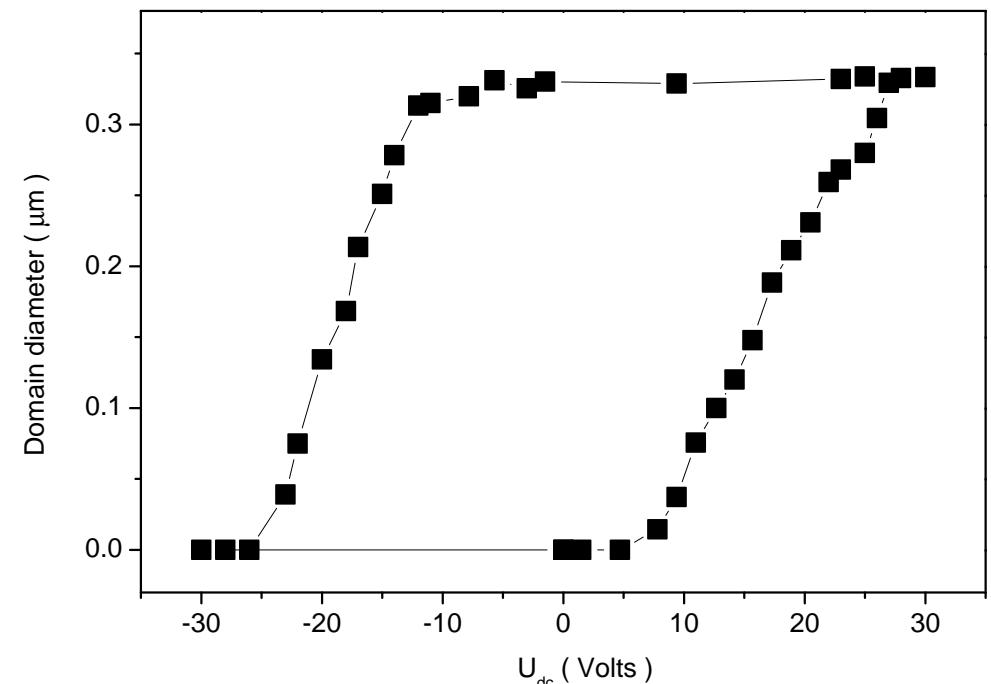


# Vertical vs. horizontal hysteresis

*piezoresponse vs.  $V_{dc}$*

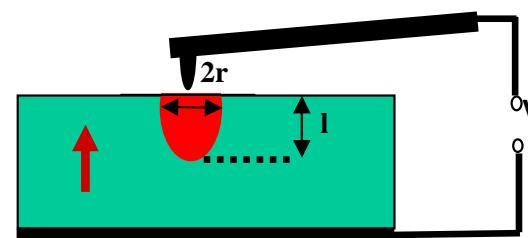


*domain size vs.  $V_{dc}$*

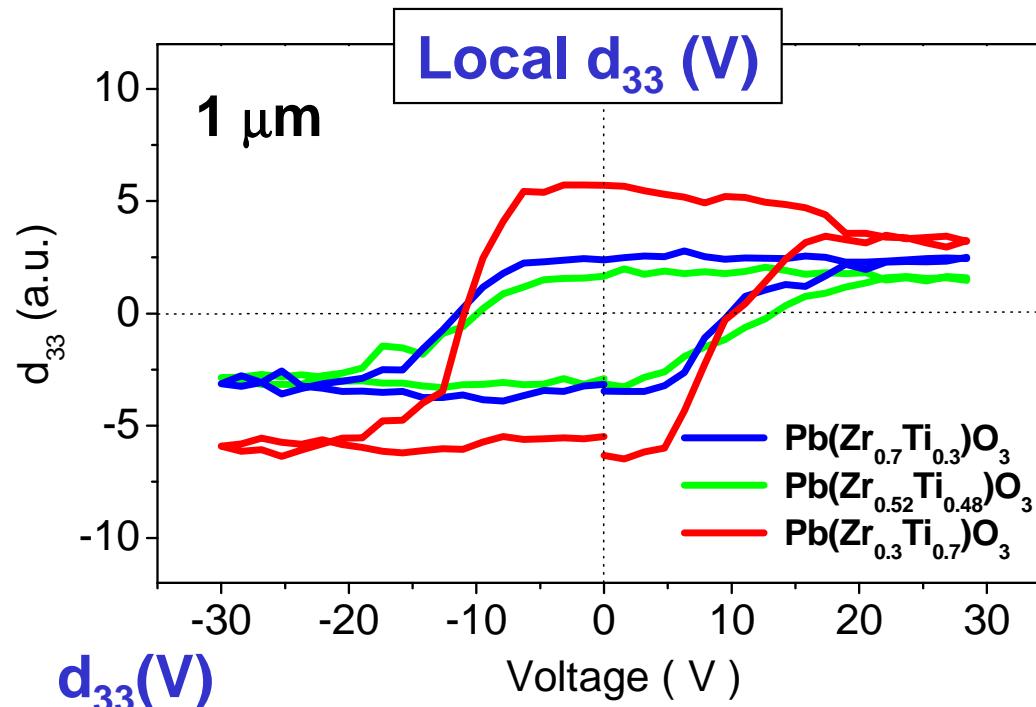


$$V_{\text{out}} = \alpha d_{33} V_{\text{ac}} f[I(V_{\text{dc}})]$$

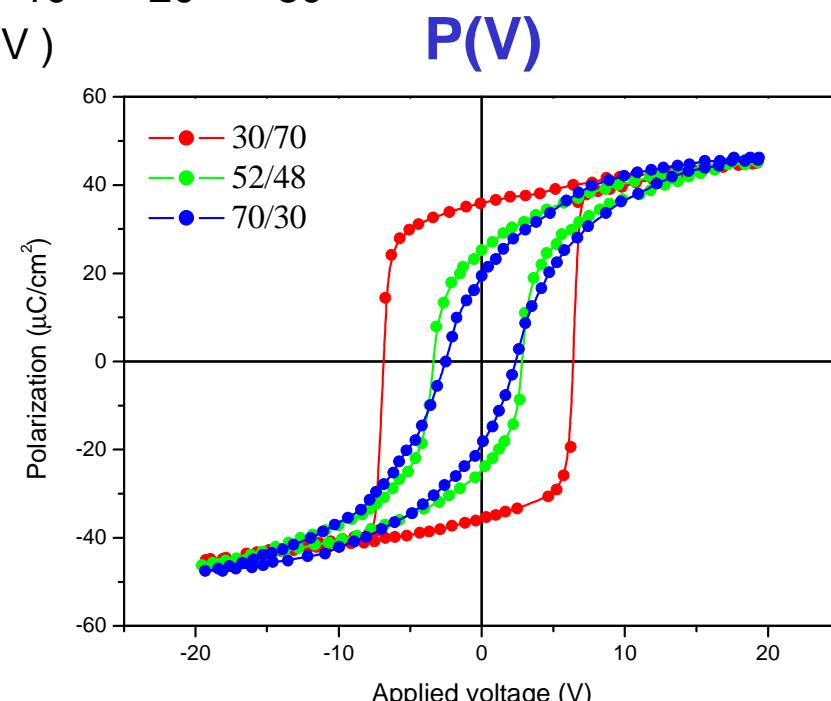
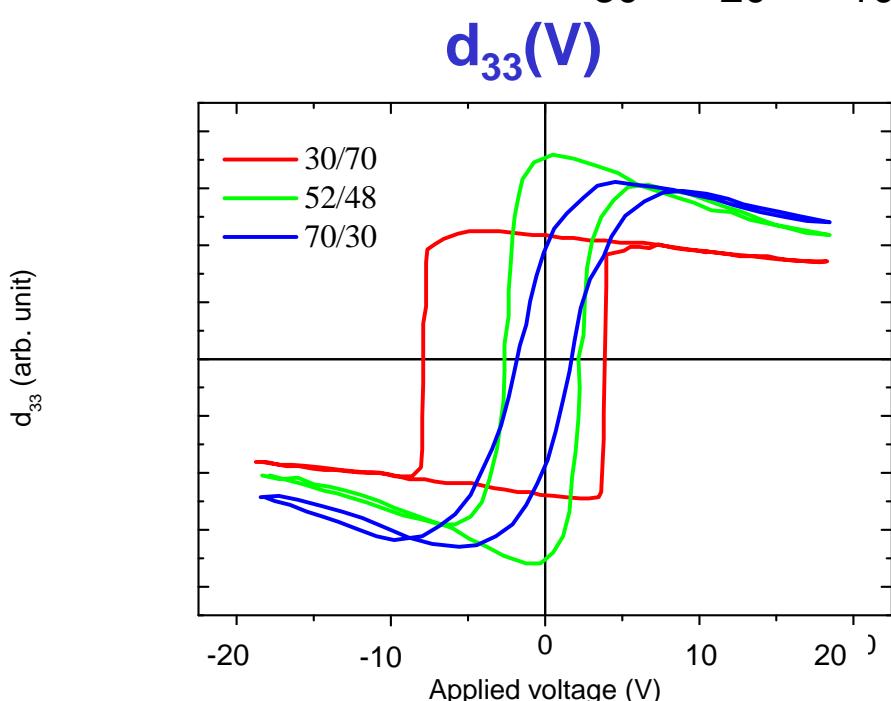
$$r(V_{\text{dc}}) = [S(V_{\text{dc}})/\pi]^{1/2}$$



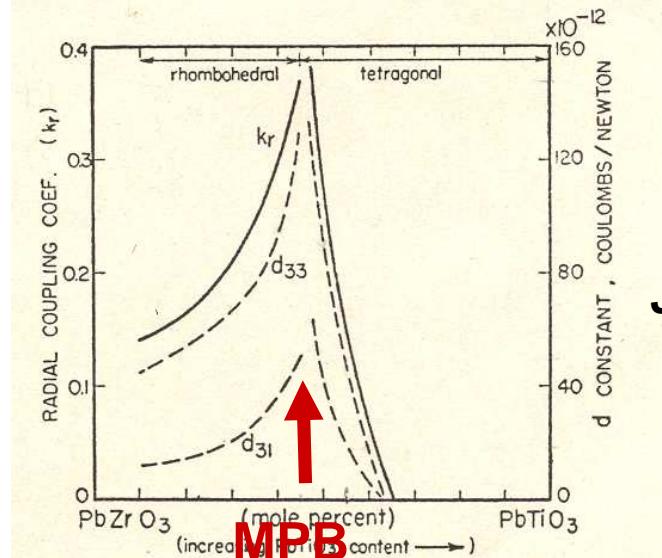
# *Effect of PZT composition*



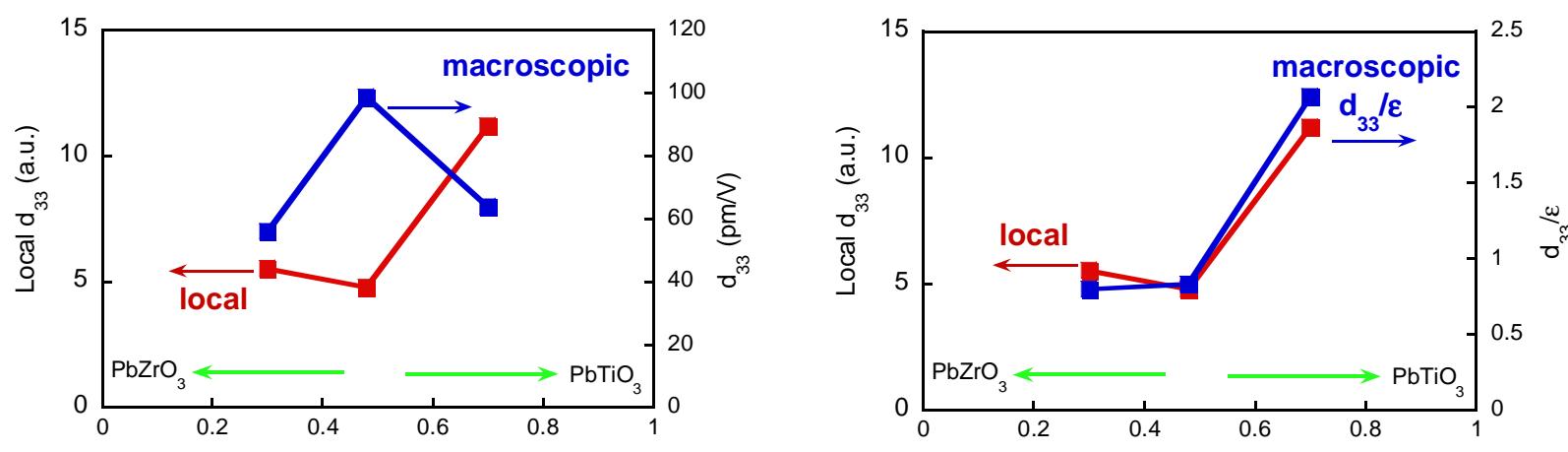
PZT films  
1  $\mu\text{m}$ , sol-gel  
INOSTEK



# Effect of dielectric constant



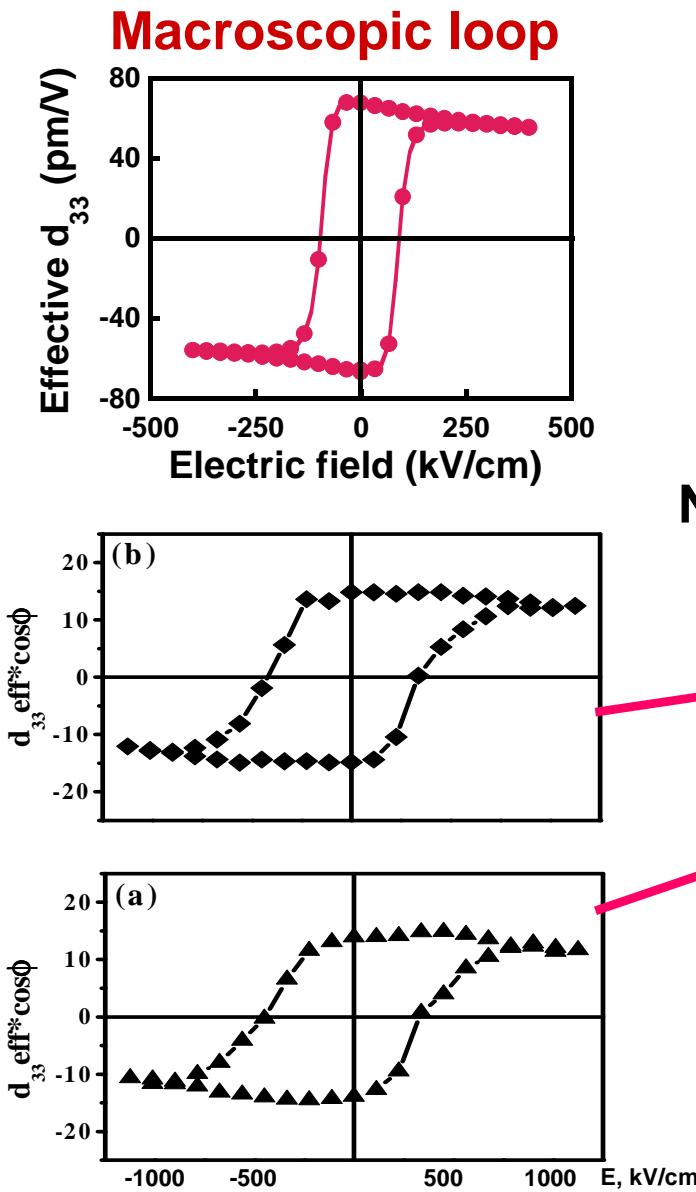
Jaffe et al., J. Appl. Phys. (1954)



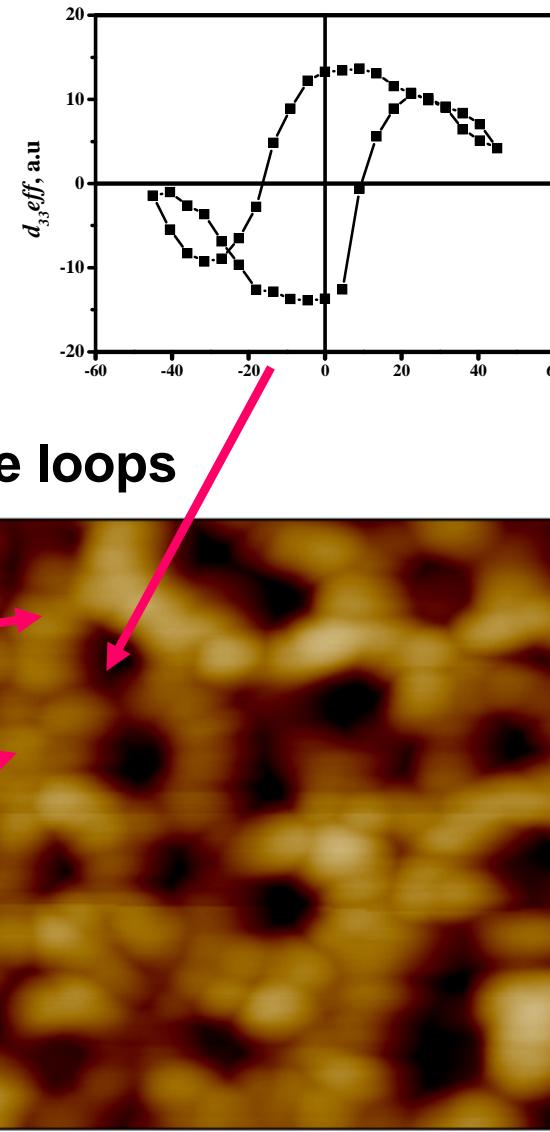
■ Local  $d_{33}$  is inversely proportional to dielectric constant

$$(d_{33})_{eff} \propto \frac{A_{1\omega}}{V_{ac}} = \frac{A_{1\omega}}{1 + \frac{\delta\epsilon}{d\epsilon_\delta}}$$

# *Effect of topography*



**Nanoscale loops**



# Domains in Airvillius-type materials

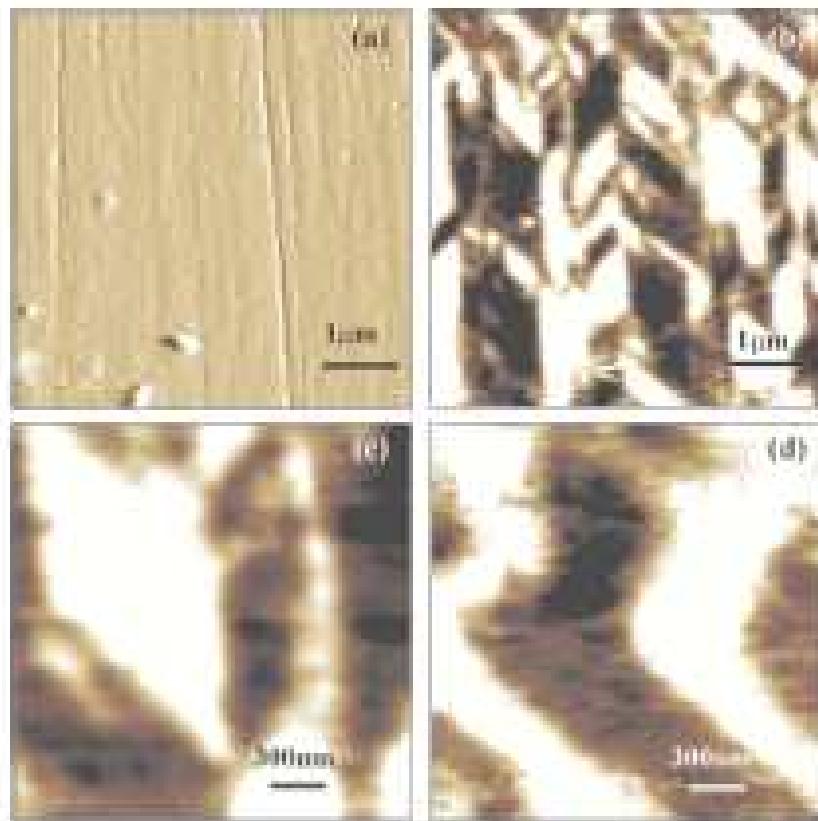
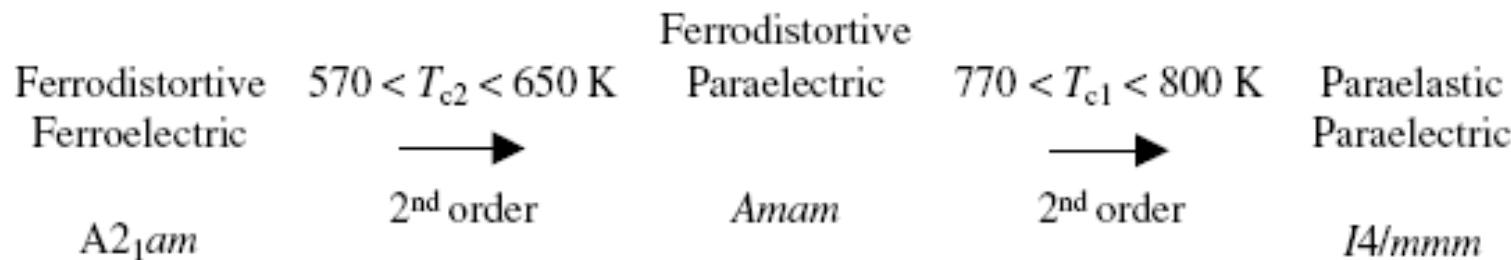


FIG. 2. (Color online) Topography (a) and lateral piezoresponse images (b)–(d) obtained for the (001) face of SBT crystal. The image (d) was acquired after rotating the sample by 90° relative to its position in image (c).

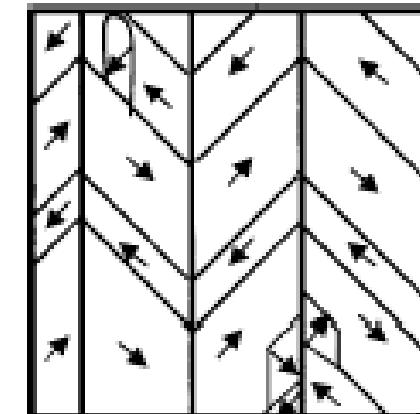
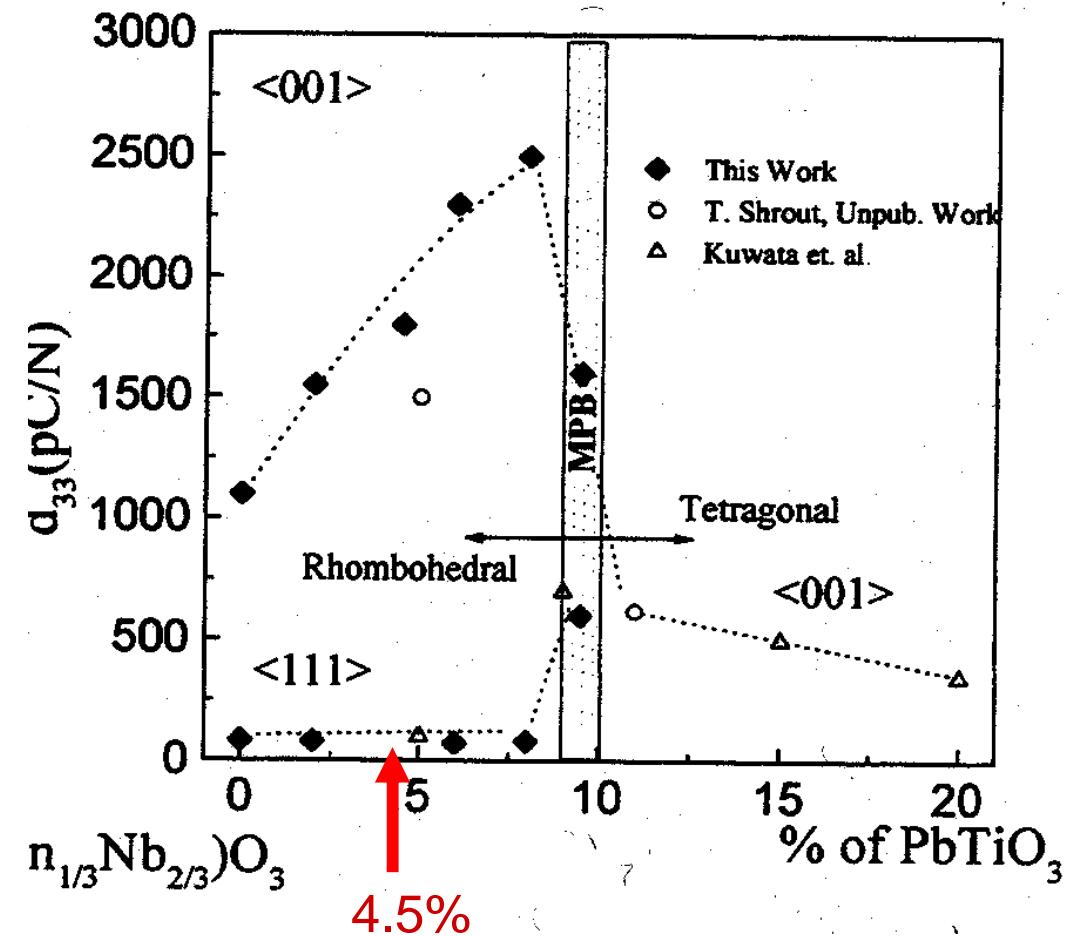
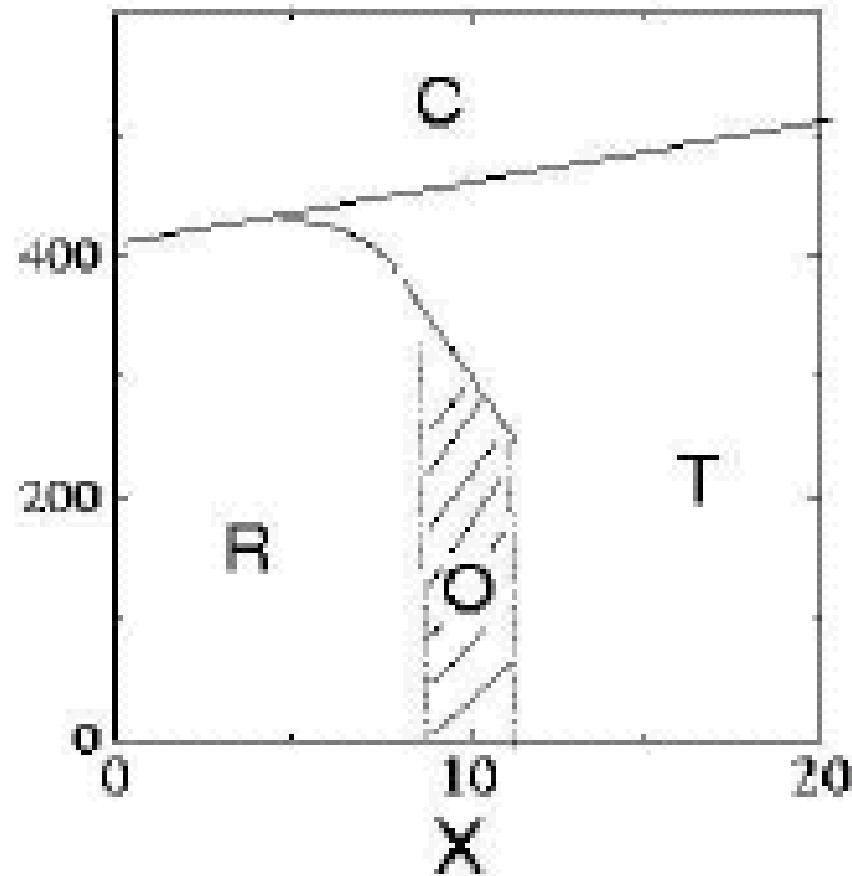


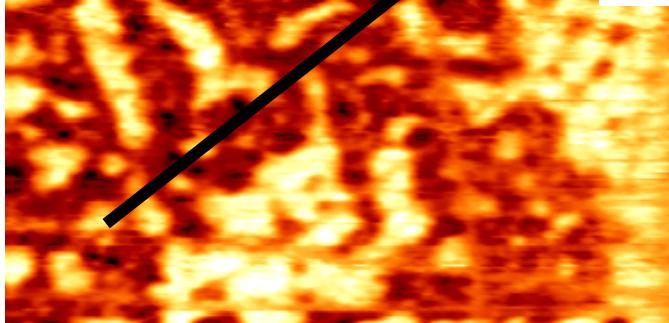
FIG. 3. Reconstructed domain structure of the SBT single crystal. The cross section of domain pattern in the  $a\bar{b}$  plane is depicted. Arrows show the polarization directions.

# *Phase diagram of PZN-PT*

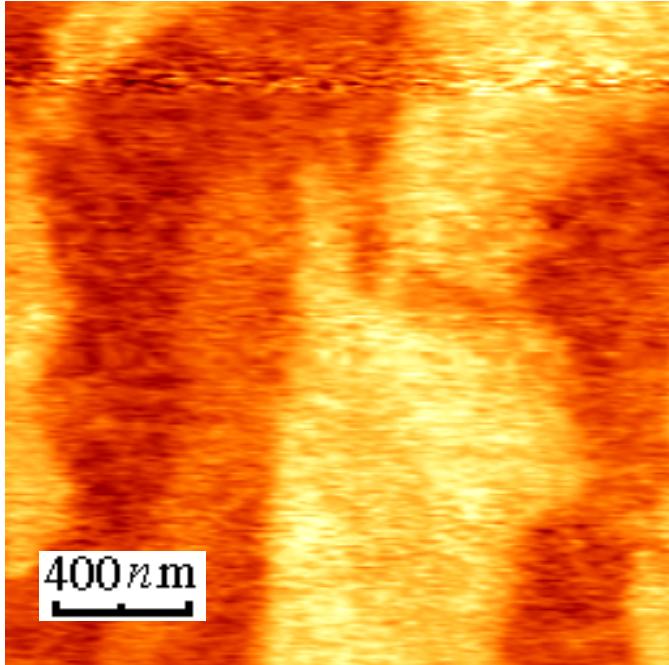
PZN-PT



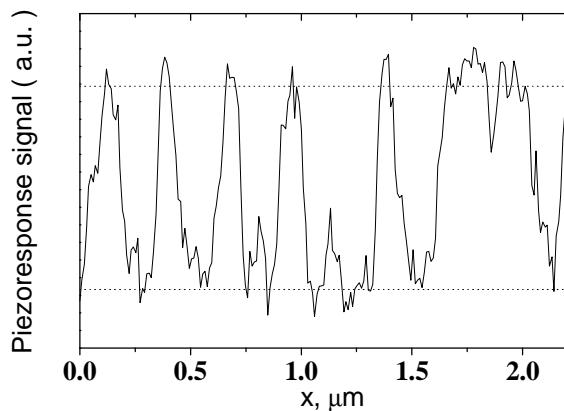
(001)



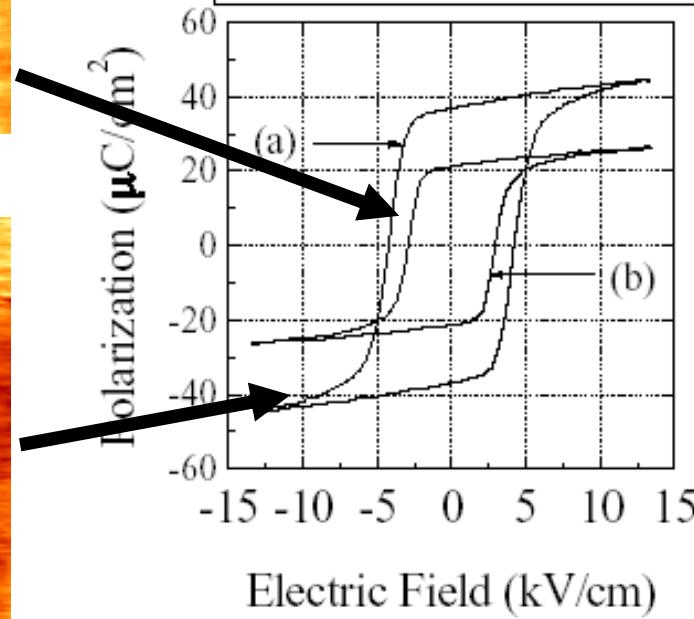
(111)



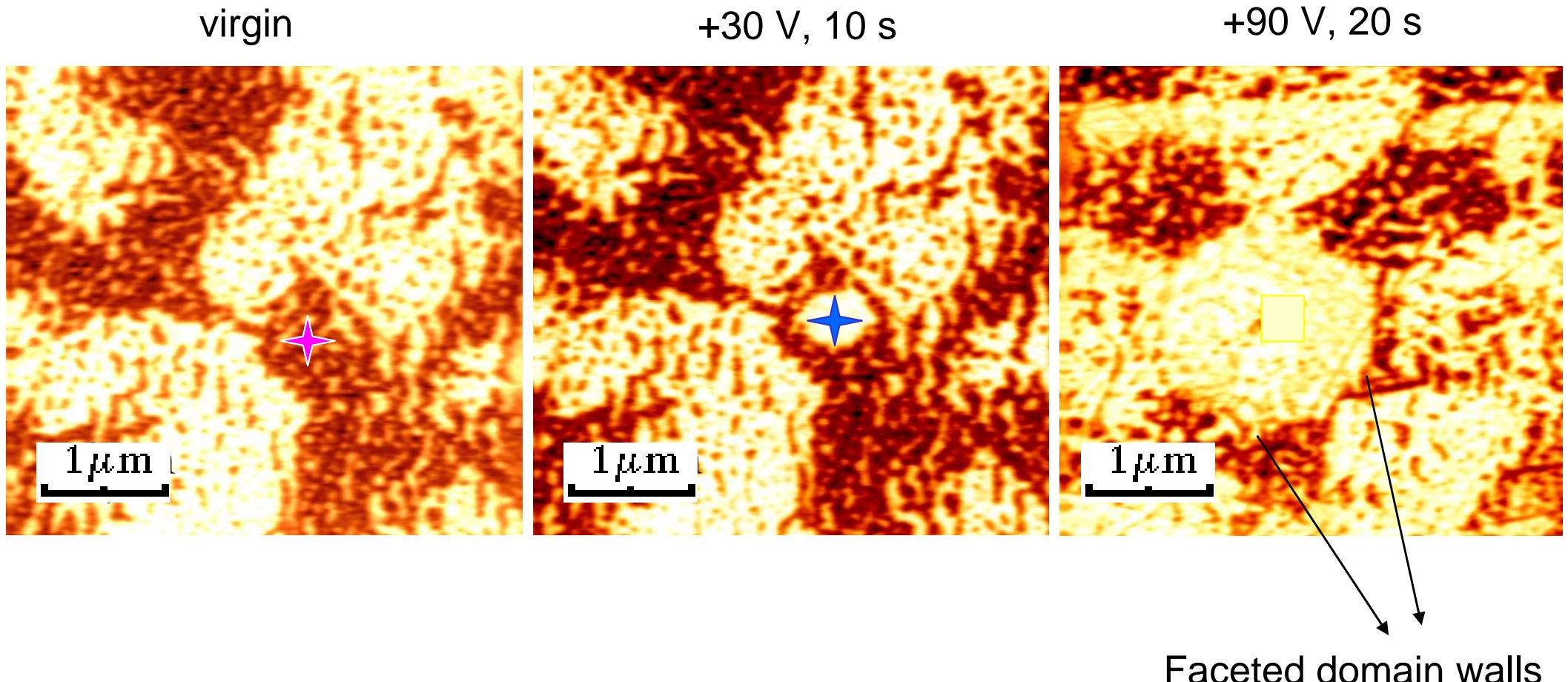
## Polar nanodomains



(a) (111) oriented crystal  
(b) (001) oriented crystal

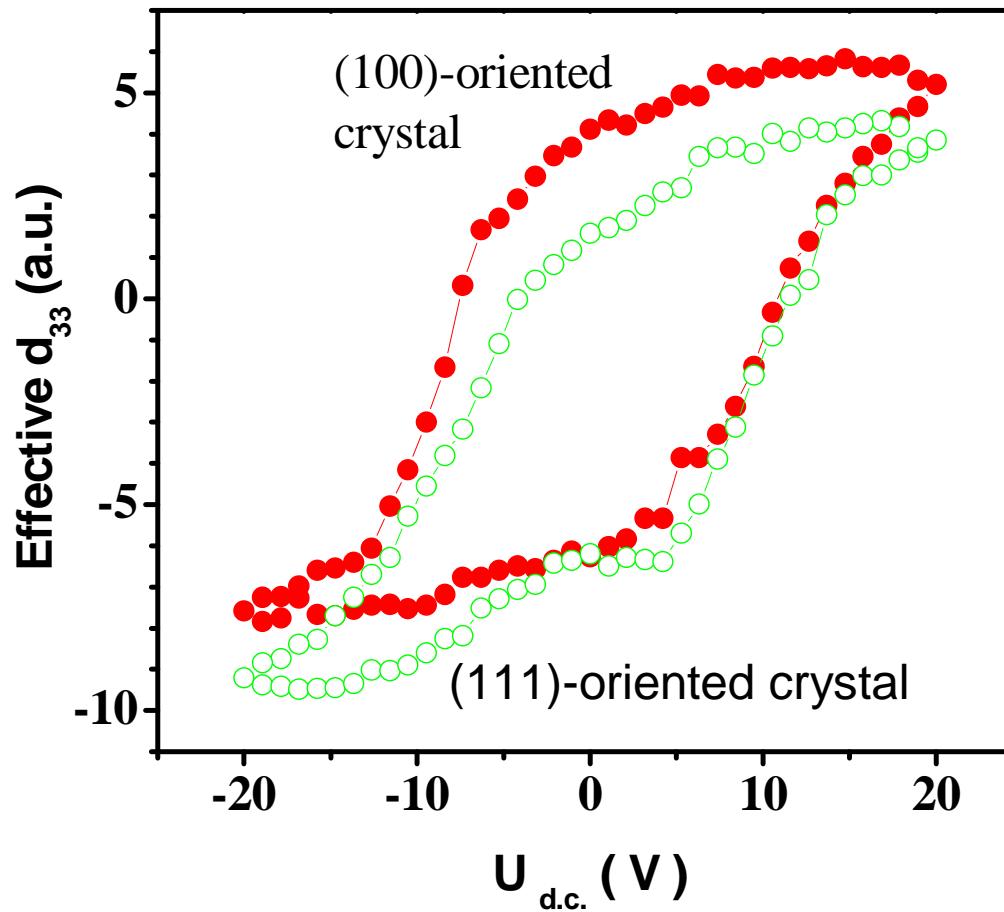


# *Poling by the tip*



- Poling does not remove nanodomains
  - “Normal” micron-sized domains can be written

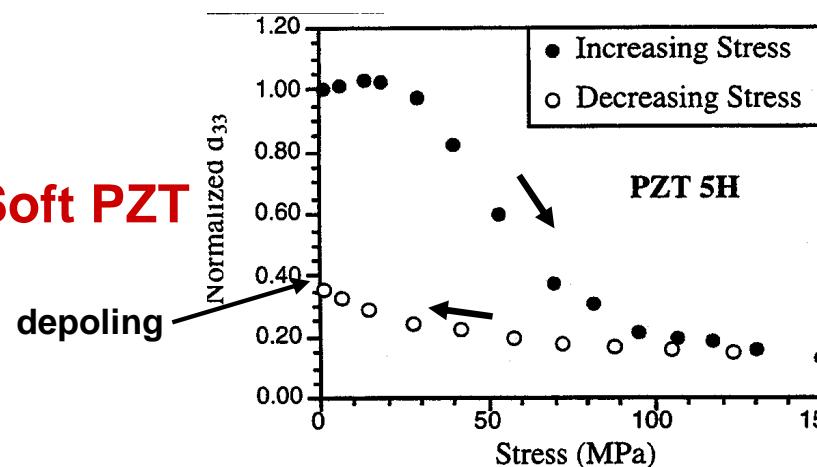
# *Hysteresis comparison*



- Local properties are comparable → giant piezoresponse is related to nanodomains

# Mechanical stress effect on hysteresis

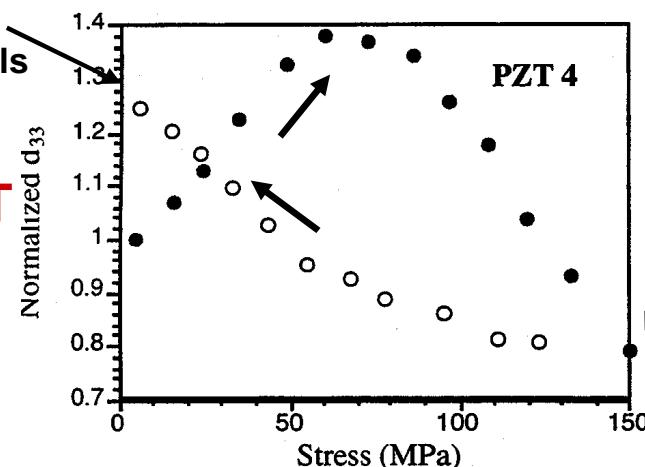
**Soft PZT**



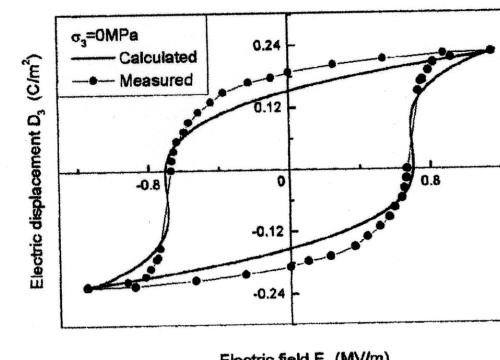
ferroelastic switching

Depinning of domain walls

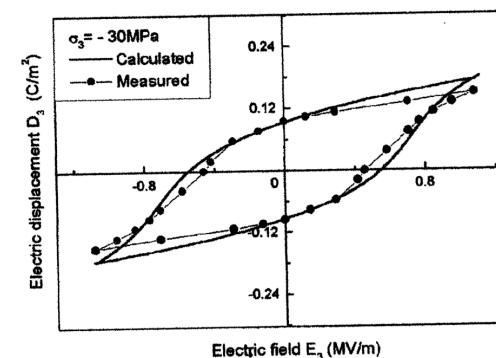
**Hard PZT**



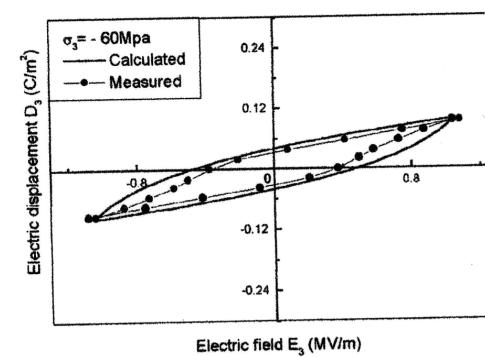
Zang et al, J. Mater. Res. 12, 236 (1997)



-30 MPa



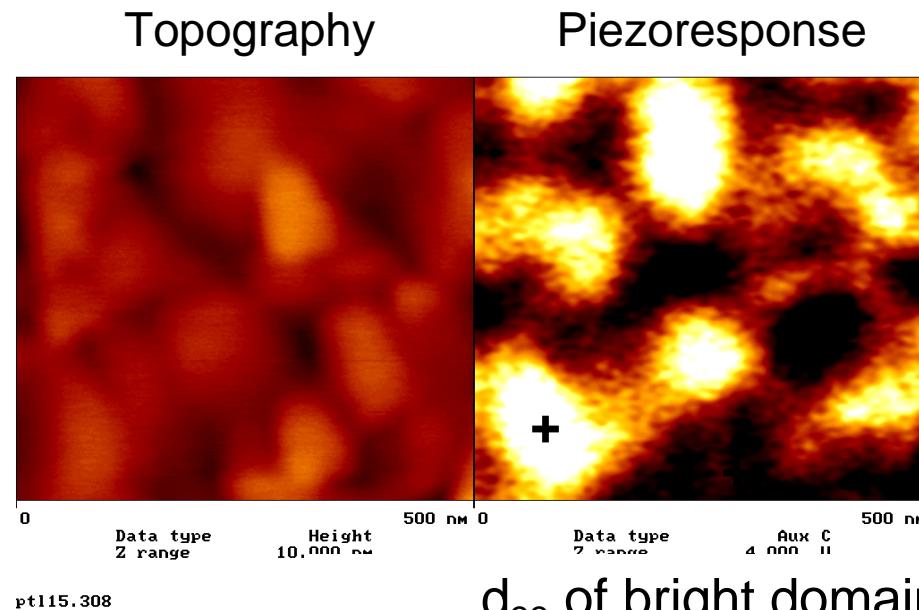
-60 MPa



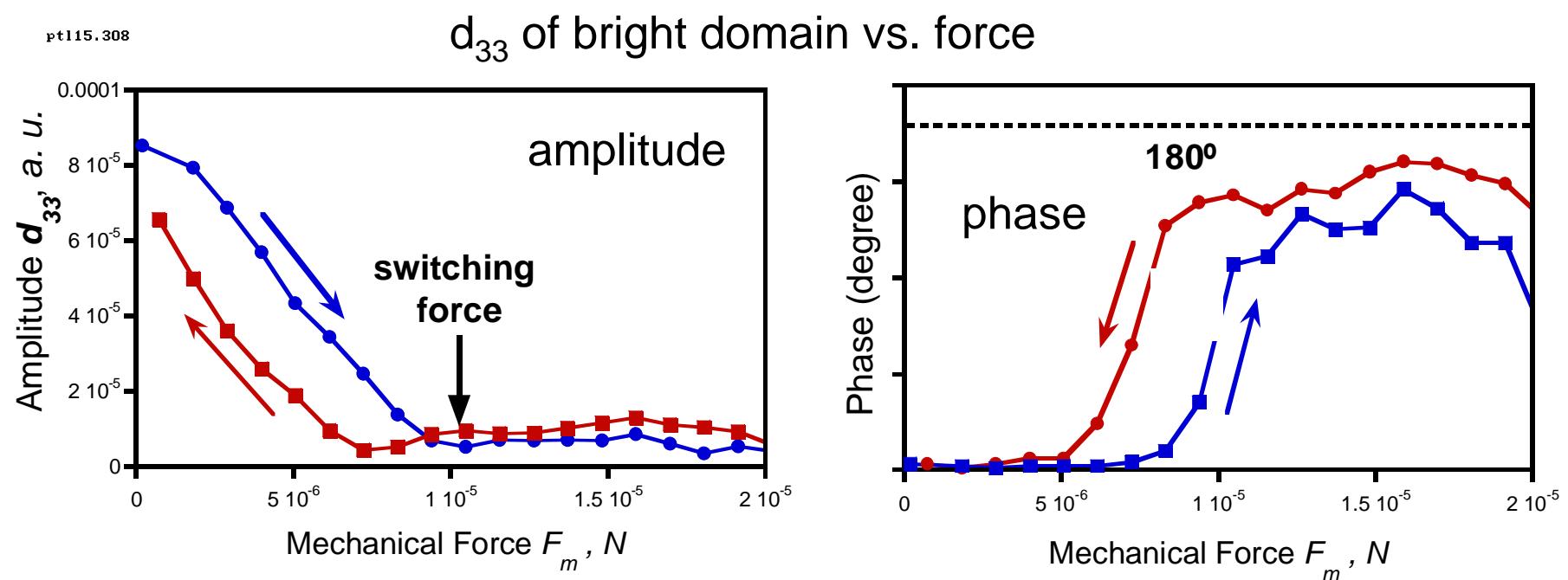
Fang & Li, J. Mater. Sci. 34, 4001 (1999)

- Compressive strain (|| P) results in depoling (soft PZT) through ferroelastic domain wall motion

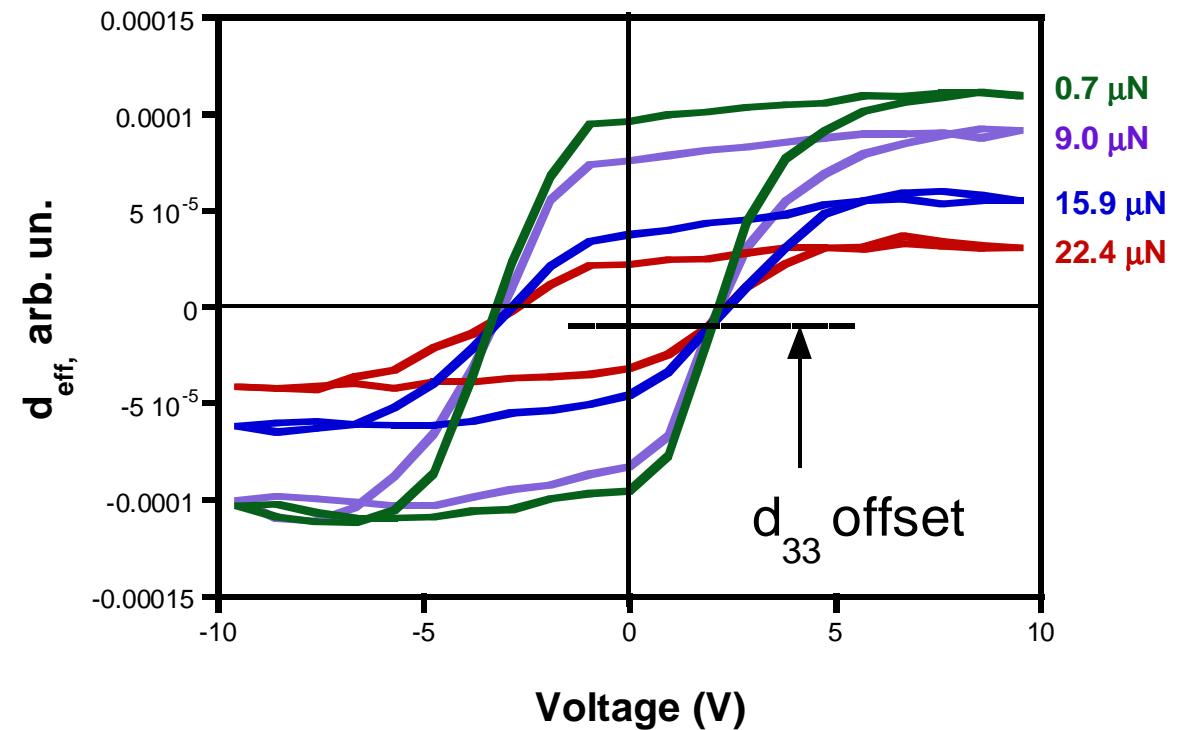
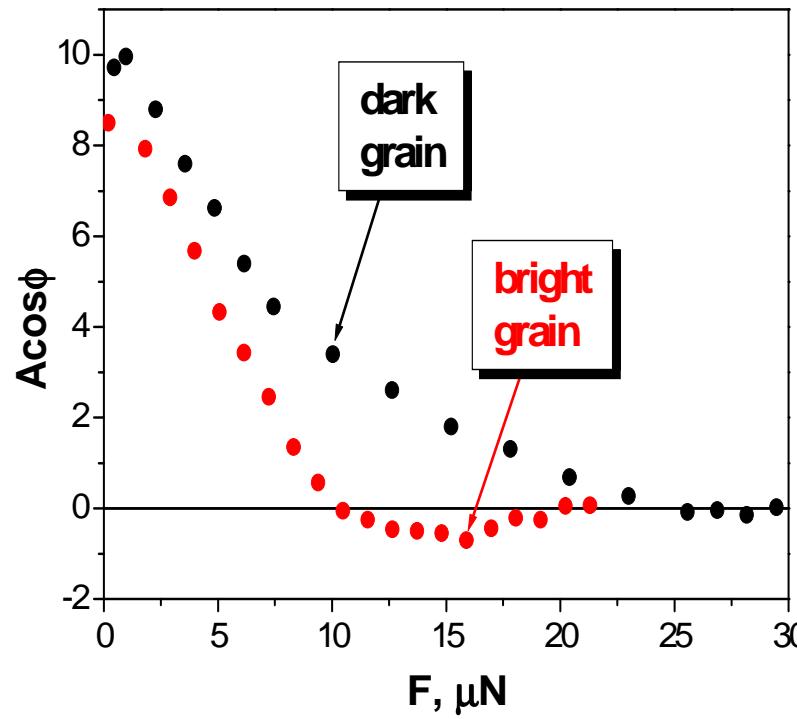
# Stress by PFM: $PbTiO_3$ -La films



- Sol-gel spin-on technique
- Thickness  $\approx 340$  nm
- (001)/(100) preferred orientation
- Grain size 50-150 nm
- High piezoelectric anisotropy

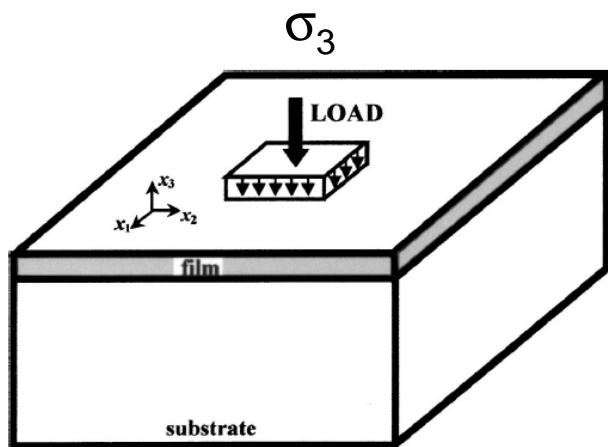


# Hysteresis suppression



- $d_{33}$  suppression depends on the initial contrast
- Hysteresis is suppressed with strong polarization offset

# Non-linear thermodynamic approach



$$S_1 = S_2 = S_m = \frac{b^* - a_0}{a_0} = \text{const}$$

- misfit strain

$$S_6 = 0$$

Renormalization leads to

$$\left\{ \begin{array}{l} S_m^\sigma = S_m - (s_{11} + 2s_{12})\sigma_3, \\ T_\sigma = T - 2\varepsilon_0 C(Q_{11} + 2Q_{12})\sigma_3. \end{array} \right.$$

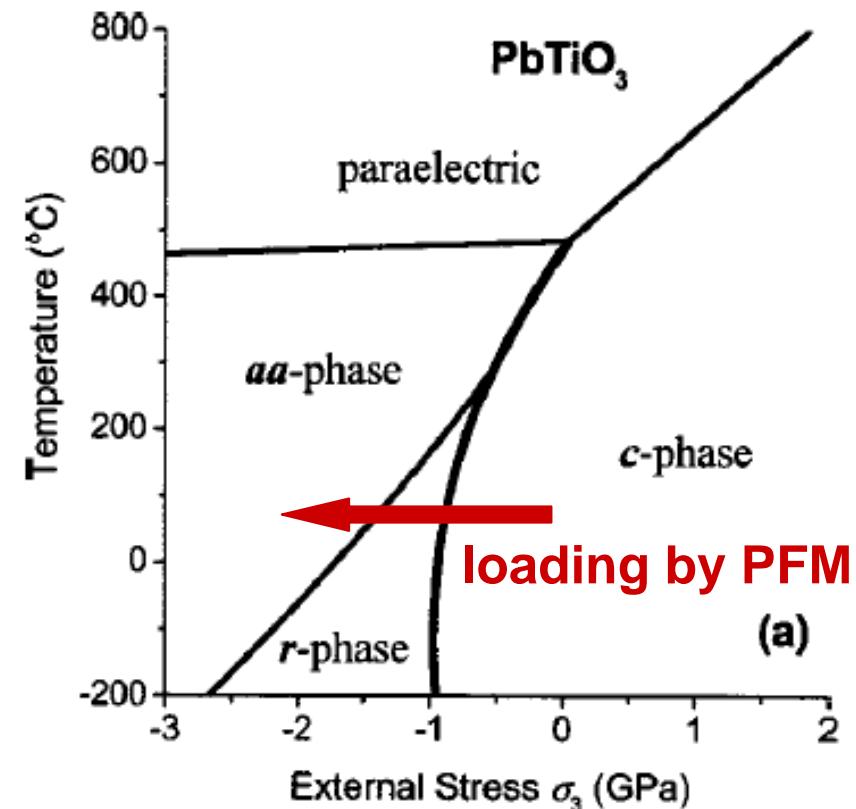
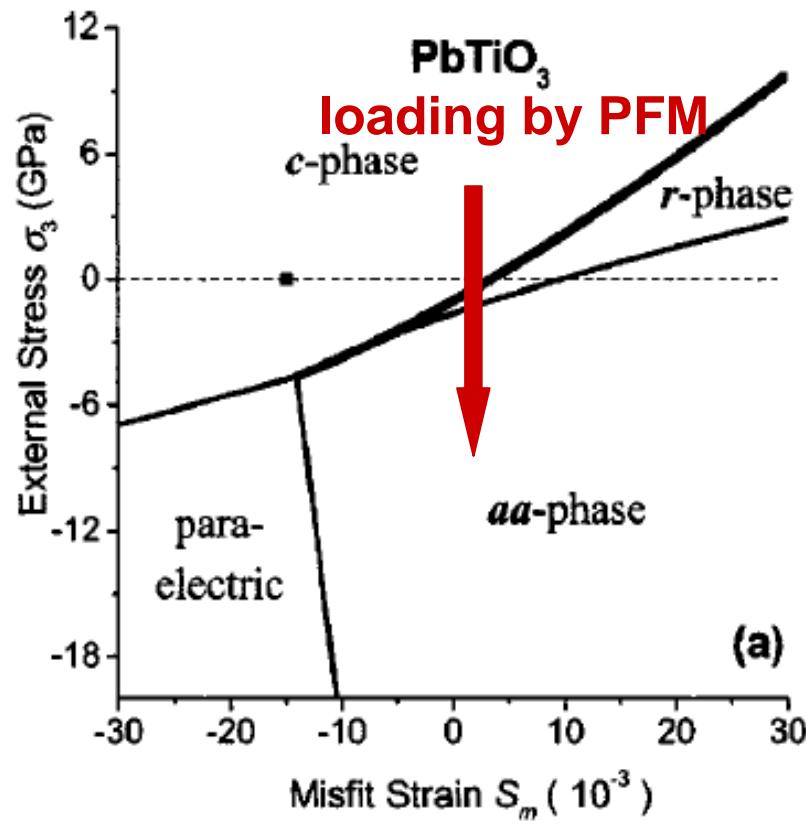
## Modified thermodynamic potential

$$\begin{aligned} \tilde{G} &= G + S_1\sigma_1 + S_2\sigma_2 + S_6\sigma_6 \\ \tilde{G} &= \frac{(S_m - s_{12}\sigma_3)^2}{s_{11} + s_{12}} - \frac{1}{2}s_{11}\sigma_3^2 + a_1^*(P_1^2 + P_2^2) + a_3^*P_3^2 \\ &\quad + a_{11}^*(P_1^4 + P_2^4) + a_{12}^*P_1^2P_2^2 + a_{13}^*(P_1^2 + P_2^2)P_3^2 + a_{33}^*P_3^4 \\ &\quad + a_{111}(P_1^6 + P_2^6 + P_3^6) + a_{112}[P_1^4(P_2^2 + P_3^2) + P_2^4(P_1^2 + P_3^2) \\ &\quad + P_3^4(P_1^2 + P_2^2)] + a_{123}P_1^2P_2^2P_3^2, \end{aligned}$$

where  $a_1^*, a_{ij}^*, a_{ijk}^*$  are dielectric stiffnesses at constant stress

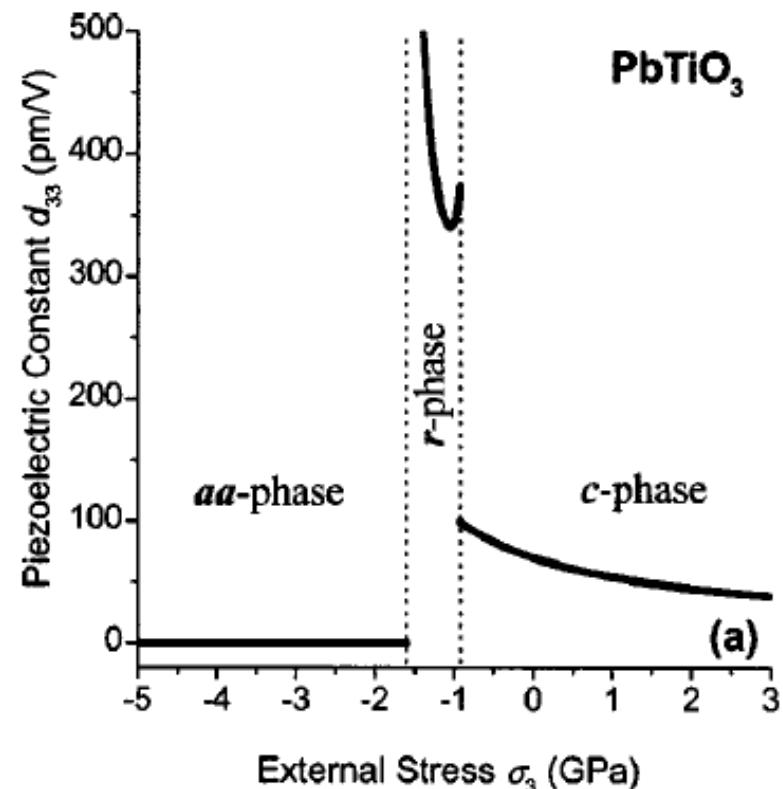
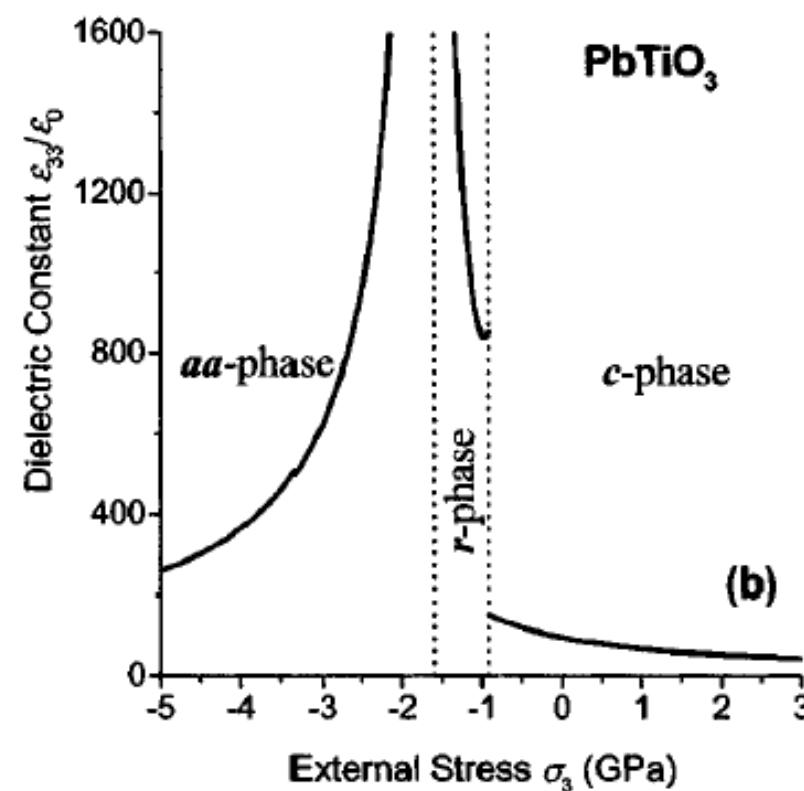
$$\text{e. g. } a_1^* = a_1 - \frac{Q_{11} + Q_{12}}{s_{11} + s_{12}} S_m + \frac{Q_{11}s_{12} - Q_{12}s_{11}}{s_{11} + s_{12}} \sigma_3$$

# Phase diagrams under stress



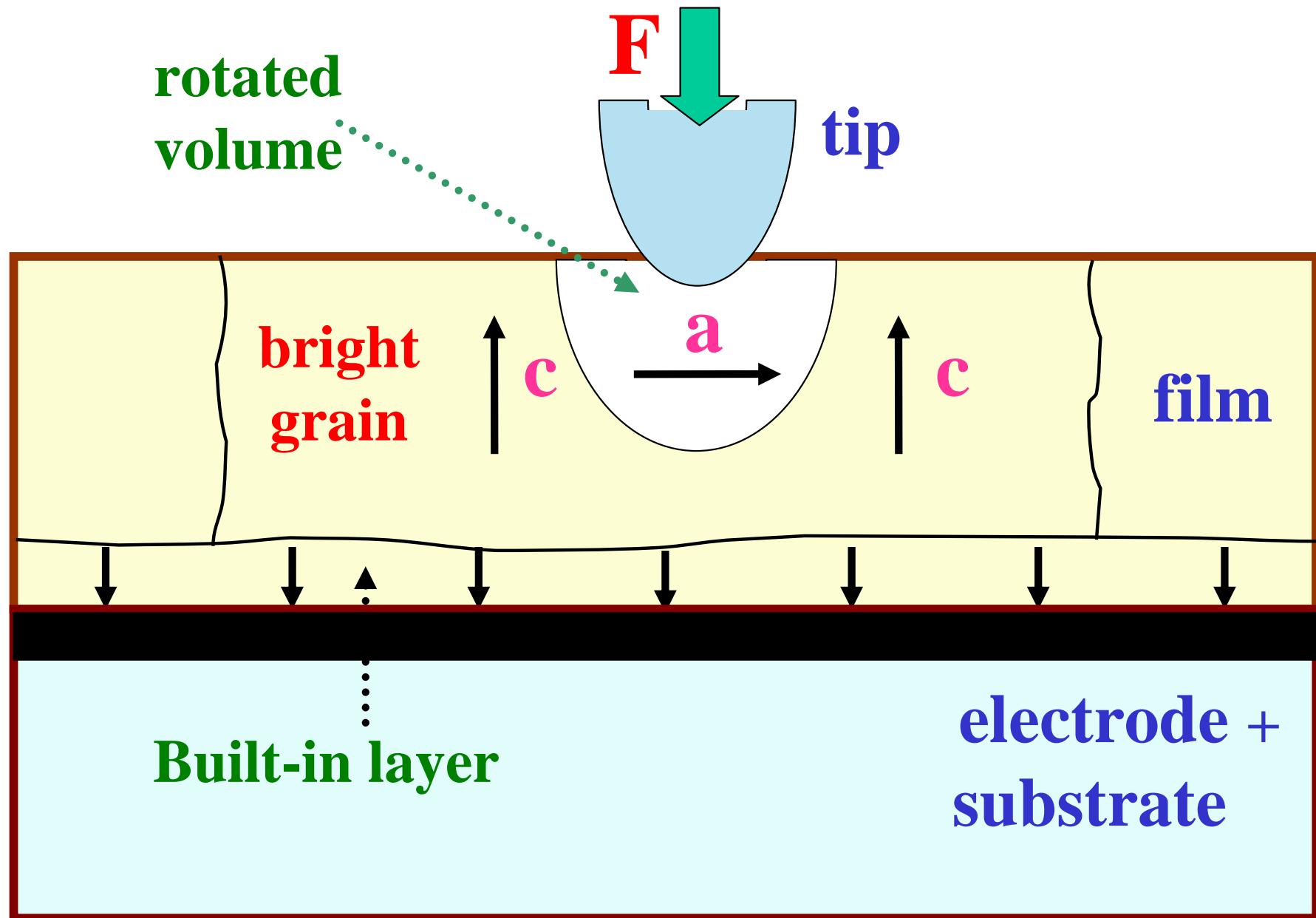
- Phase transition into *r*- and *aa*-phases is possible under high compressive stress

# *Local parameters under stress*

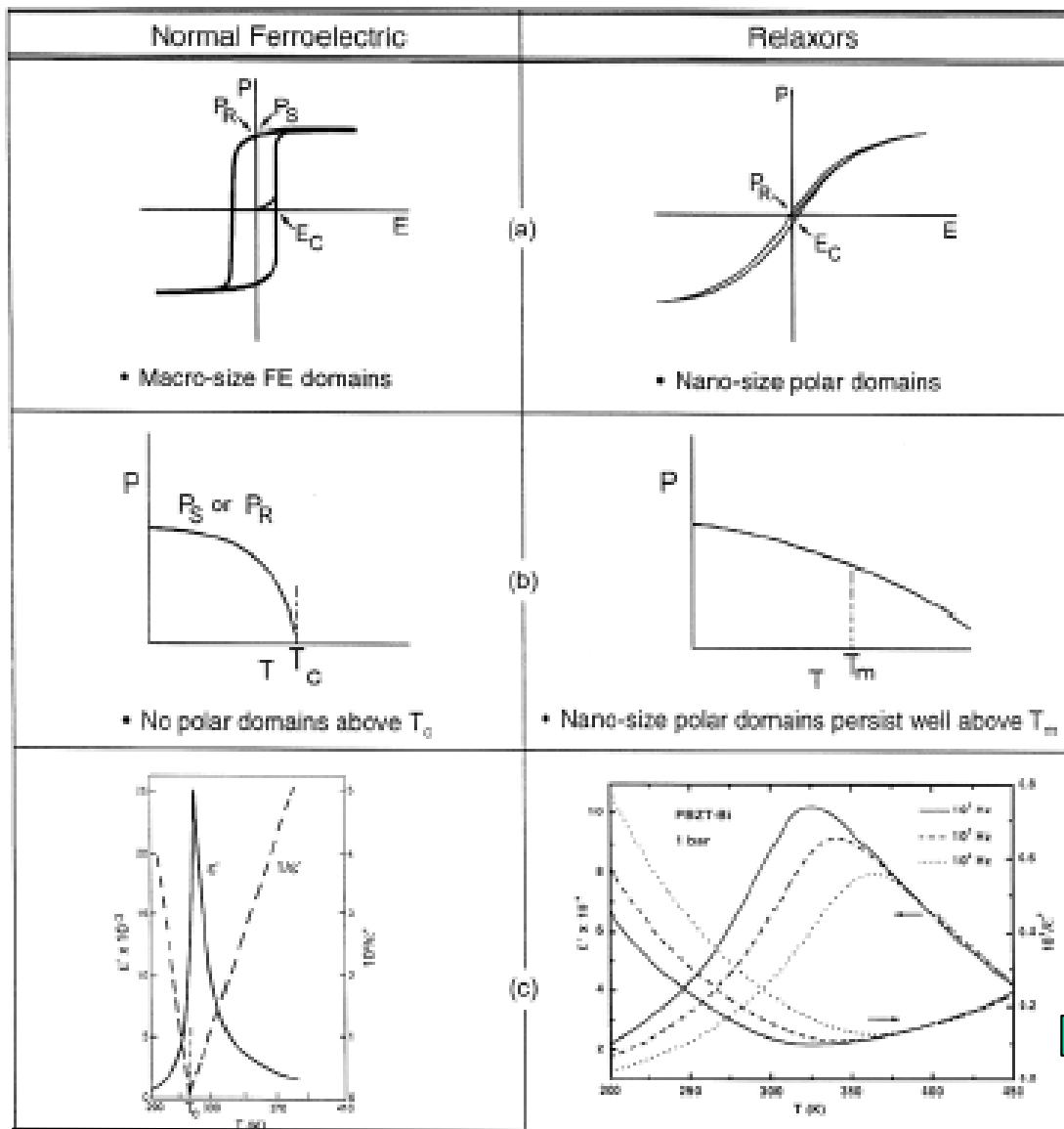


- Proper stress design of microdevices may enhance materials properties near stress-induced phase transitions

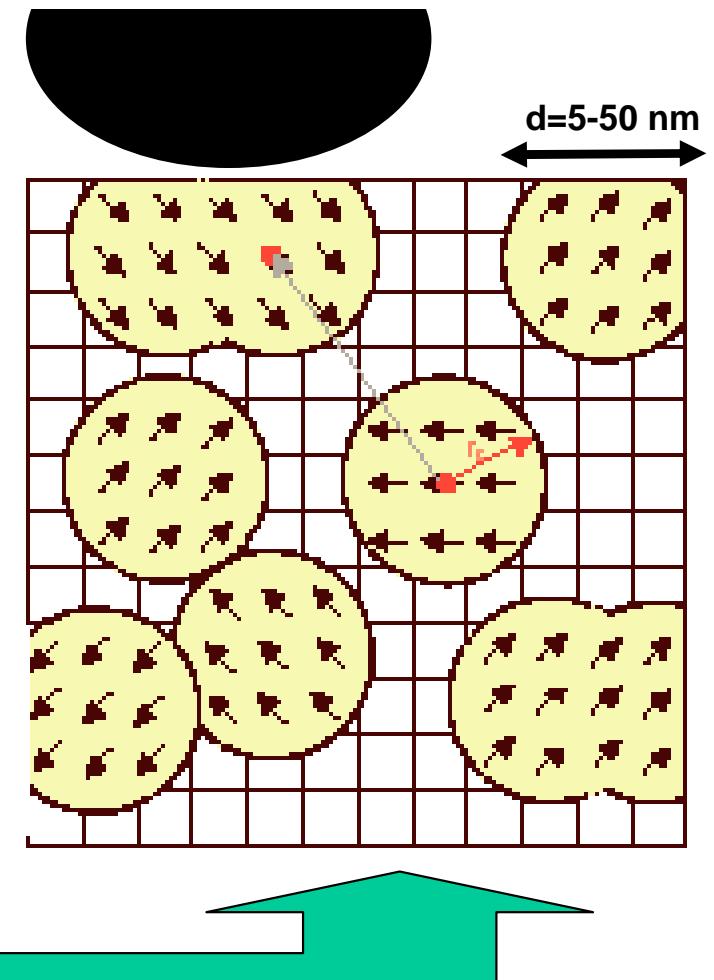
# *Tentative explanation of asymmetry*



# Mesoscopic disorder in relaxors



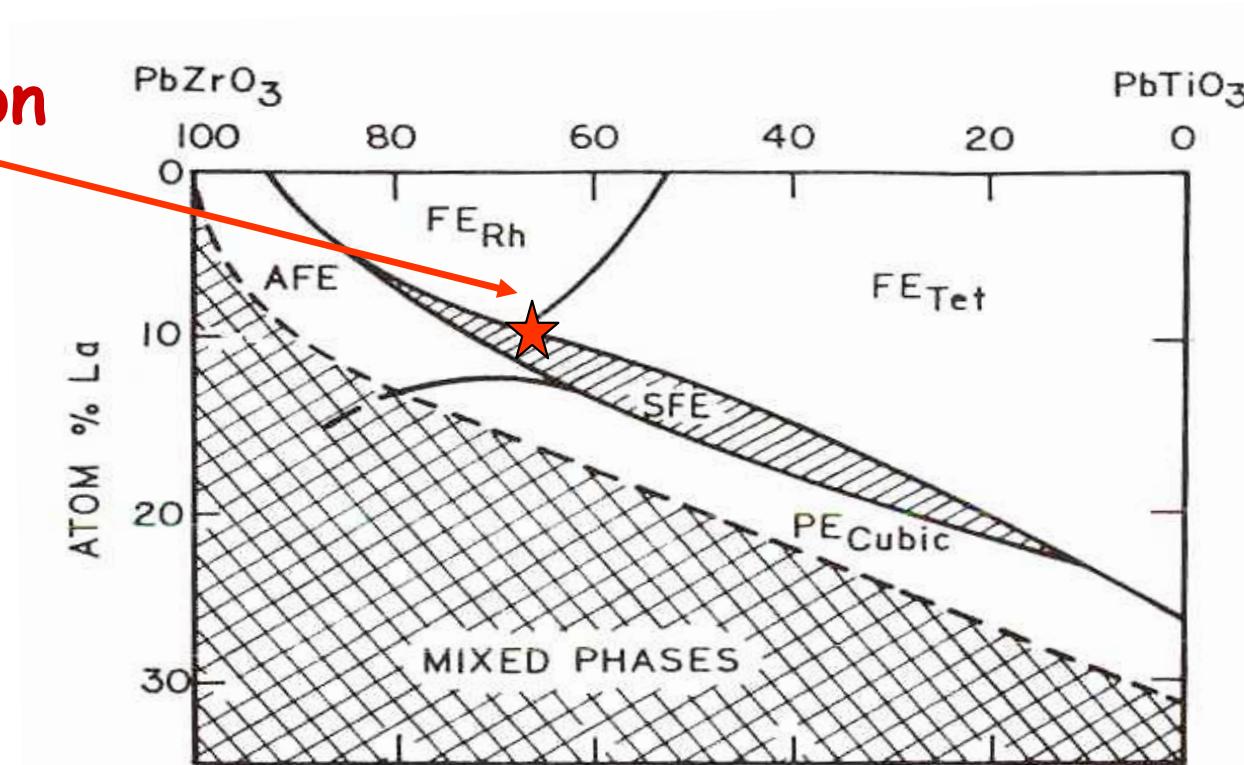
AFM tip



Polar clusters  
(nanodomains)

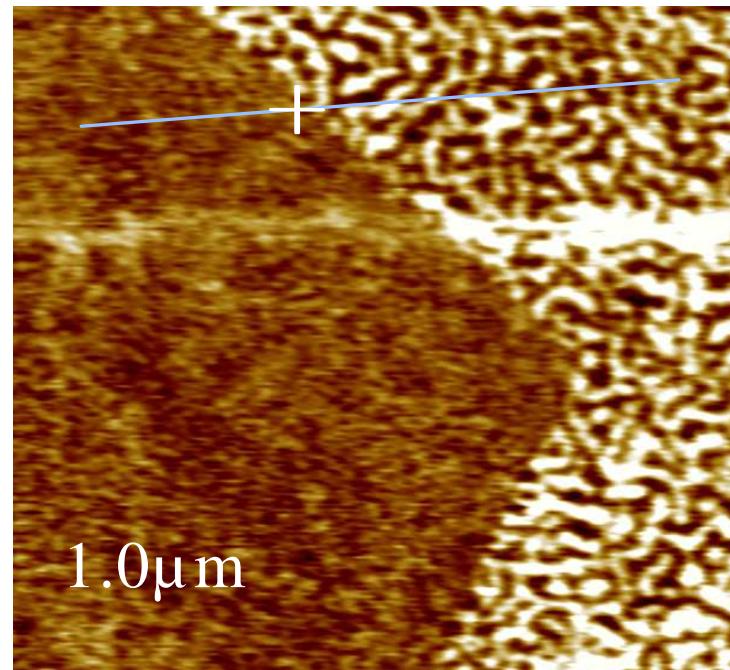
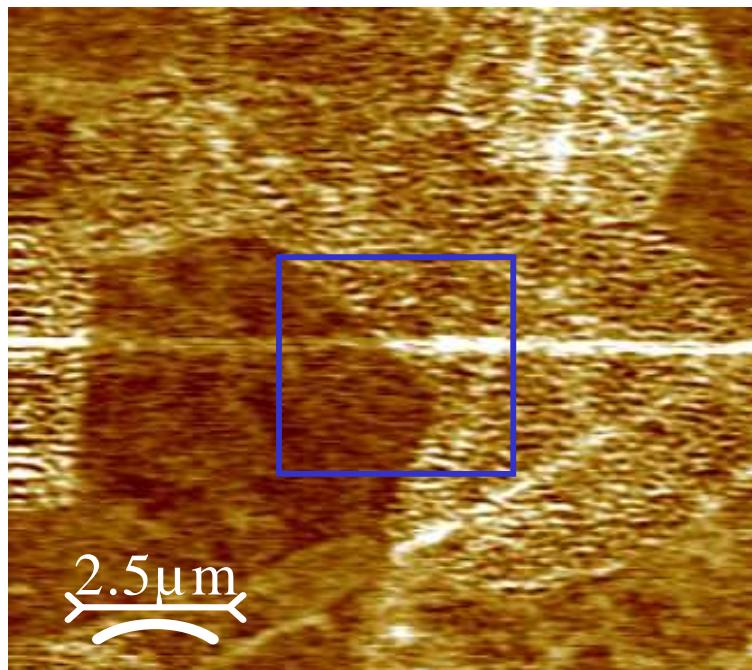
# *Phase diagram of $(Pb,La)(Zr,Ti)O_3$*

studied  
composition

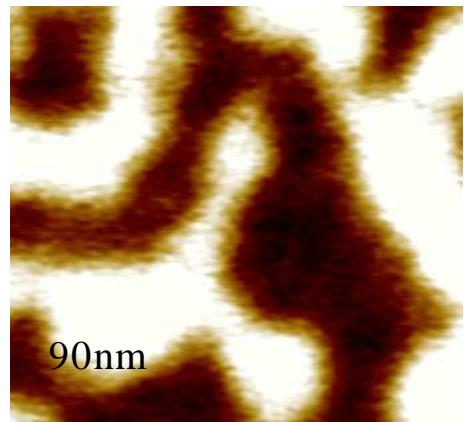


- ❖ Relaxor “ergodic” cubic phase at room temperature
- ❖ No piezoelectric activity at large scale

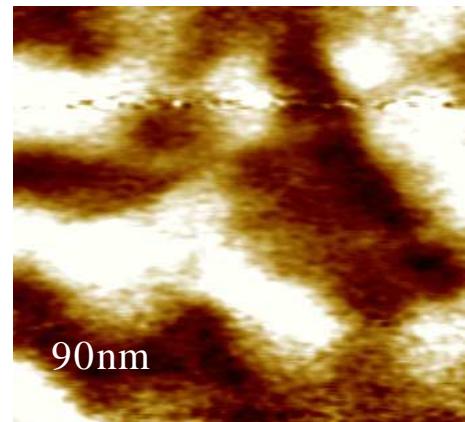
# *Domains in virgin ceramics*



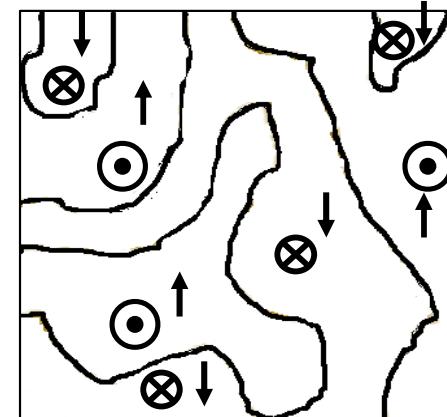
vertical piezoresponse



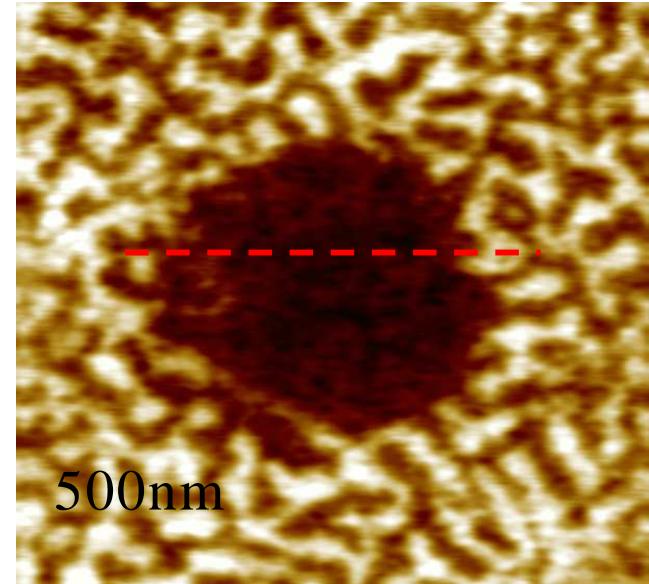
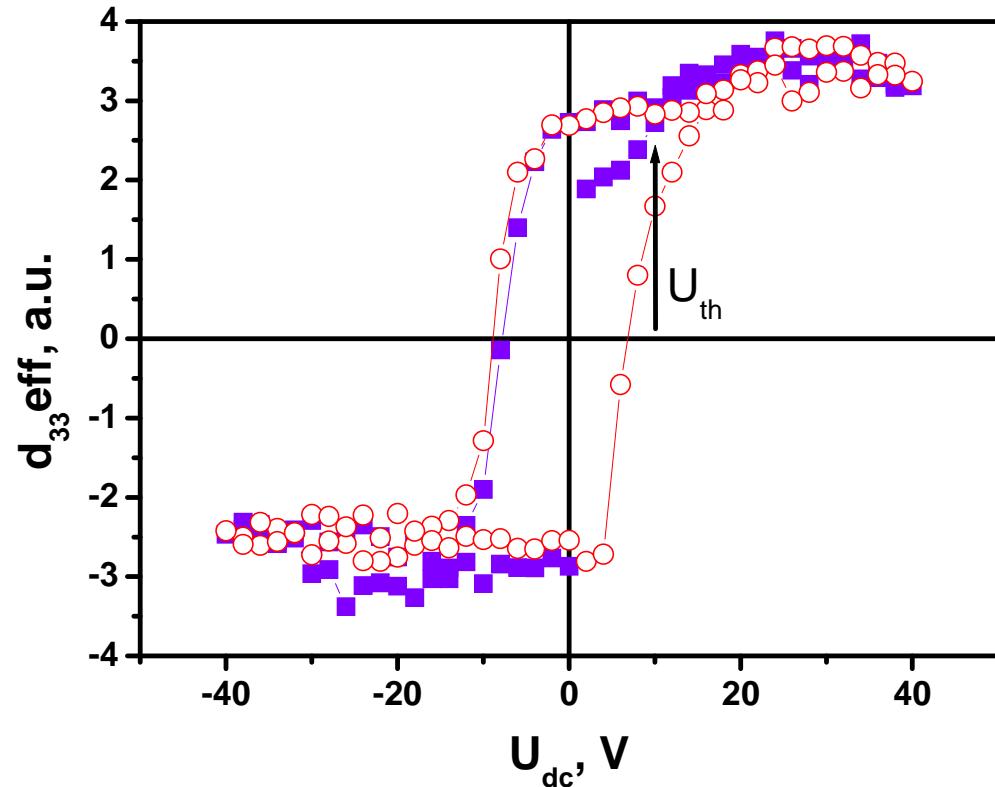
lateral piezoresponse



reconstructed domains

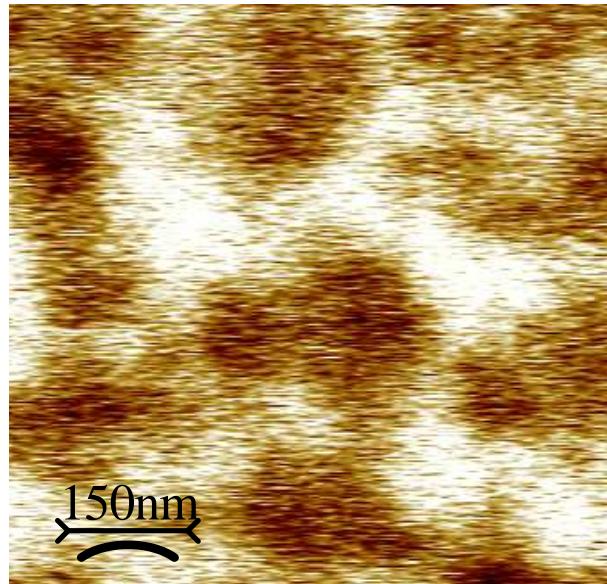


# *Poling and tip-induced phase transition*

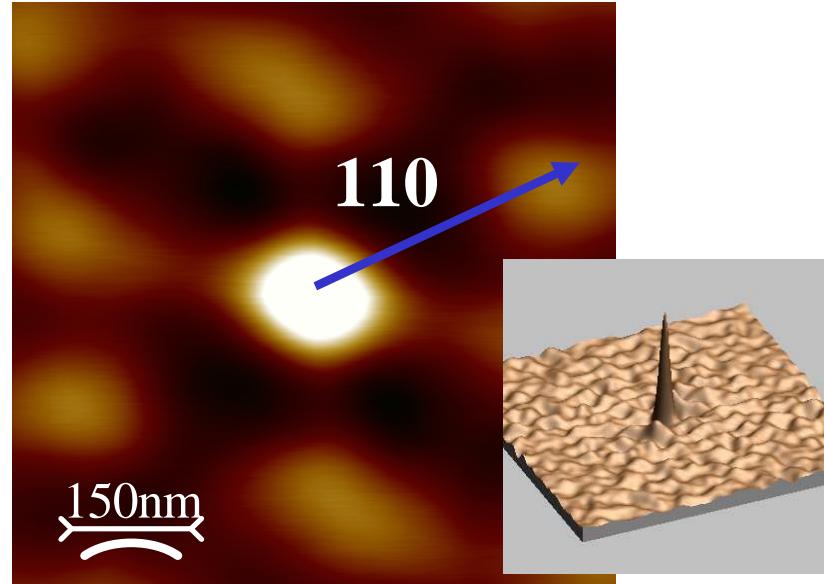


- ❖ Threshold-like increase of piezoresponse signifies bias-induced phase transition
- ❖ Further hysteresis is due to evolution of ferroelectric phase

# Polar nanoclusters in PMN-10%PT

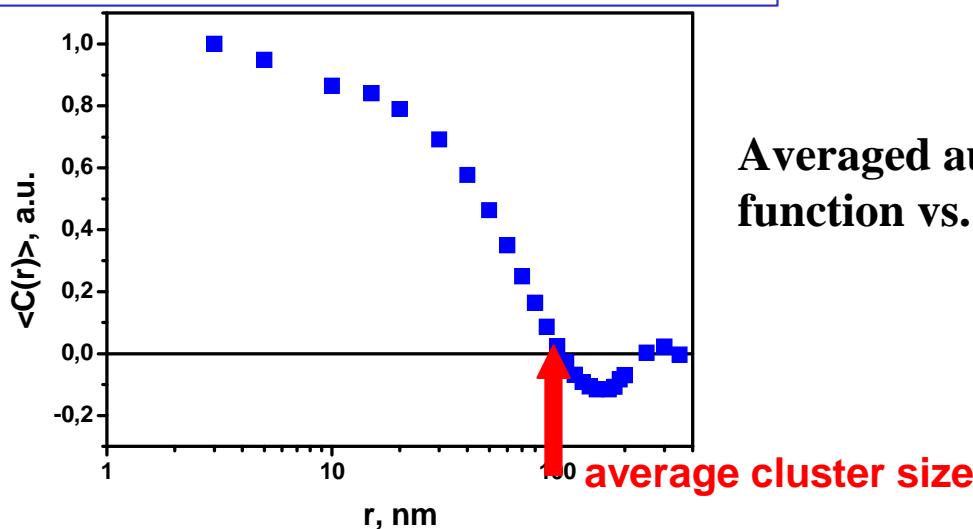


piezoresponse

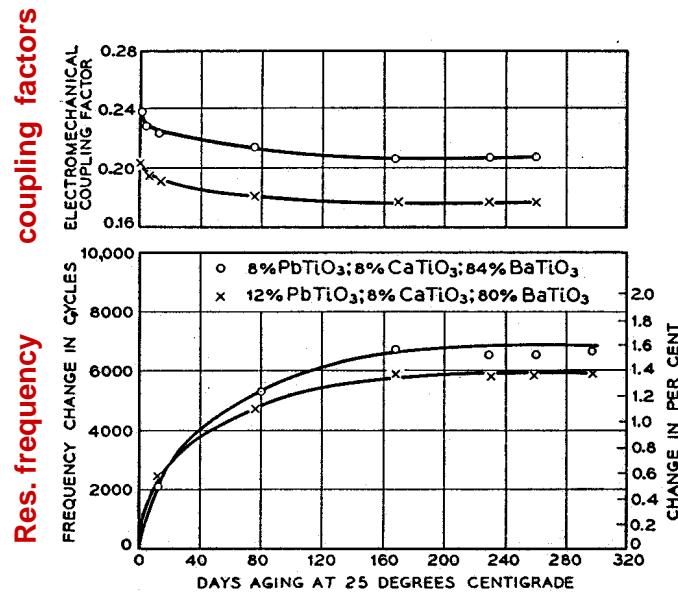


autocorrelation image of  $d_{33}$

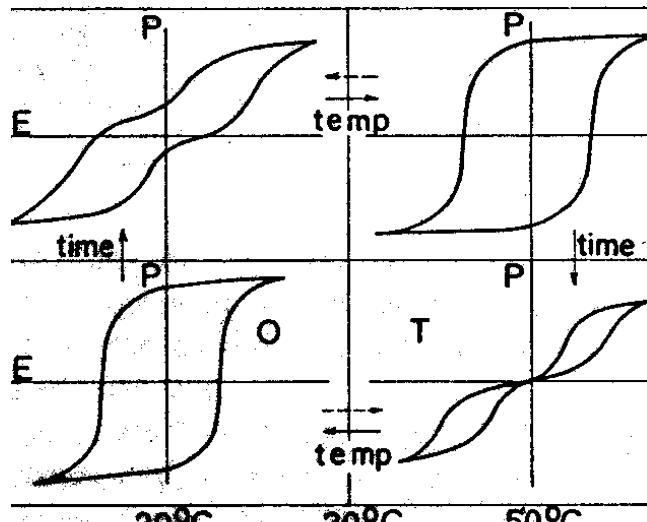
$$C(r_1, r_2) = \sum_{x,y} D(x, y)D(x + r_1, y + r_2)$$



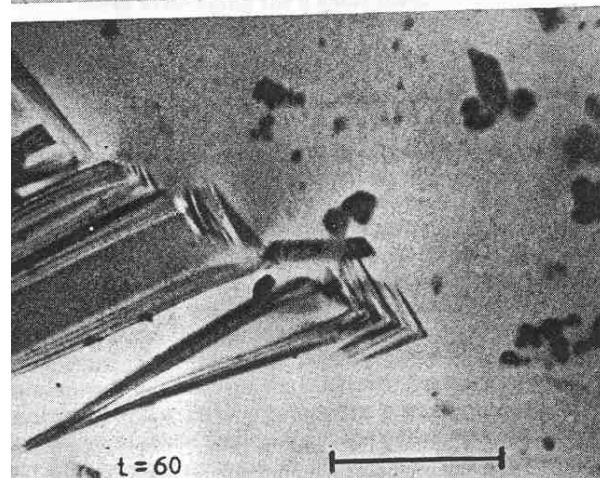
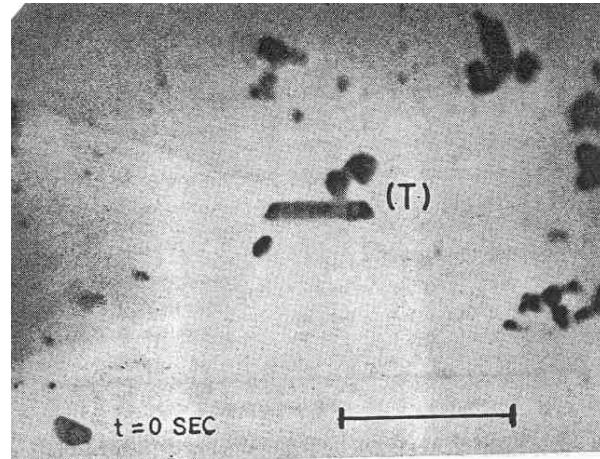
# Aging in ferroelectrics



Mason (1954)

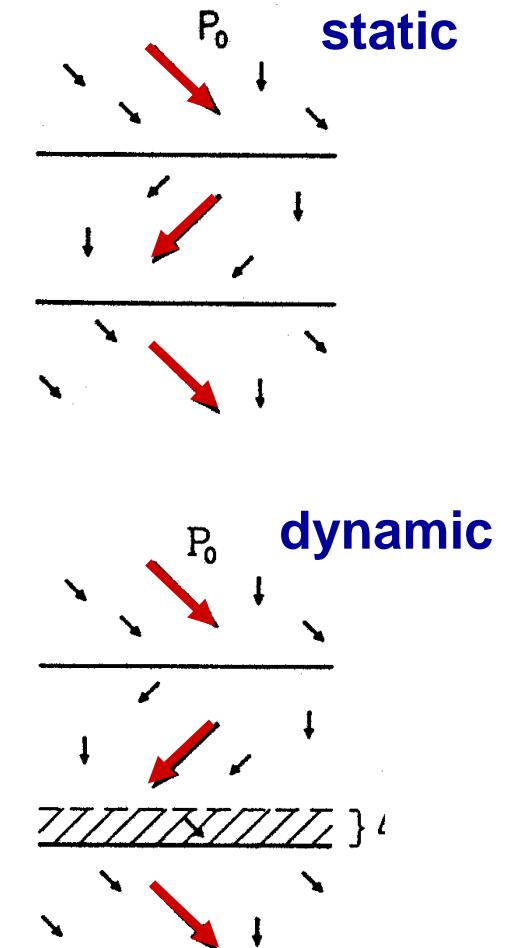


Jonker (1972)



Bradt & Ansell (1968)

Formation of new  
90° domain walls

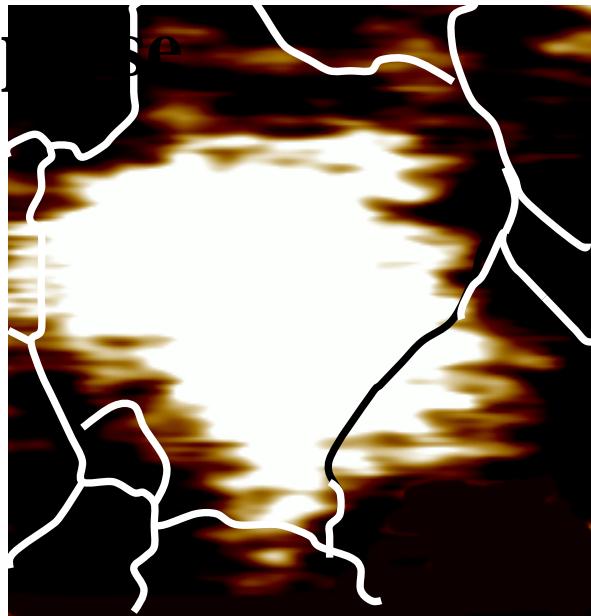


Robels & Arlt (1992)

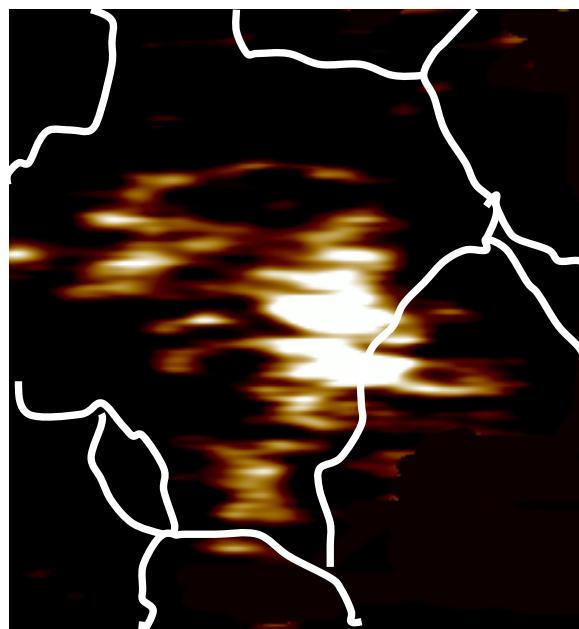
Stabilization of  
existing walls by  
polar defects

# *Piezorelaxation at the nanoscale*

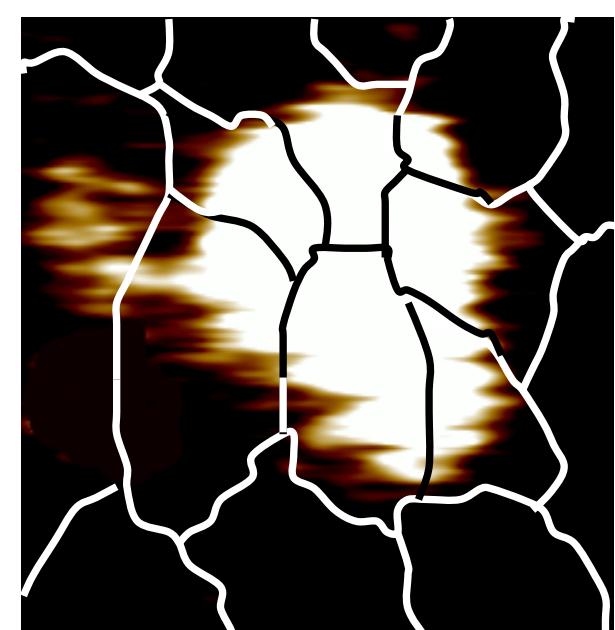
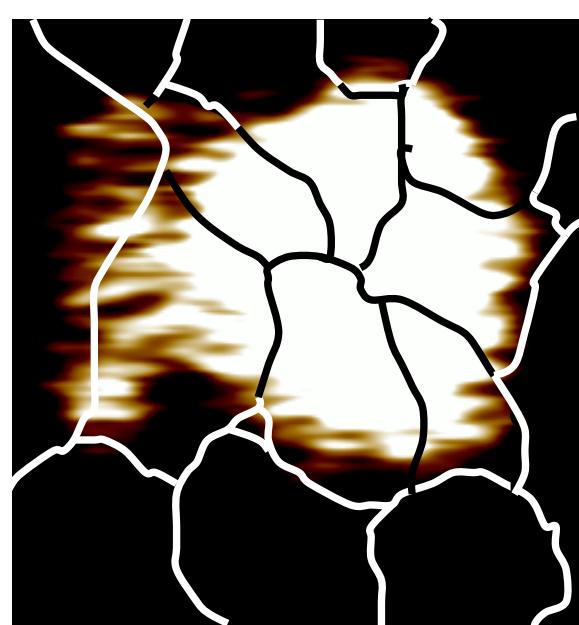
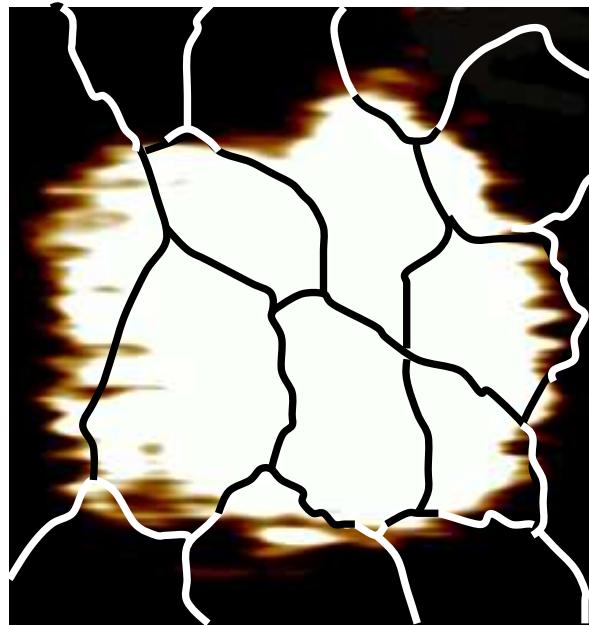
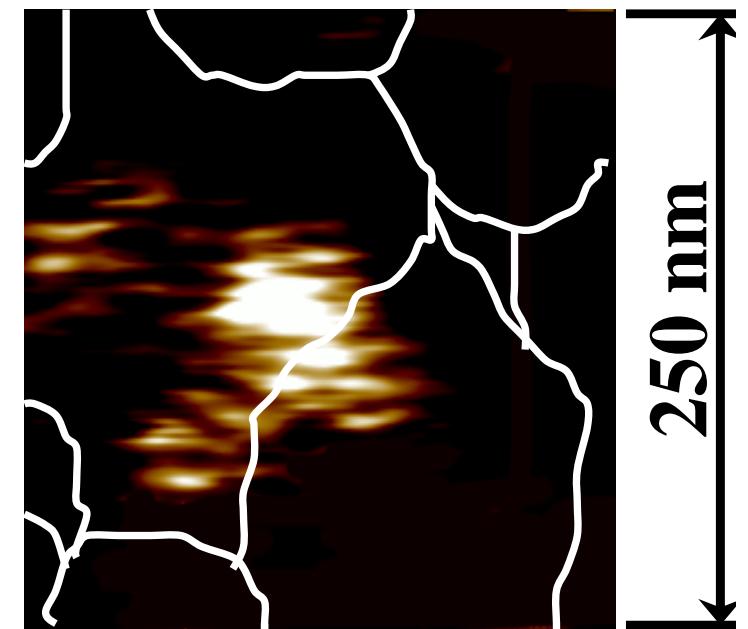
5 min after dc



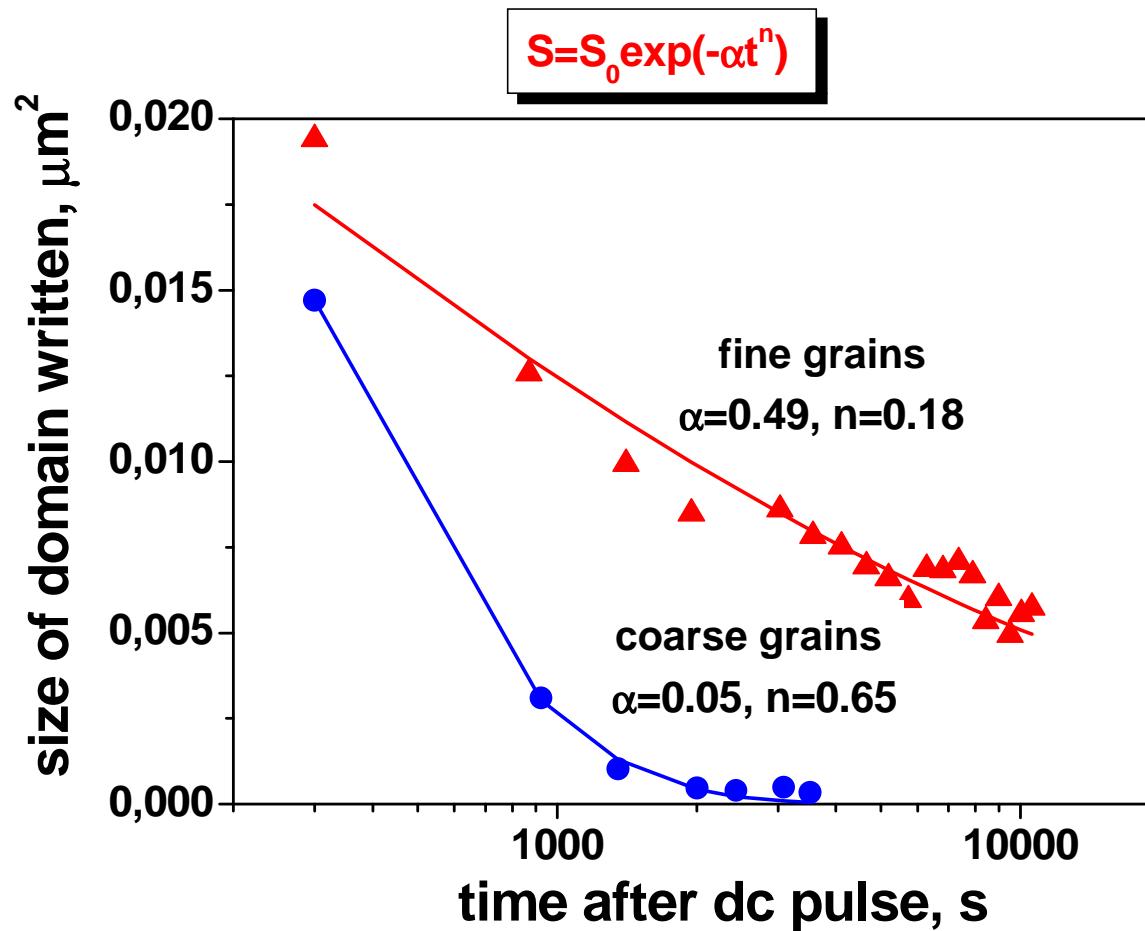
20 min



90 min

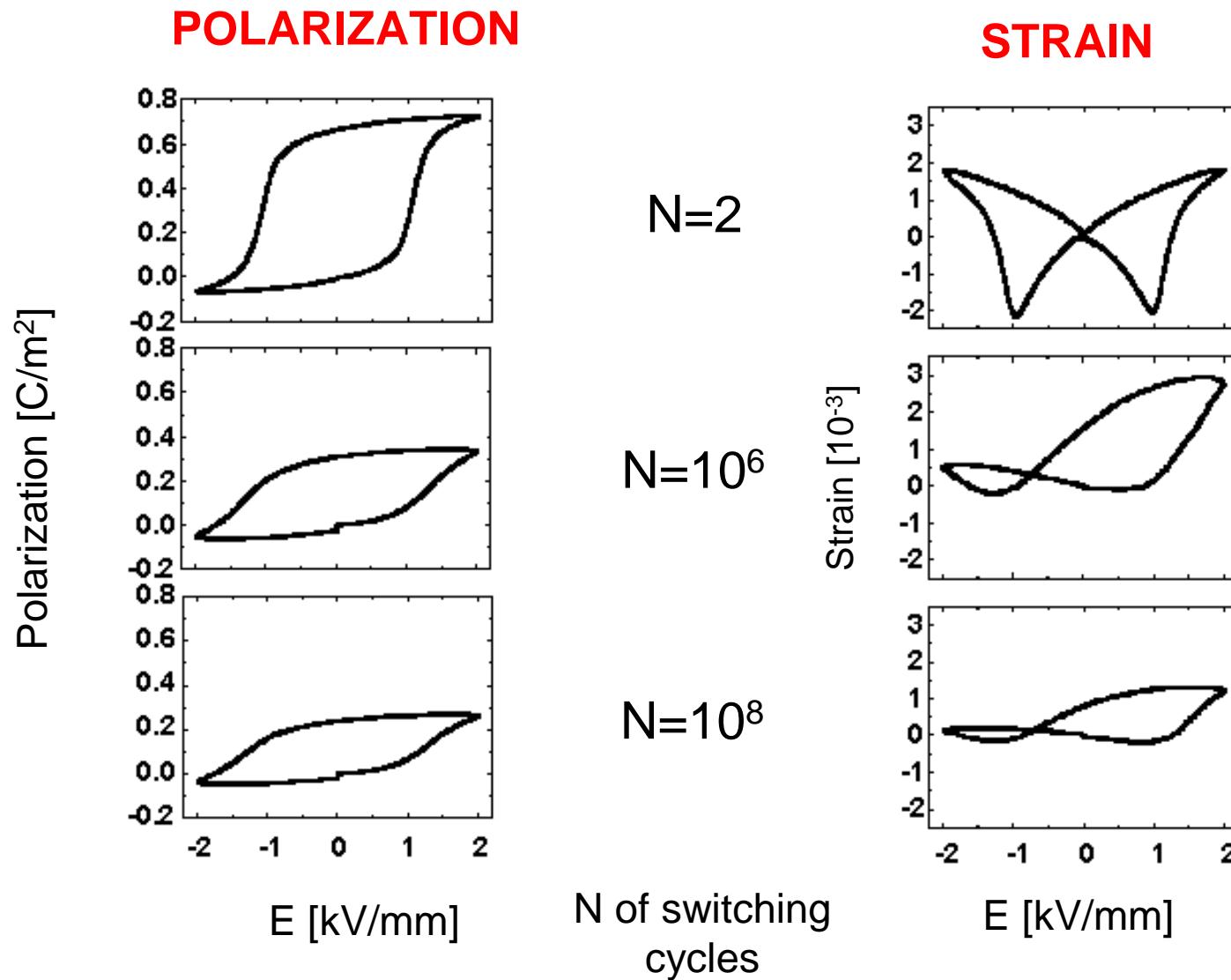


# Piezorelaxation vs. grain size



- Stretched exponential dependence is typical for systems with the broad spectrum of relaxation times

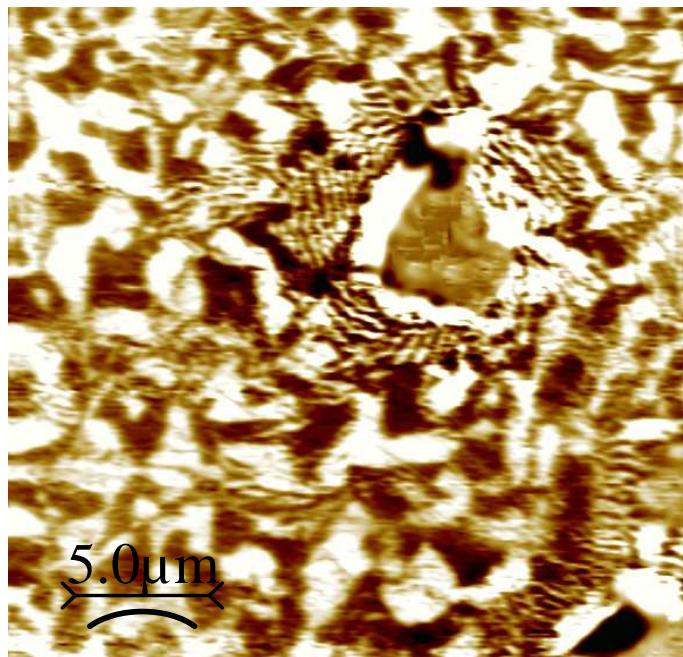
# *Bipolar fatigue in piezoceramics*



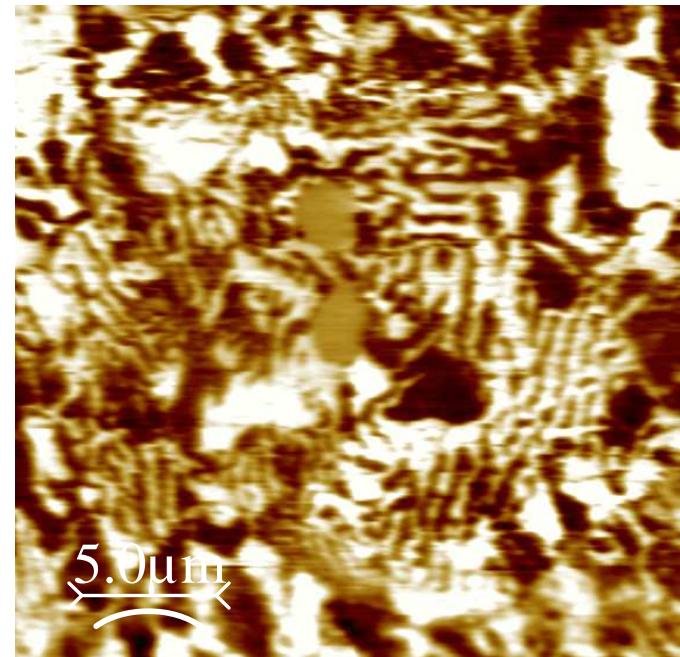
- Strain asymmetry as a result of fatigue-induced polarization offset

# *Fatigued vs. virgin piezoceramics (PIC151)*

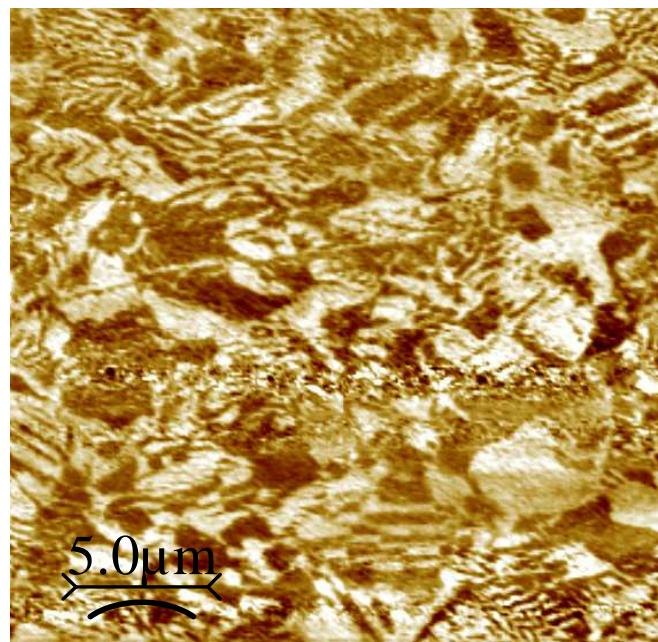
**N=0**



**N=10<sup>6</sup>**

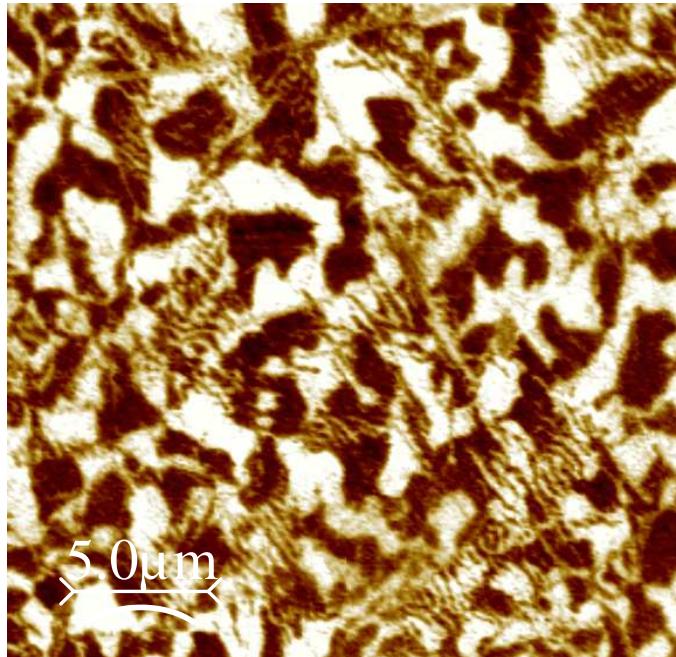


**N=5x10<sup>7</sup>**

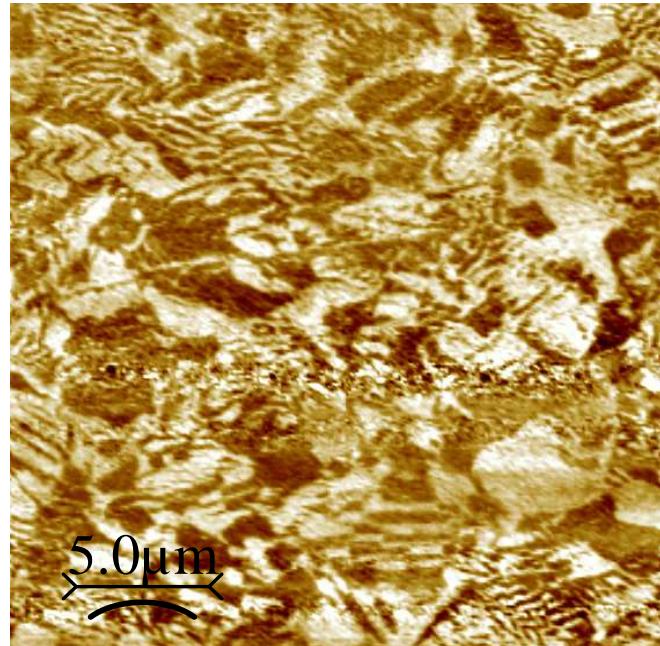


# *Effect of annealing*

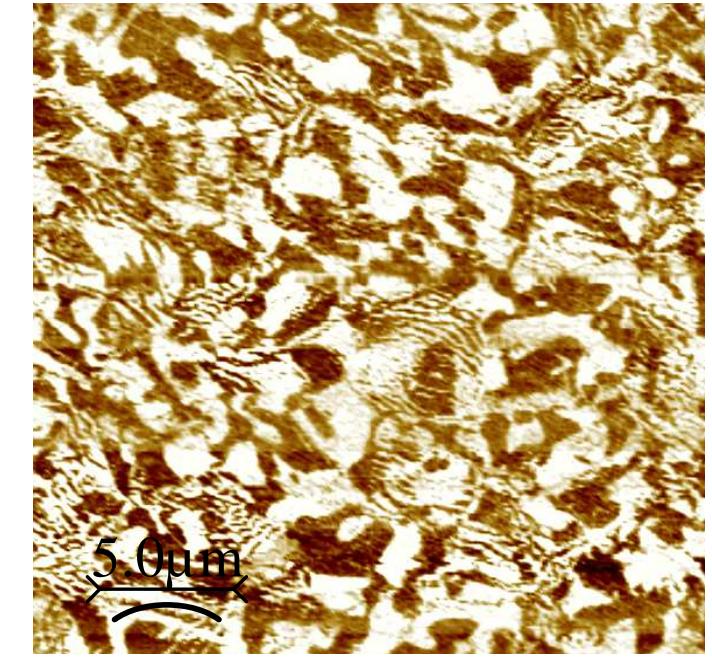
virgin sample



fatigued



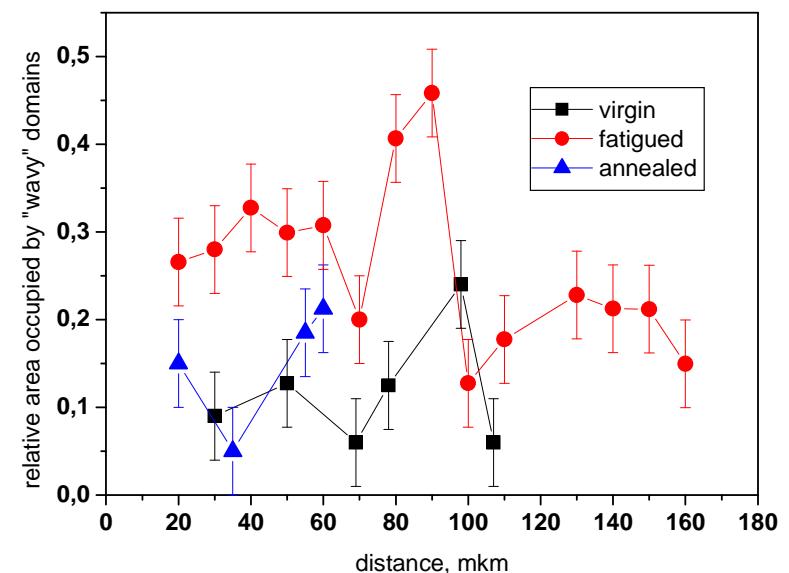
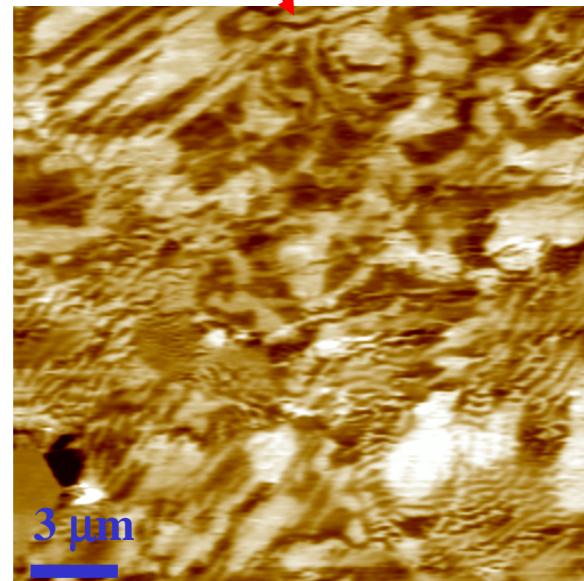
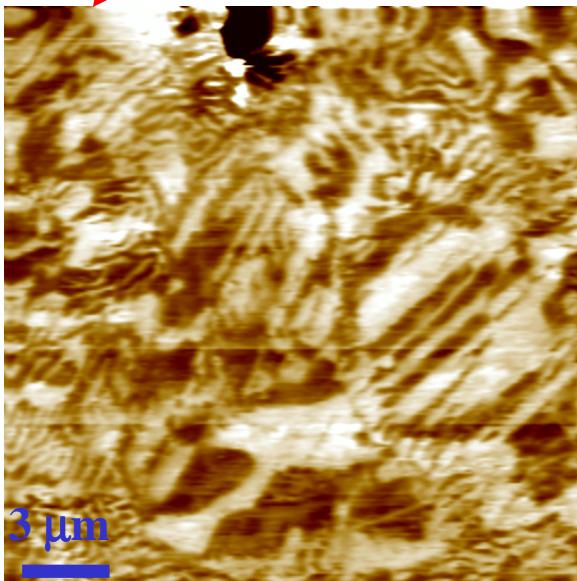
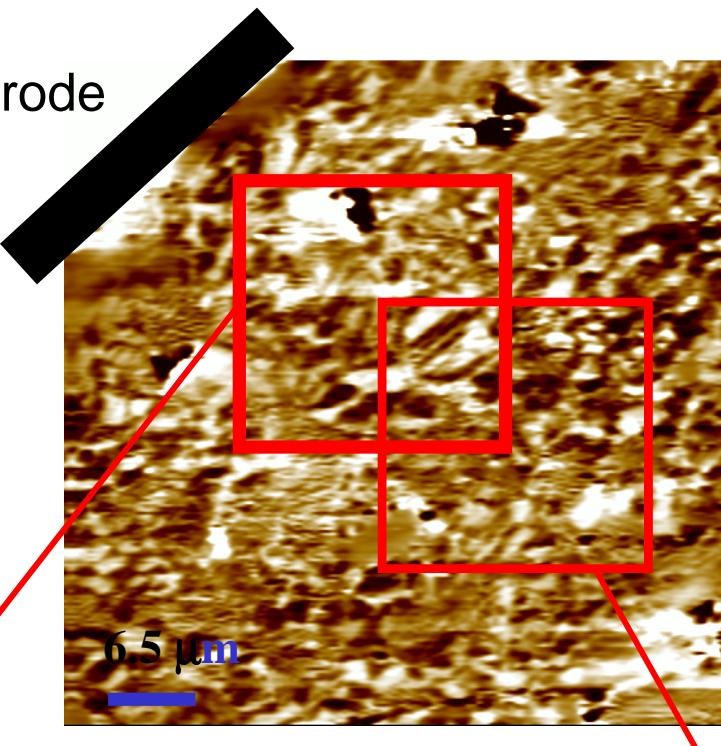
annealed



After annealing (**T=700 °C, 10 h**) the initial ratio between  $180^\circ$  and ferroelastic domains is restored along with the rejuvenation of switchable polarization.

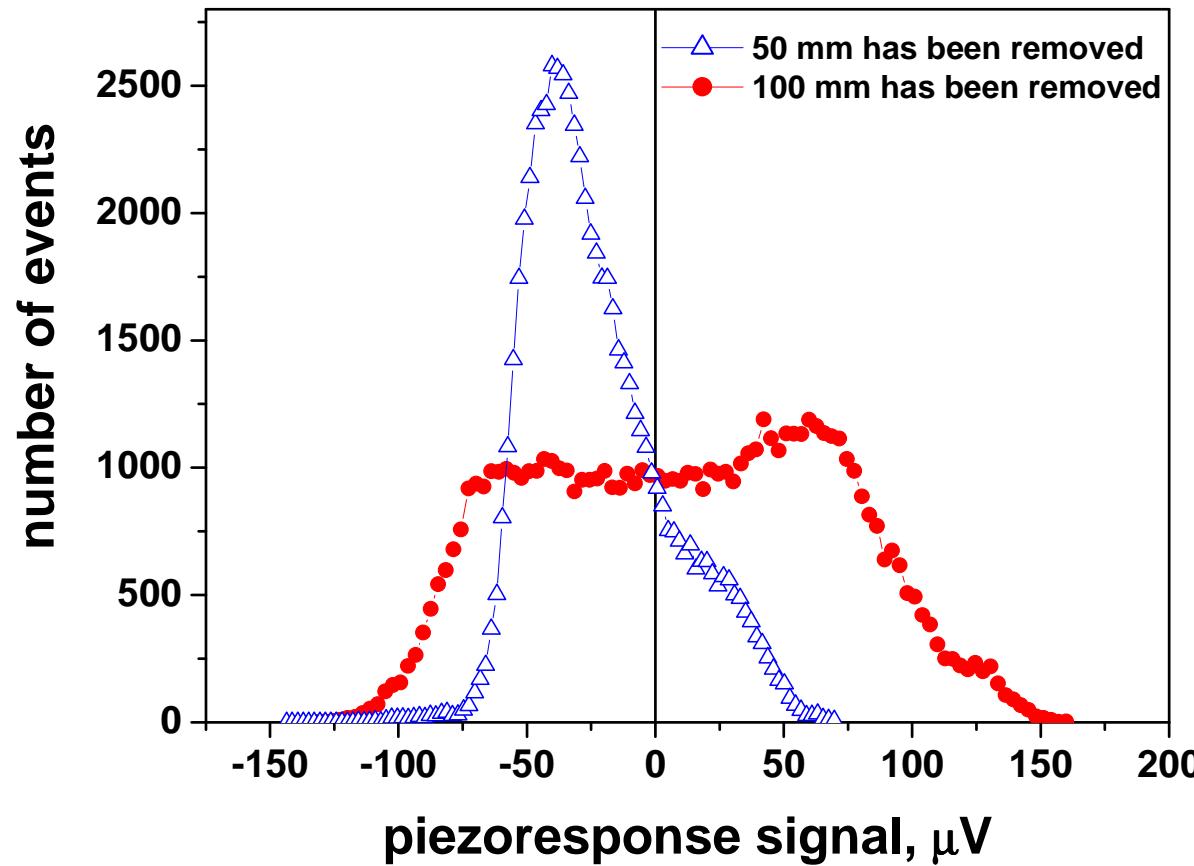
# *Concentration of stripe domains vs. depth*

Electrode



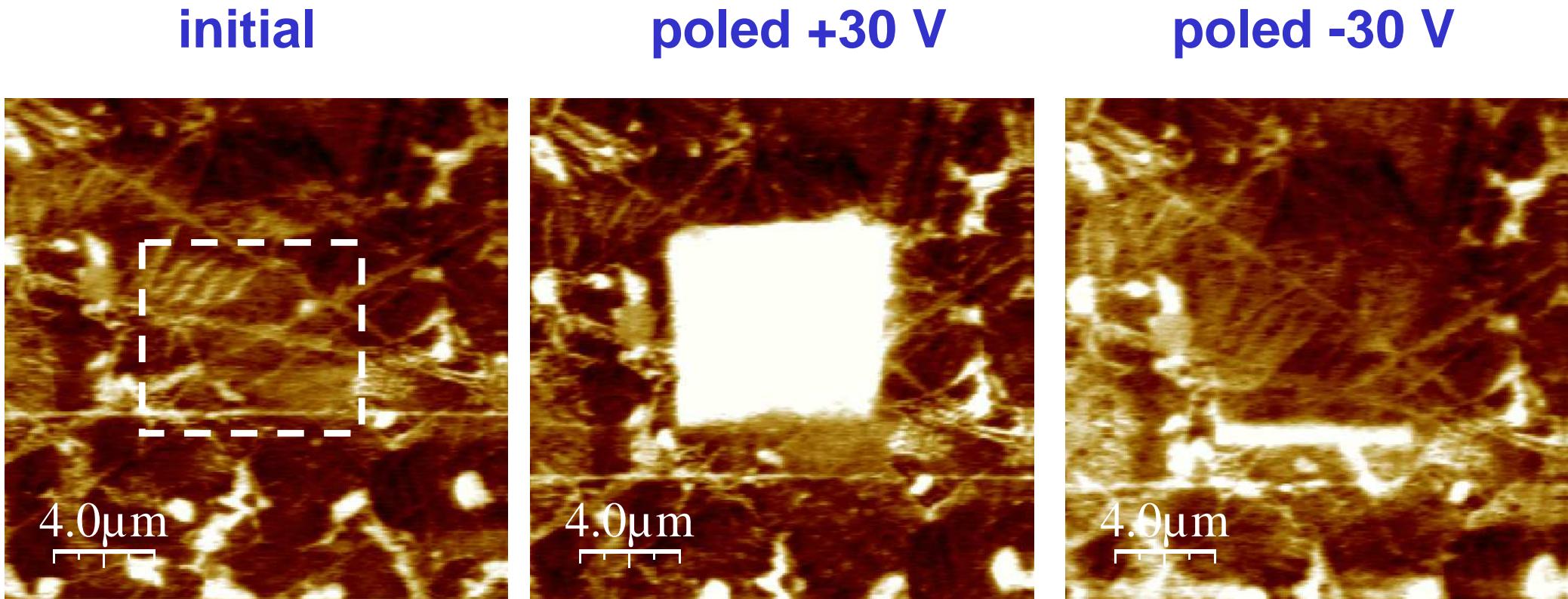
- Concentration of stripe domains decreases with the distance from the electrode
- A maximum at about 80-100  $\mu\text{m}$  below the electrode

# *Domain distribution vs. depth*



- In fatigued areas the domain histogram is highly asymmetrical demonstrating a preference of polarization state (domain head at electrode)
- The symmetrical domain distribution is restored with depth

# *Local polarization switching vs. fatigue*



- Local poling is different for both directions of poling field
- Pinning mechanism might be applicable for fatigue in ceramics

# *Conclusions:*

- *Miniaturization requires integration of piezoelectric/ferroelectric materials using Si-based technology*
- *Laser interferometry is an indispensable technique for the piezoelectric measurements in thin films down to nm thickness*
- *Piezoresponse Force Microscopy (PFM) can be used for integrated piezoelectrics provided the mechanism of the tip-surface interaction is understood*
- *PFM is suitable for the local study of polarization reversal, aging and fatigue in ferroelectrics including relaxors*