

# Surface Impedance in Spin-triplet Superconducting Proximity Structures



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## Collaborators



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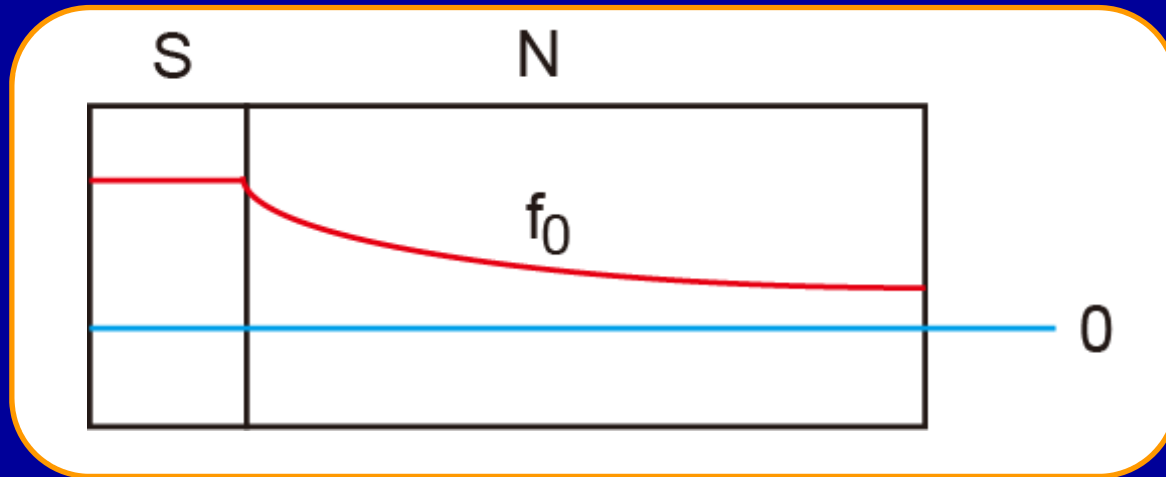


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# Introduction

Proximity effect



$$f_0 \approx e^{-x/\xi_T}$$

$$\xi_T = \sqrt{\frac{D_0}{4\pi T}}$$

Tunneling spectra of NS  
Josephson current in SNS

What happen on N attached to spin-triplet superfluids?

# Surface Impedance

Response to electromagnetic field

$$Z = \frac{4\pi}{c} \frac{E}{H} = R + iX$$

reflects various properties of quasiparticle states

R: surface resistance

loss of microwave power due to normal carrier

X: reactance

response of superconducting carriers

# Usual electromagnetic property

London's Eq.

$$\mathbf{J} = -\frac{c}{4\pi\lambda^2}\mathbf{A}$$

$$\lambda^2 = \frac{mc^2}{4\pi n_s e^2}$$

Together with a Maxwell's Eq.

$$\nabla \times \mathbf{H} = \frac{4\pi}{c}\mathbf{J} \quad (\text{static case})$$

$$\nabla \times \nabla \times \mathbf{H} = -\frac{1}{\lambda^2}\mathbf{H} \quad (\mathbf{H} = \nabla \times \mathbf{A})$$

$$\left(\nabla^2 - \frac{1}{\lambda^2}\right)\mathbf{H} = 0$$

Magnetic field usually decays in superconductor.

# Unusual electromagnetic property

Expression of  $\lambda$  in linear response theory

$$\mathbf{J} = -\frac{c}{4\pi\lambda^2}\mathbf{A} \quad \lambda(x)^{-2} = \frac{K}{\xi_N^2(T_c)} \frac{T}{T_c} \sum_{\omega_n > 0} F_{\omega_n}^2(x)$$

Belzig, Bruder, Schon, PRB 53, 5727(96)

F: anomalous part of Usadel Green function

even-freq. proximity

F is real number

$$\lambda^2 > 0$$

$$\left(\nabla^2 - \frac{1}{\lambda^2}\right)\mathbf{H} = 0$$

decay of H

odd-freq. proximity

F is pure imaginary

$$\lambda^2 < 0$$

$$\left(\nabla^2 + \frac{1}{|\lambda^2|}\right)\mathbf{H} = 0$$

oscillations of H

# Phenomenological theory

$$\mathbf{j} = \mathbf{j}_n + \mathbf{j}_s = (\sigma_1 - i\sigma_2)\mathbf{E} \quad \omega \sim 10\text{GHz}$$

$$\sigma_1 = \sigma_0 \frac{n_n}{n} \quad \text{for } \omega\tau \ll 1$$

$$\sigma_2 = \sigma_0 \omega\tau \frac{n_n}{n} + \frac{c^2}{4\pi\lambda_L^2\omega} \quad R + iX = \left[ \frac{4\pi i\omega}{c^2(\sigma_1 - i\sigma_2)} \right]^{1/2}$$

$$\sigma_0 = \frac{ne^2\tau}{m}$$

Drude conductivity  
in normal state

$$\frac{c^2}{4\pi\lambda_L^2\omega} = \sigma_0 \frac{n_s}{n} \frac{1}{\omega\tau}$$

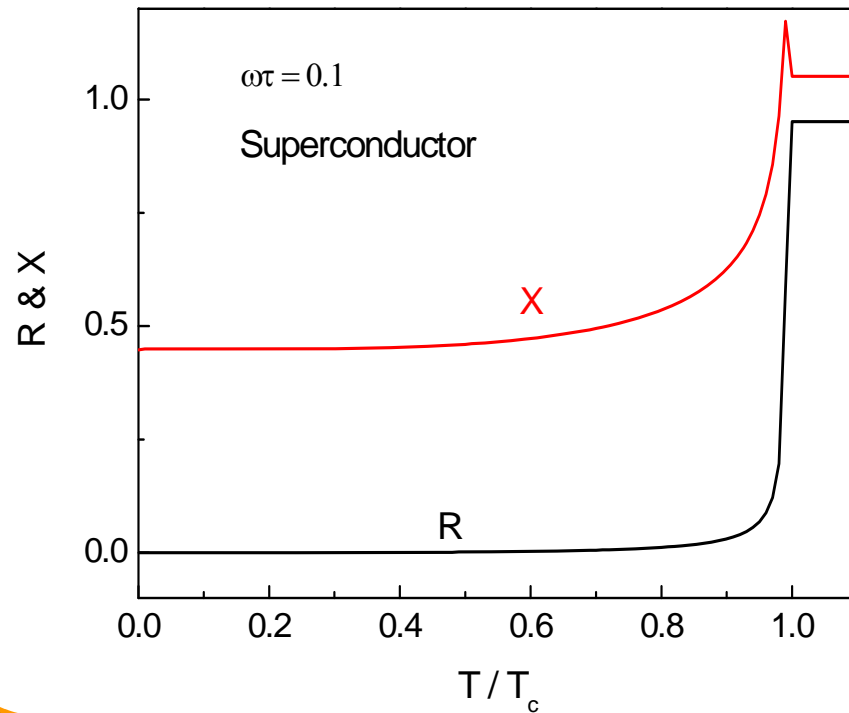
$$n = n_n + n_s$$

A simple relation

$$T < T_c \quad R < X$$

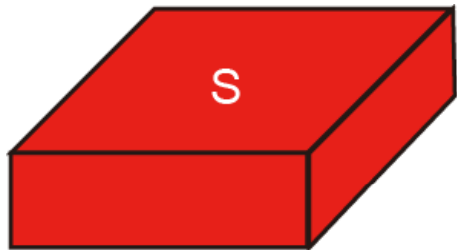
$$T \geq T_c \quad R \approx X$$

# Theoretical Results



$$n_s = n \frac{\Delta(T)}{\Delta(0)}$$

$$n_n = n - n_s$$

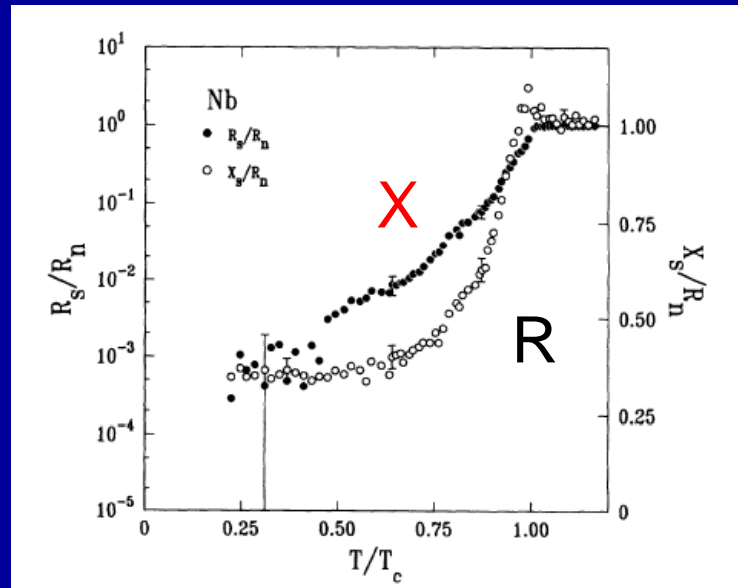


A simple relation

$$T < T_c \quad R < X$$

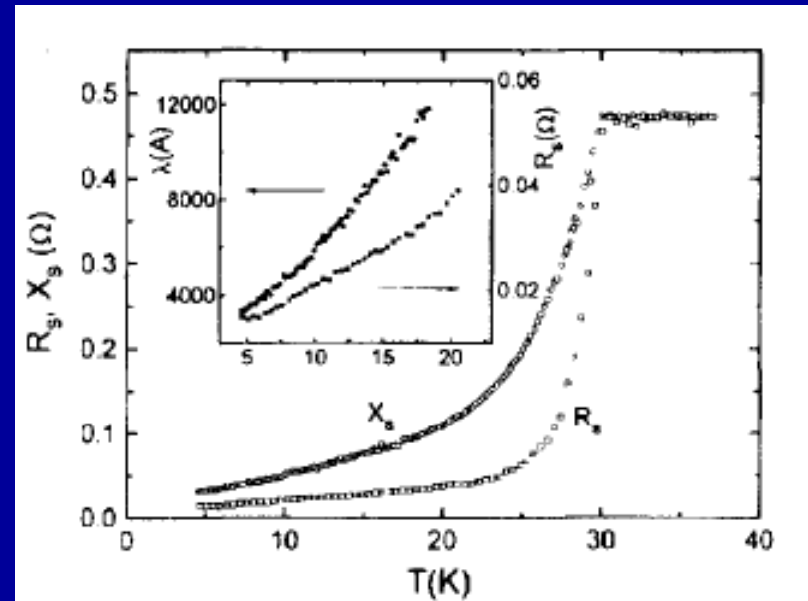
# Experimental results

Nb

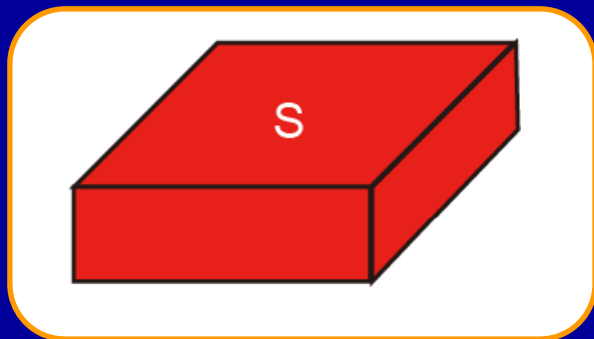


Klein, Nicol, Holczer, Gruner,  
PRB 50, 6307(94)

YBCO



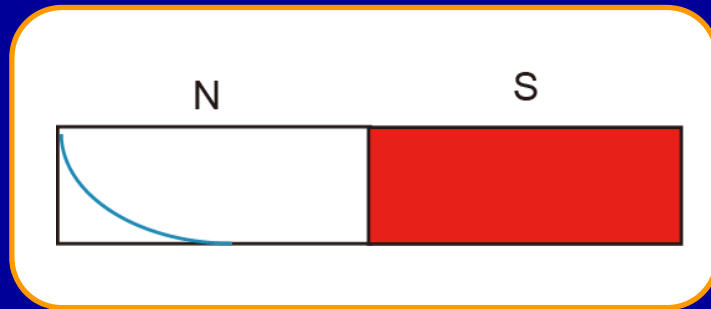
Trunin, Zhukov, Emel'chenko, Naumenko  
JETP Lett, 64, 832 (97)



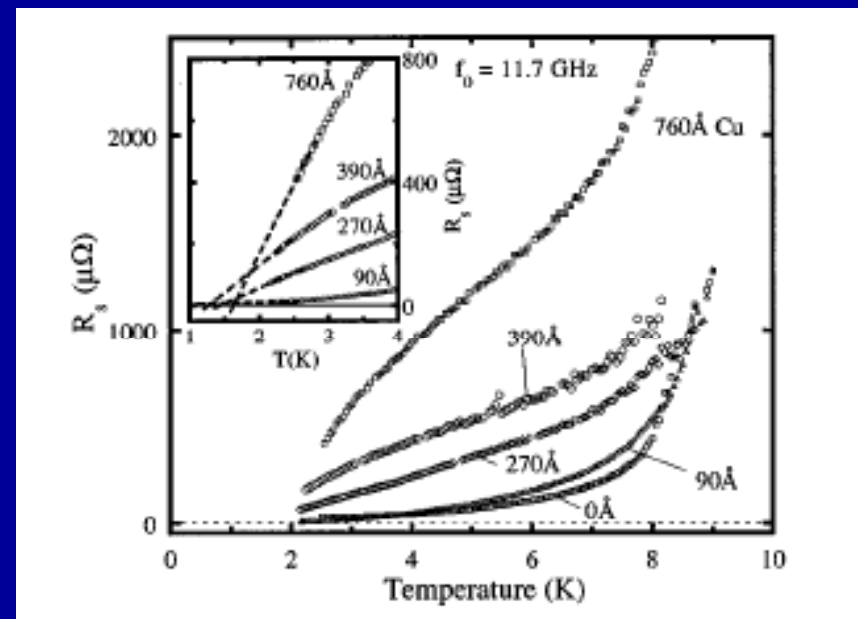
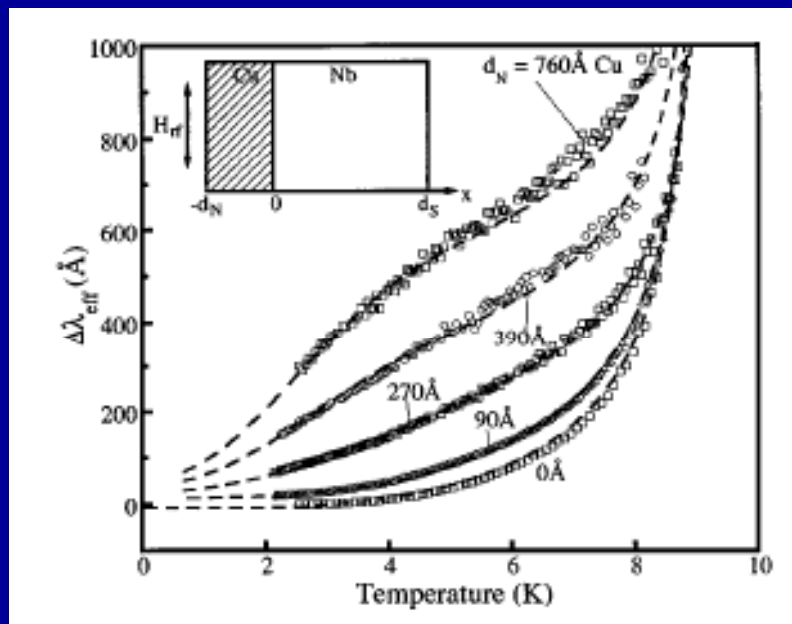
$$R < X$$



# Normal proximity effect



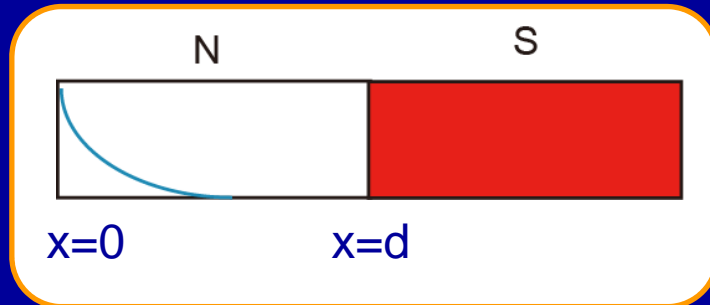
$$R < X$$



Pambianchi, Chen, Anlage, PRB 54, 3508(96)

# Proximity structures

## Strategy



Maxwell's Eqs.

$$E = \frac{\partial_x H}{\sigma}$$

$$H = -i \frac{c}{\omega} \partial_x E$$

Continuity of electromagnetic field

$$Z_{NS} = Z_N \frac{Z_S \cos(k_n d) + i Z_N \sin(k_n d)}{Z_N \cos(k_n d) + i Z_S \sin(k_n d)}$$

For  $0 < x < d$

$$H = y(H_1 e^{ik_n x} + H_2 e^{-ik_n x}) e^{i\omega t}$$

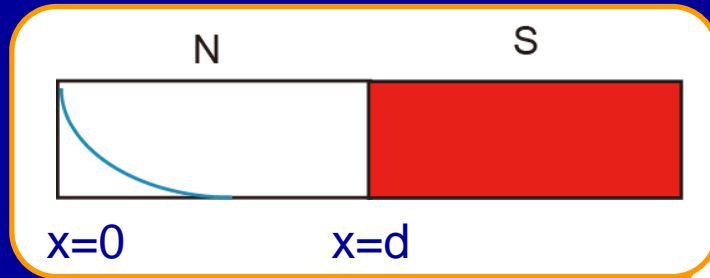
$$E = z(E_1 e^{ik_n x} + E_2 e^{-ik_n x}) e^{i\omega t}$$

For  $d < x$

$$H = y H_s e^{-ik_s x} e^{i\omega t}$$

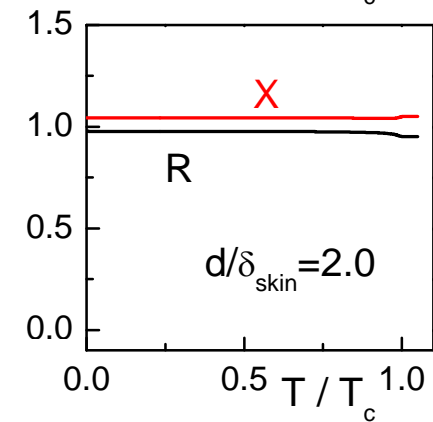
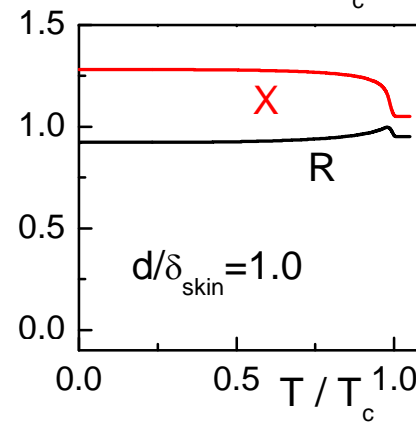
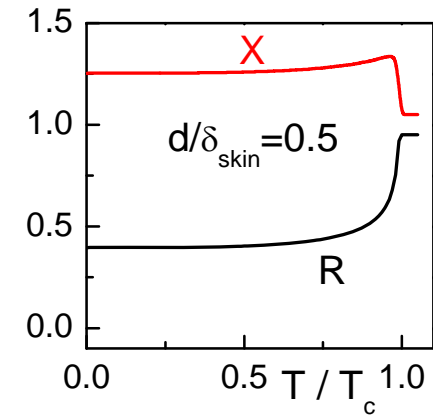
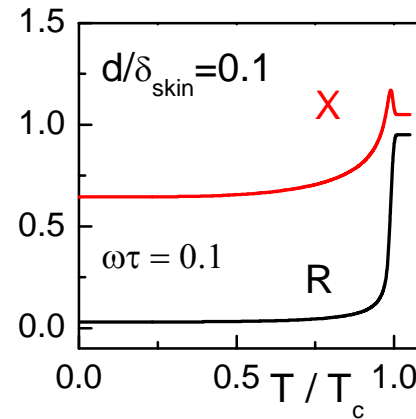
$$E = z E_s e^{-ik_s x} e^{i\omega t}$$

# I : zero proximity effect



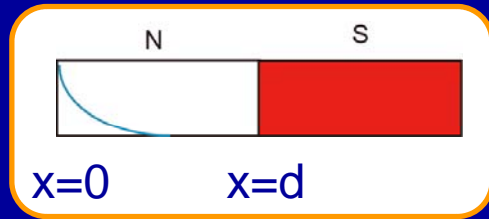
$$n_s = 0 \text{ in N}$$

$$R < X$$



No proximity effect

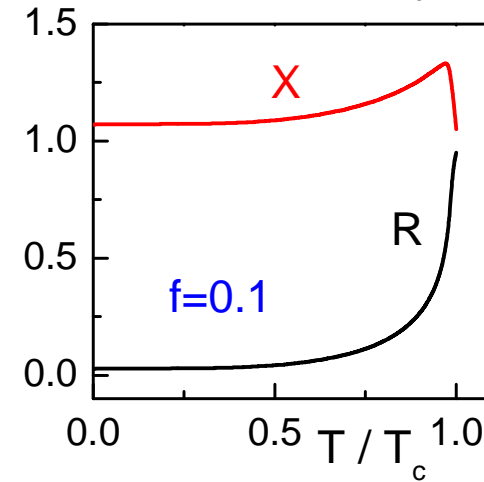
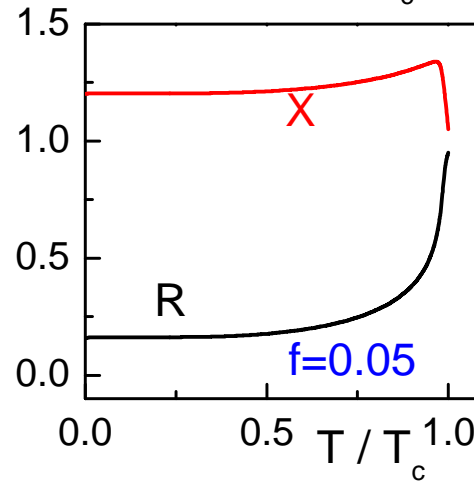
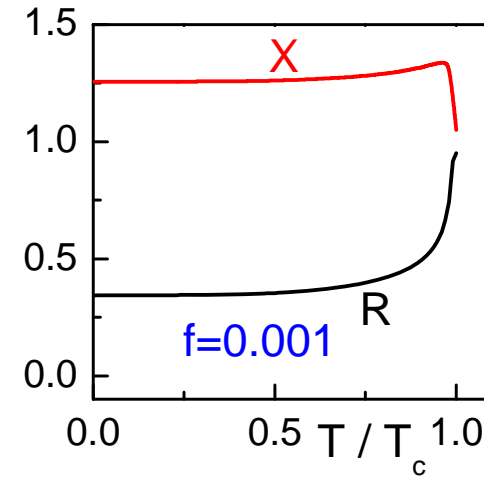
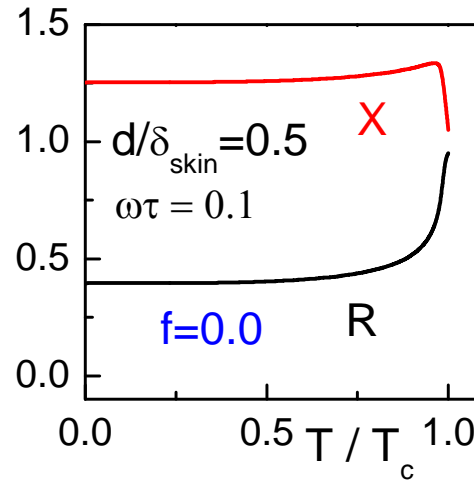
# II : usual proximity effect



f: strength of proximity effect

$$n_s = f n \frac{\Delta(T)}{\Delta(0)}$$

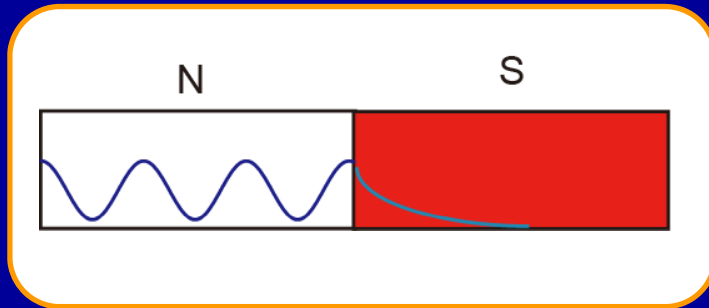
$$R < X$$



usual proximity effect

# Anomalous proximity effect

Spin-triplet p-wave



Tanaka, Asano, Golubov, Kashiwaya,  
PRB 72, 140503R (05);  
PRB 73, 059901 (E) (06).

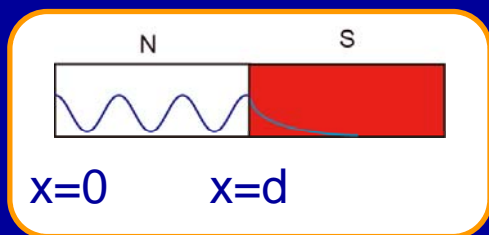
$\lambda^2 < 0$  due to odd-frequency

$$\sigma_1 = \frac{n_n e^2 \tau}{m}$$
$$\sigma_2 = \frac{n_n e^2 \tau}{m} \omega \tau + \frac{c^2}{4\pi \lambda_L^2 \omega}$$

$$\frac{c^2}{4\pi \lambda_L^2 \omega} = \sigma_0 \frac{n_s}{n} \frac{1}{\omega \tau}$$

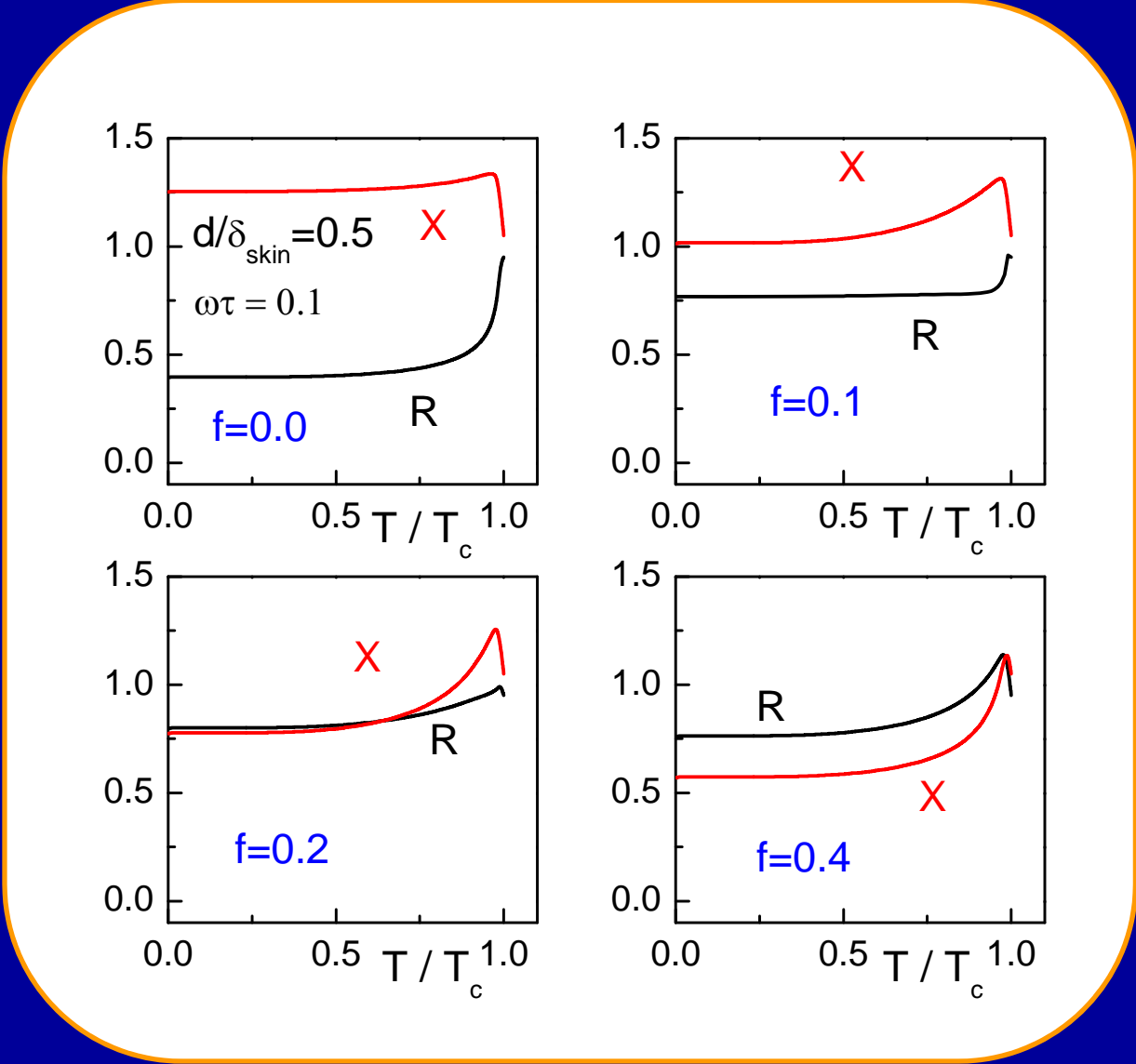
negative density

# III : anomalous proximity effect



$$n_s = -f n \frac{\Delta(T)}{\Delta(0)}$$

$$R > X$$

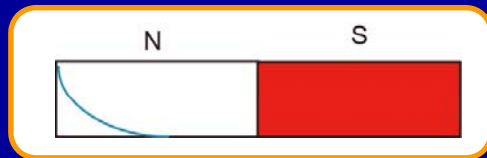


# Summary

surface impedance in NS proximity structures

$$Z = R + iX$$

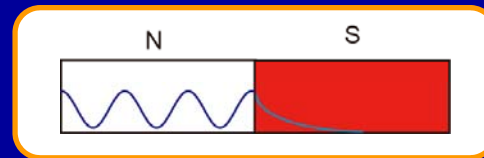
A phenomenological theory indicates  
spin-singlet



even freq. pairs

$$R < X$$

spin-triplet



odd freq. pairs

$$R > X \quad : \text{unusual relation}$$

Beyond the phenomenological theory....

Mattice & Bardeen, Phys. Rev. 111, 412(1958)

Nam, Phys. Rev. 156, 470(67); 156, 487(1967)

Abrikosov, Gor'kov, Dzyaloshinskii, Text (1962)

Trunin and Golubov, (2003)

# Unusual Impedance

$$\sigma_1 = \frac{n_n e^2 \tau}{m}$$
$$\sigma_2 = \frac{n_s e^2}{m\omega} = \frac{1}{\mu_0 \omega \lambda^2}$$
$$R + iX = \left[ \frac{i\omega\mu_0}{\sigma_1 - i\sigma_2} \right]^{1/2}$$

In singlet proximity structures (even freq. pair)

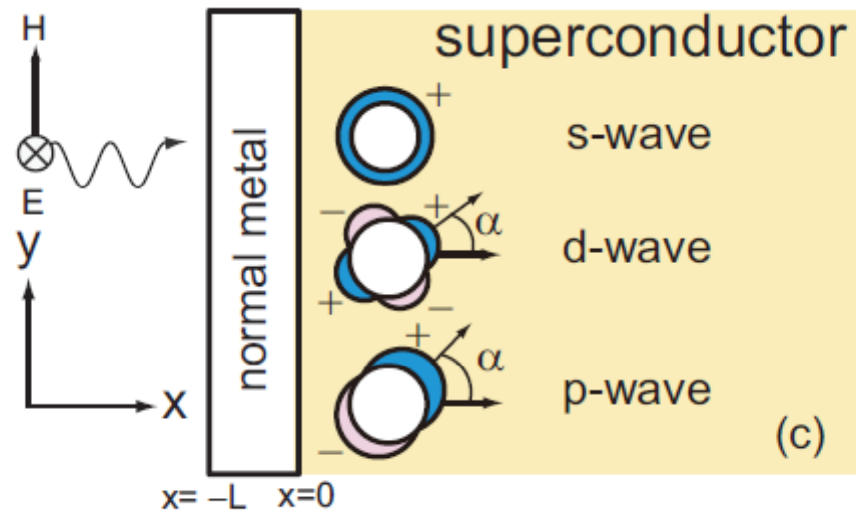
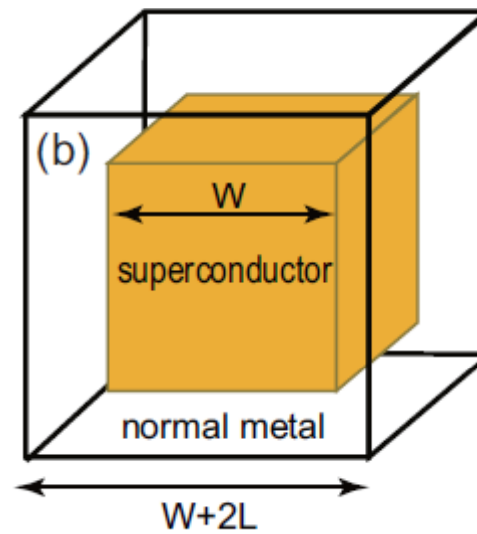
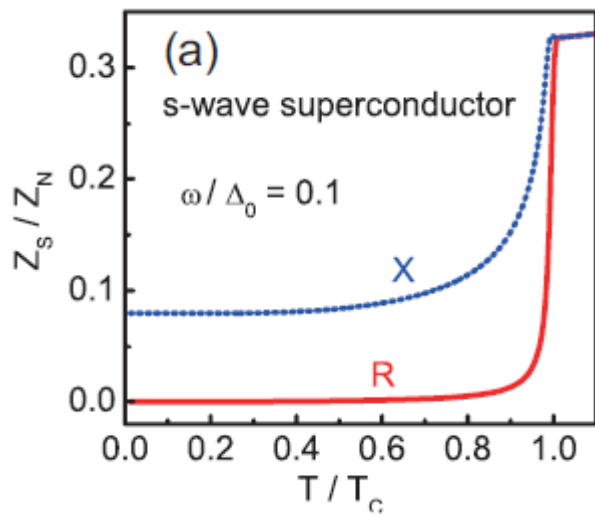
$$\lambda^2 > 0 \quad R < X$$

In triplet proximity structures (odd freq. pair)

$$\lambda^2 < 0 \quad \sigma_2 \text{ changes its sign!}$$

$$R > X$$





Formalism :

$$\left. \frac{\partial \theta(x, \epsilon)}{\partial x} \right|_{x=0} = \frac{1}{L} \frac{R_D}{R_B} \frac{\langle F \rangle}{T_B},$$
$$\langle F \rangle = \int_{-\pi/2}^{\pi/2} d\gamma \frac{T_N \cos \gamma (f_s \cos \theta_0 - g_s \sin \theta_0)}{(2 - T_N) + T_N (g_s \cos \theta_0 + f_s \sin \theta_0)},$$

$$\sigma_{NS}(\omega) = \sigma_1 + i\sigma_2,$$

$$\frac{\sigma_1}{\sigma_0} = \frac{1}{2\omega} \int_{-\infty}^{\infty} d\epsilon [J(\epsilon + \omega) - J(\epsilon)] K_1,$$

$$\frac{\sigma_2}{\sigma_0} = \frac{1}{2\omega} \int_{-\infty}^{\infty} d\epsilon [J(\epsilon + \omega) K_2 + J(\epsilon) K_3],$$

$$K_1 = f_I(\epsilon) f_I(\epsilon + \omega) + g_R(\epsilon) g_R(\epsilon + \omega),$$

$$K_2 = f_R(\epsilon) f_I(\epsilon + \omega) - g_I(\epsilon) g_R(\epsilon + \omega),$$

$$K_3 = f_R(\epsilon + \omega) f_I(\epsilon) - g_I(\epsilon + \omega) g_R(\epsilon),$$

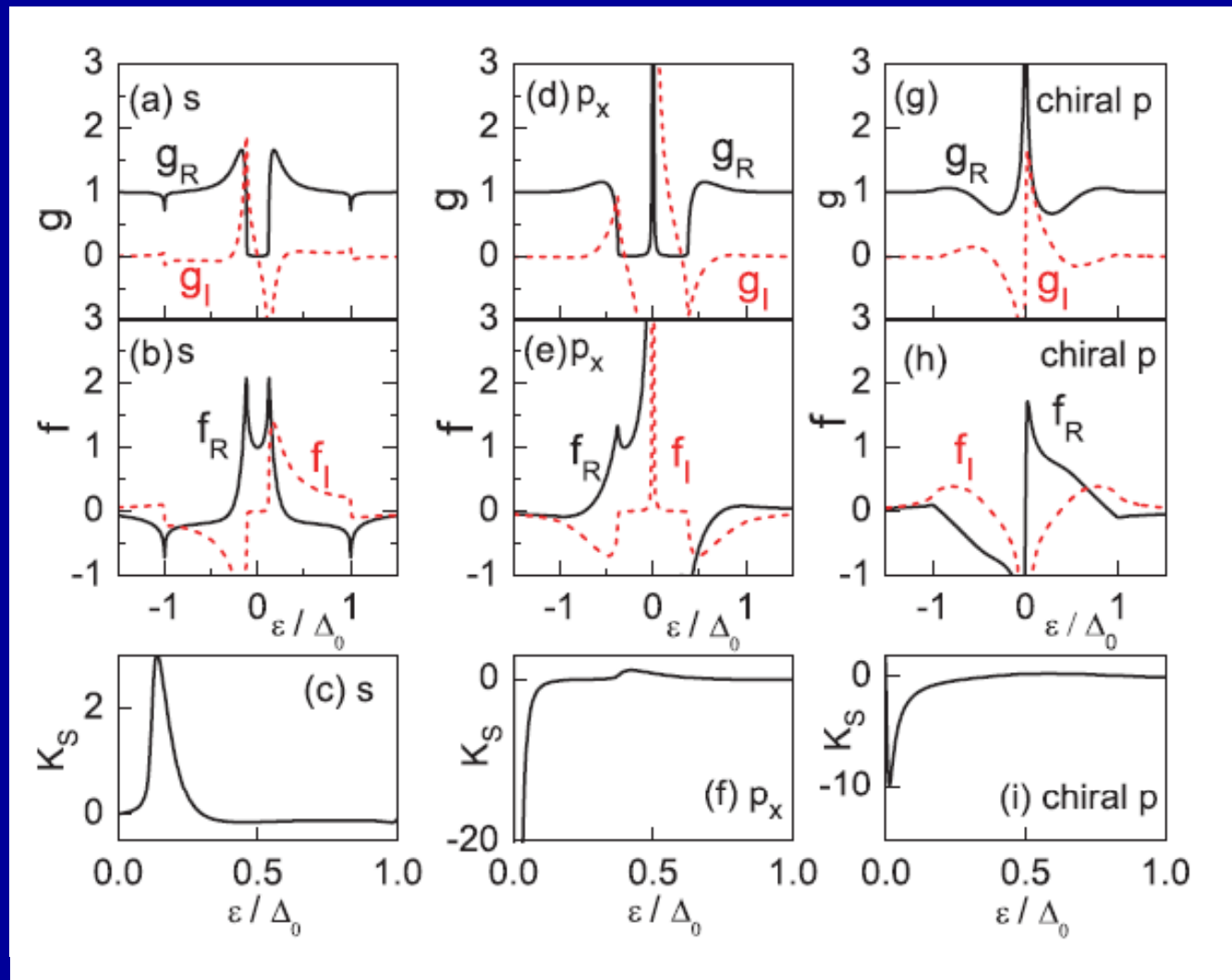
with  $J(\epsilon) = \tanh(\epsilon/2T)$  and

$$g_R(\epsilon) = \text{Re}[g(\epsilon)], \quad g_I(\epsilon) = \text{Im}[g(\epsilon)],$$

$$f_R(\epsilon) = \text{Re}[f(\epsilon)], \quad f_I(\epsilon) = \text{Im}[f(\epsilon)].$$

$$Z_{NS}(\omega) = R - iX = -i \sqrt{\frac{2i\omega}{\Delta_0} \frac{\sigma_0}{\sigma_{NS}(\omega)}} Z_N,$$

# The results of calculations: Green functions



# The results of calculations: impedance

