

Josephson supercurrent through a topological insulator surface state

*Nb/Bi<sub>2</sub>Te<sub>3</sub>/Nb junctions*

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# Acknowledgements

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## **MESA+ Institute for Nanotechnology, University of Twente, The Netherlands**

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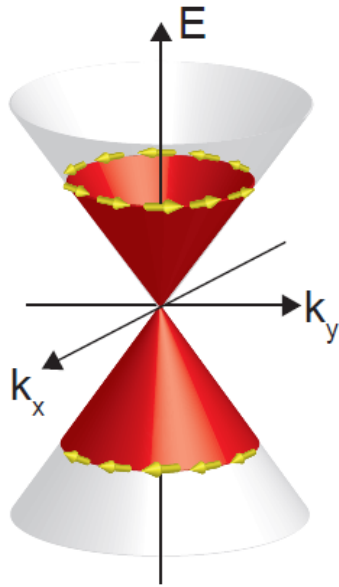
V. K. Guduru

## **University of Wollongong, Australia**

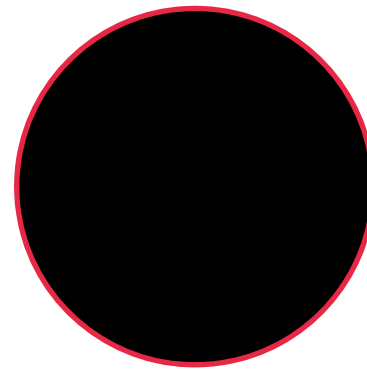
X. L. Wang

# Motivation

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TI surface states



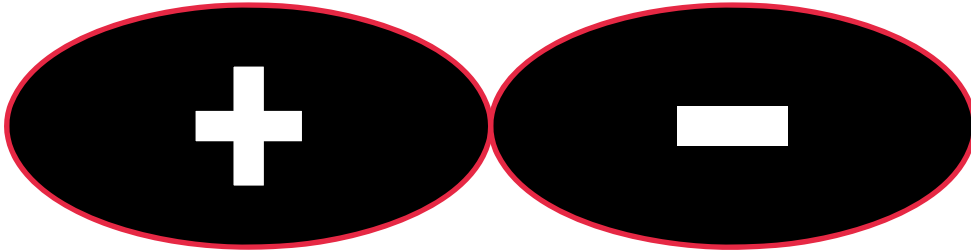
S-wave

Superconductor



# Motivation

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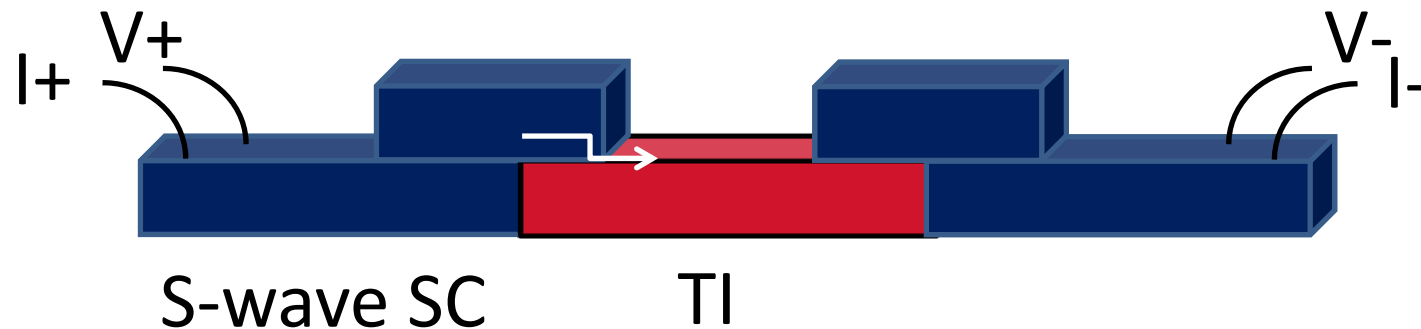
P-wave Superconductor

Natural place to look for Majorana fermions  
(single zero-energy modes)

# Motivation

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Essential: supercurrent must couple to surface states  
Characterize junction



# Content

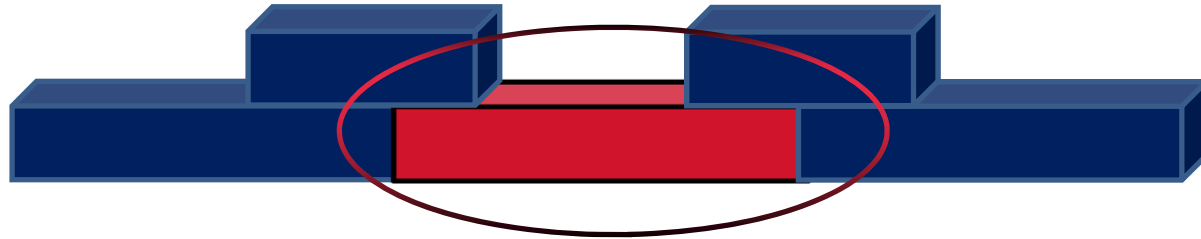
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- Part 1  $\text{Bi}_2\text{Te}_3$
- Part 2 S/TI/S junctions
- Part 3 Josephson supercurrent through the surface states

# Content

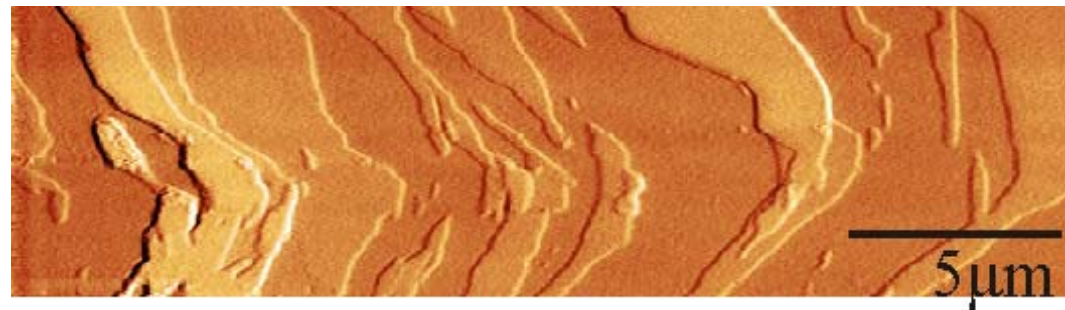
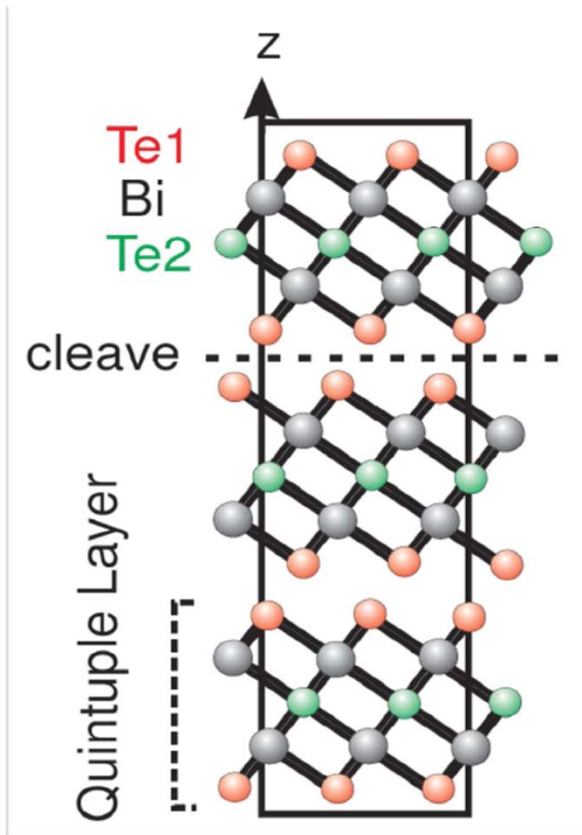
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- Part 1  $\text{Bi}_2\text{Te}_3$
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# Fabrication



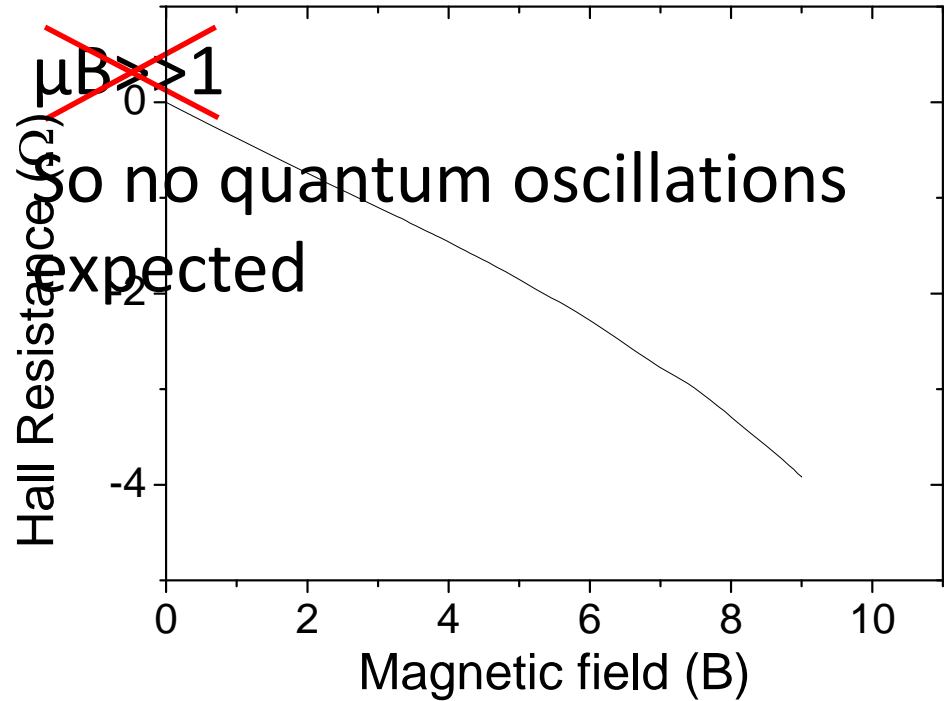
Mechanical cleavage

Photolithography contacts

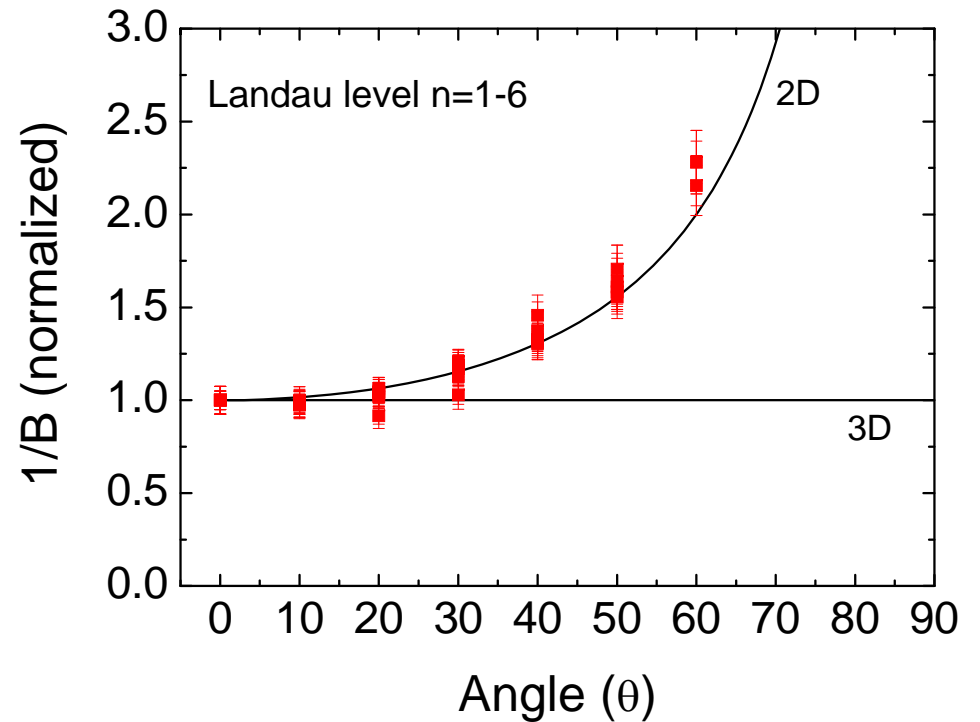
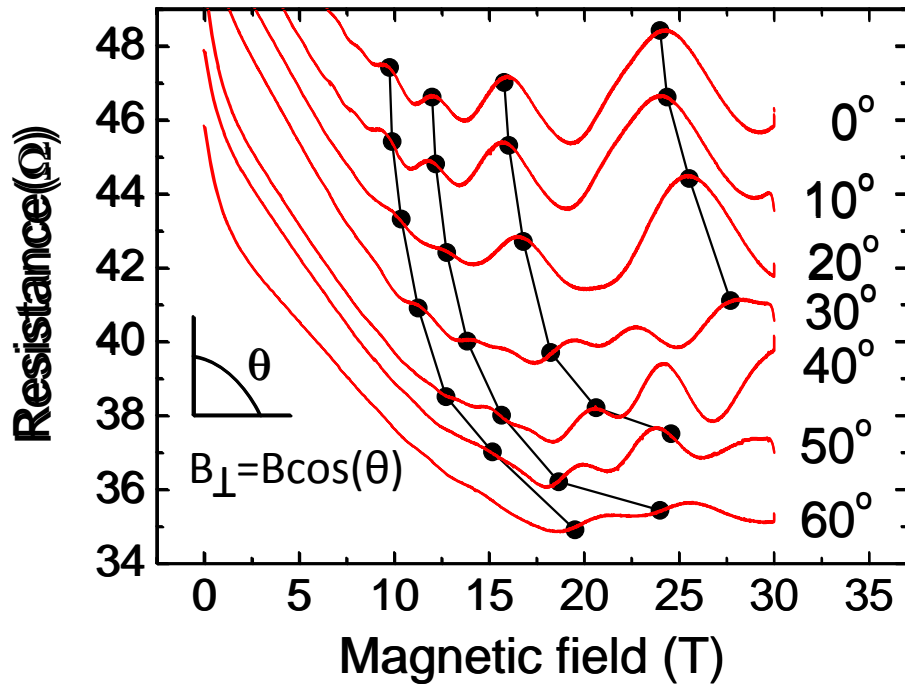
# Hall measurements

$n = 8.3 \times 10^{19} \text{ cm}^{-3}$   
 $\mu = 250 \text{ cm}^2/\text{Vs}$   
(bulk is conductive)

$l_{\text{mfp}} = 22 \text{ nm}$



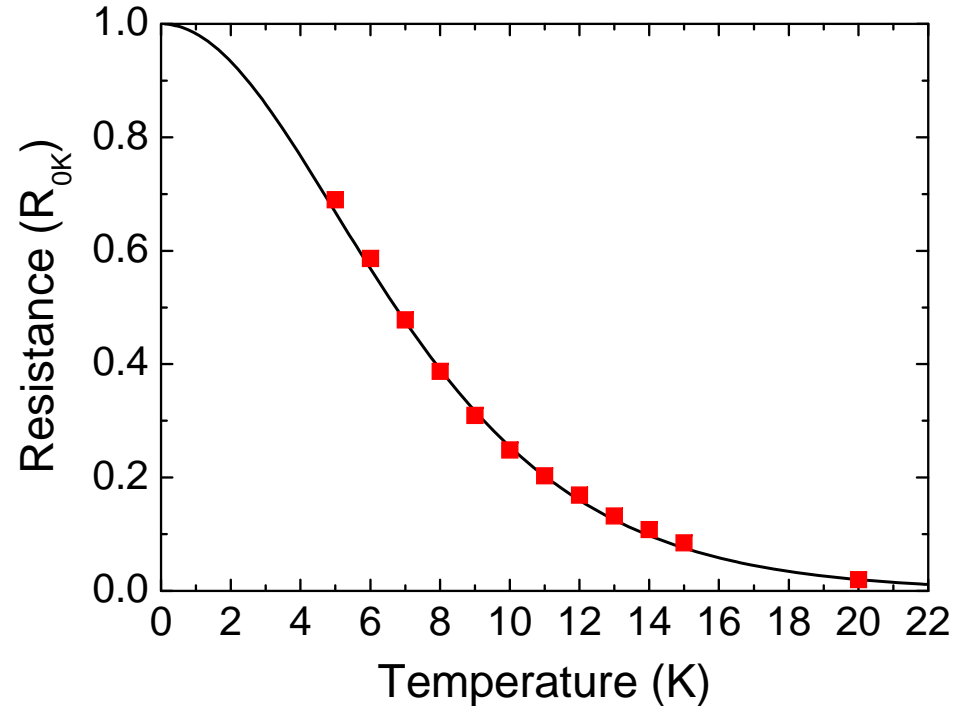
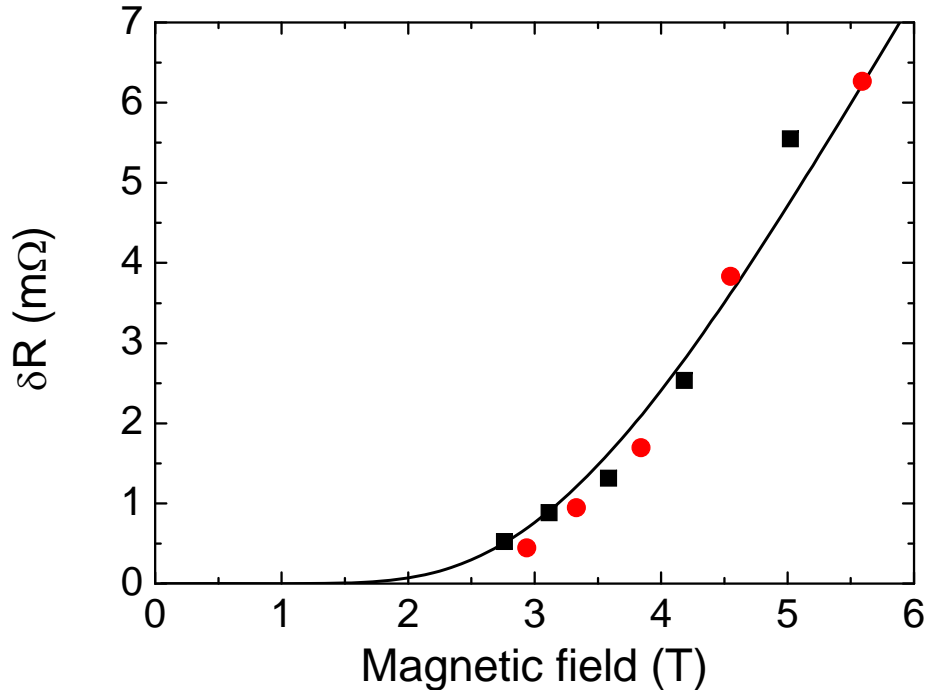
# Shubnikov-de-Haas oscillations



Left graph: Quantum oscillations @T=4.2K

Right graph: Oscillations from 2D channel

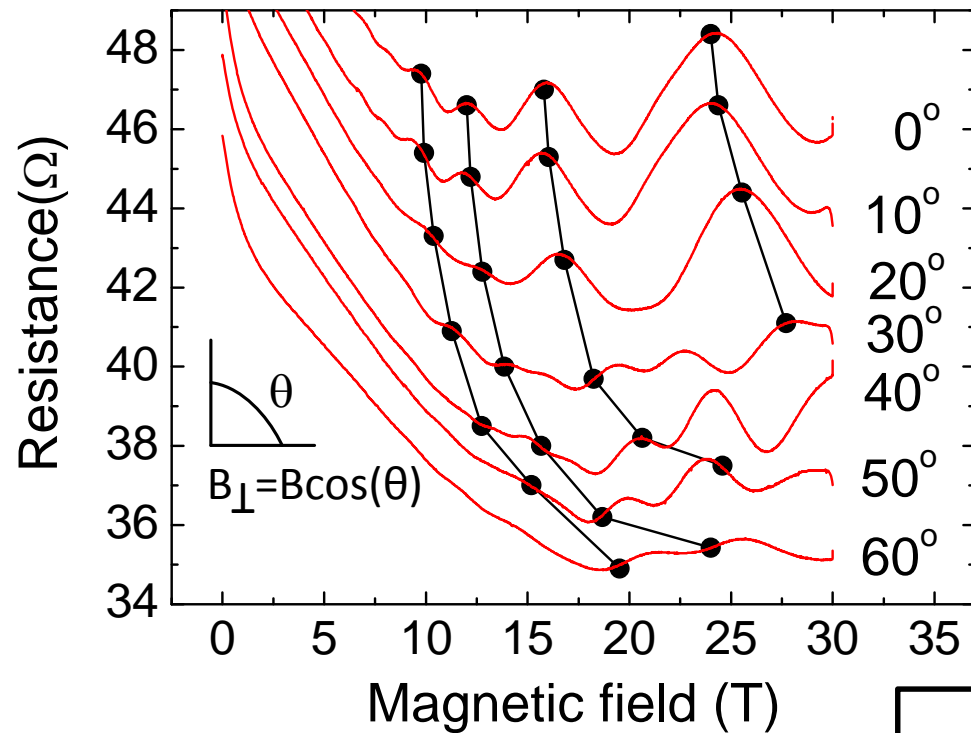
# Shubnikov-de-Haas oscillations



Left graph: Dingle temperature 1.65 K,  $\mu=8300 \text{ cm}^2/\text{Vs}$

Right graph: Effective mass  $0.16m_0$

# Shubnikov-de-Haas oscillations



$$\delta\rho_{xx} \sim \cos\left(\frac{2\pi E_f}{\hbar\omega_c} + \pi + \varphi_B\right)$$

# Lifshitz-Kosevich formalism

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$$\delta\rho_{xx} \sim \cos\left(\frac{2\pi E_f}{\hbar\omega_c} + \pi + \varphi_B\right)$$

Extrapolating till  $1/B=0$

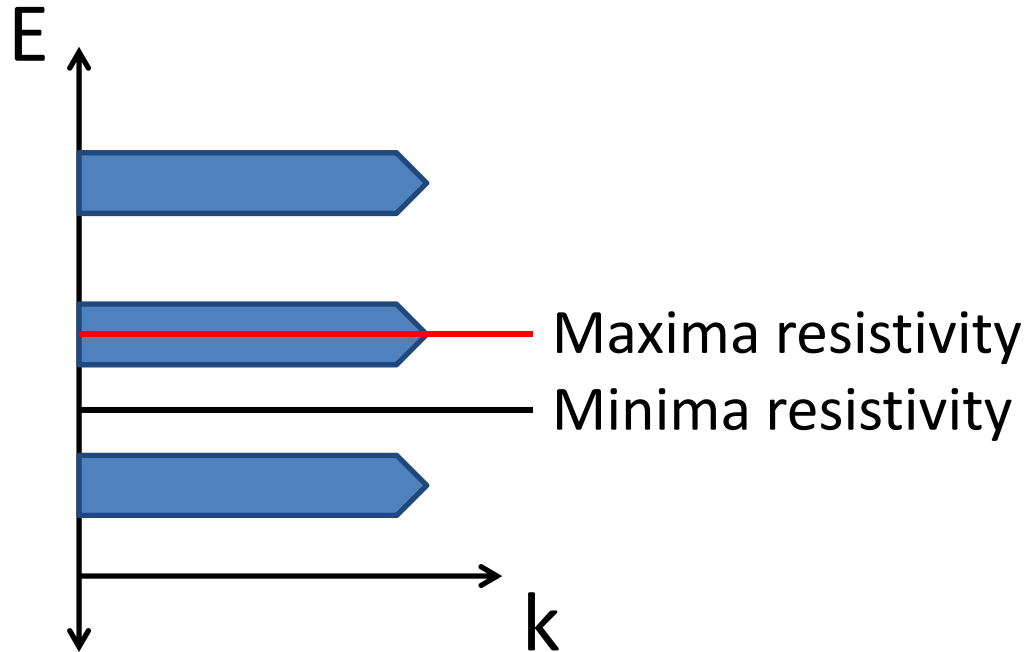
→

nth minima in resistivity

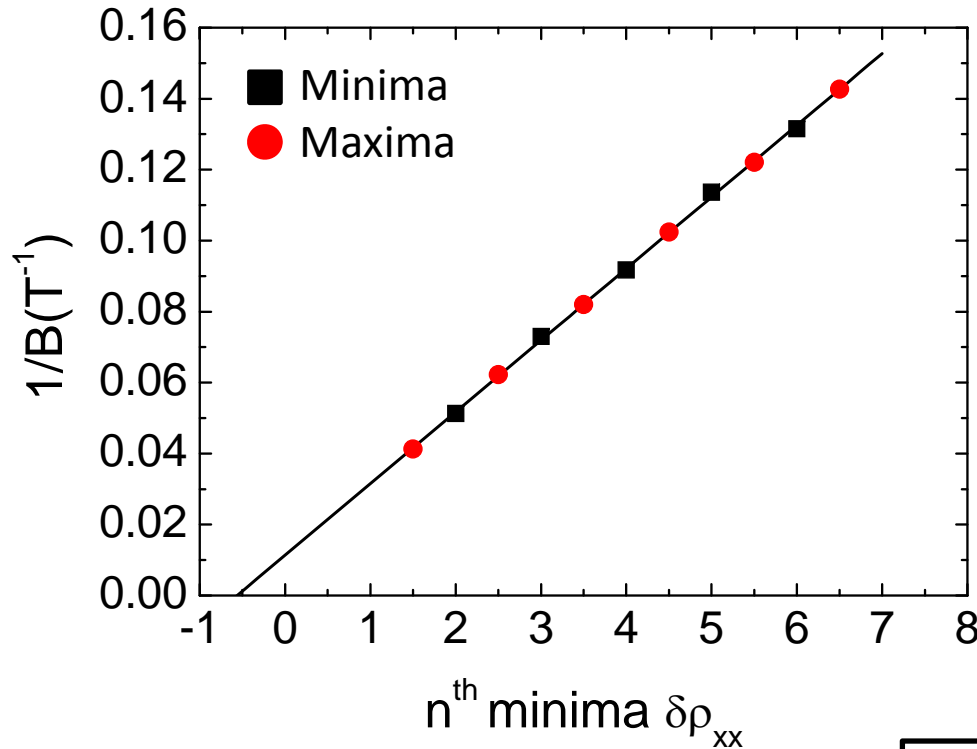
is zero at  $1/B=0$  in

'normal case' (Berry

phase=0)



# Lifshitz-Kosevich formalism



Extrapolating till  $1/B=0$

→

Berry phase is  $\pi$

$$\delta\rho_{xx} \sim \cos\left(\frac{2\pi E_f}{\hbar\omega_c} + \pi + \varphi_B\right)$$

# Lifshitz-Kosevich formalism

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Conductivity is the response function

$$\rho_{xx} = \frac{\sigma_{xx}}{\sigma_{xx}^2 + \sigma_{xy}^2}$$

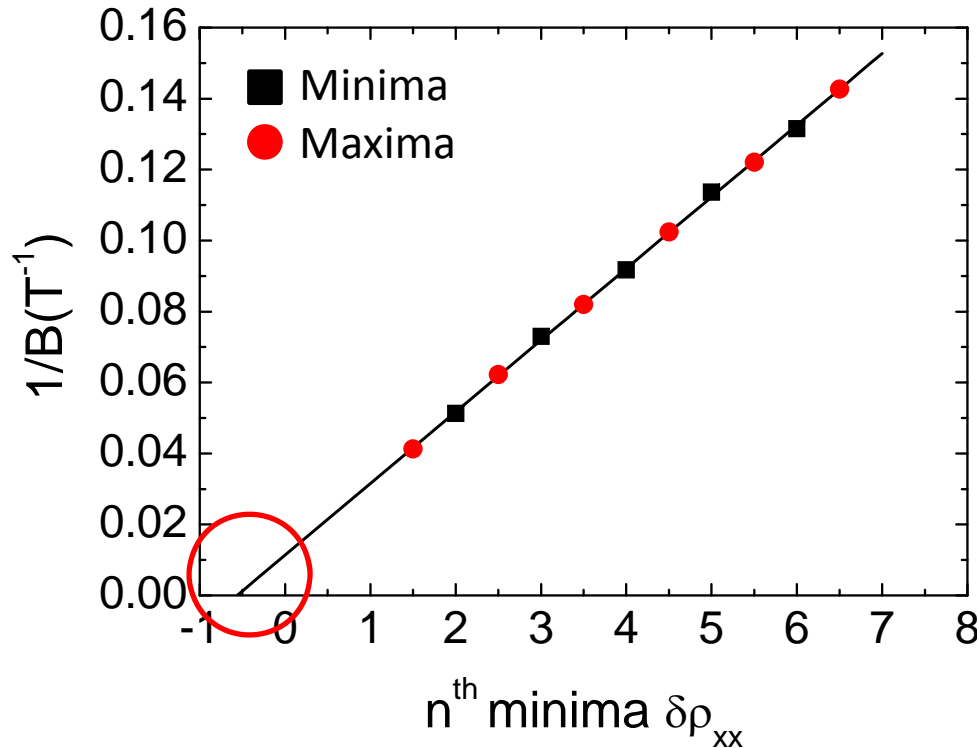
I.M. Lifshitz and A.M. Kosevich formalism applies if:

$$\frac{\delta\sigma_{xx}}{\langle\sigma_{xx}\rangle} \ll 1 \text{ or } \frac{\sigma_{xy}}{\sigma_{xx}} \gg 1 ; \text{Former is 0.01, later is 10}$$

Then  $\delta\rho_{xx} \sim \delta\sigma_{xx}$  So in normal case  $n^{\text{th}}$  minima=0 through  $1/B=0$



# Lifshitz-Kosevich formalism



- Surface states present
- @  $1/B=0$ ,  $n=-1/2$   
(Berry phase of  $\pi$ )  
→ linear dispersion relation
- $l_{\text{mfp}}=105$  nm

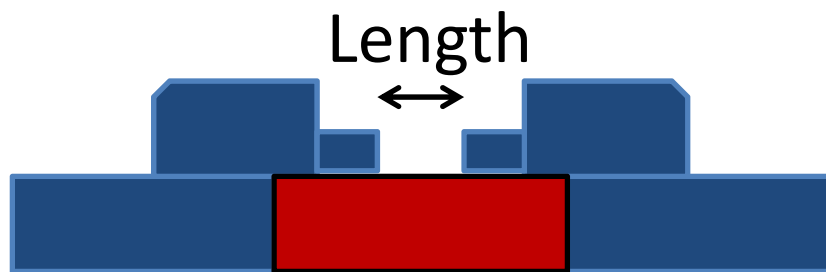
Non-trivial surface states present

# Content

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- Part 1  $\text{Bi}_2\text{Te}_3$
- Part 2 **S/TI/S junctions**
- Part 3 Josephson supercurrent through the surface states

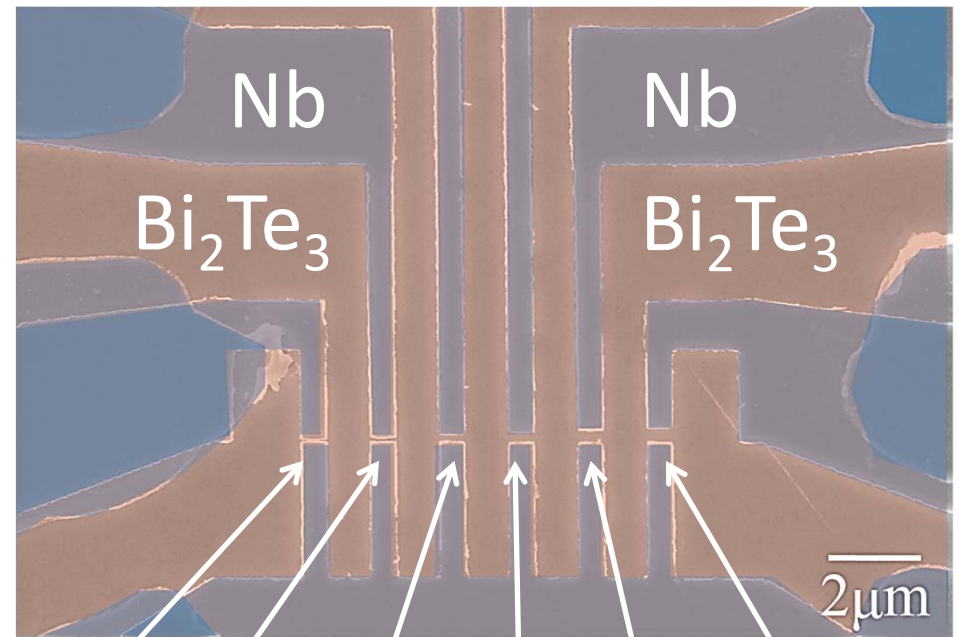
# S/TI/S junctions



Nb: blue

$\text{Bi}_2\text{Te}_3$ : red

Width=500 nm



50 100 150 200 250 300 nm

Junction lengths

# Josephson supercurrent

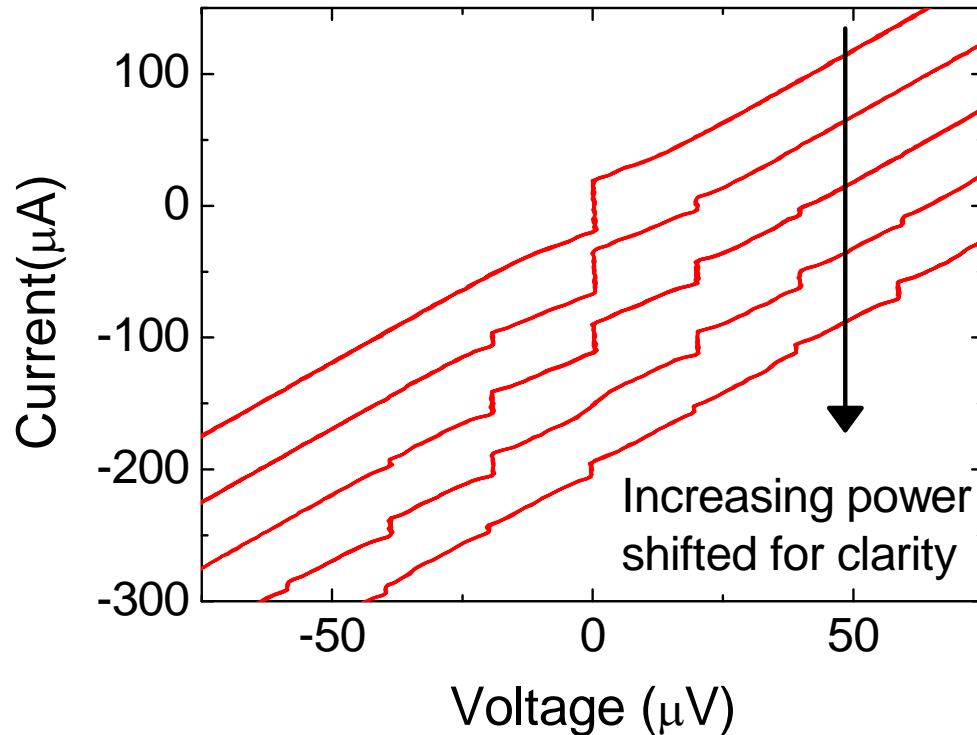
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Hallmarks for a Josephson junction:

- 1) Shapiro steps
- 2) Modulation  $I_c$  versus B-field

# First hallmark – Shapiro steps

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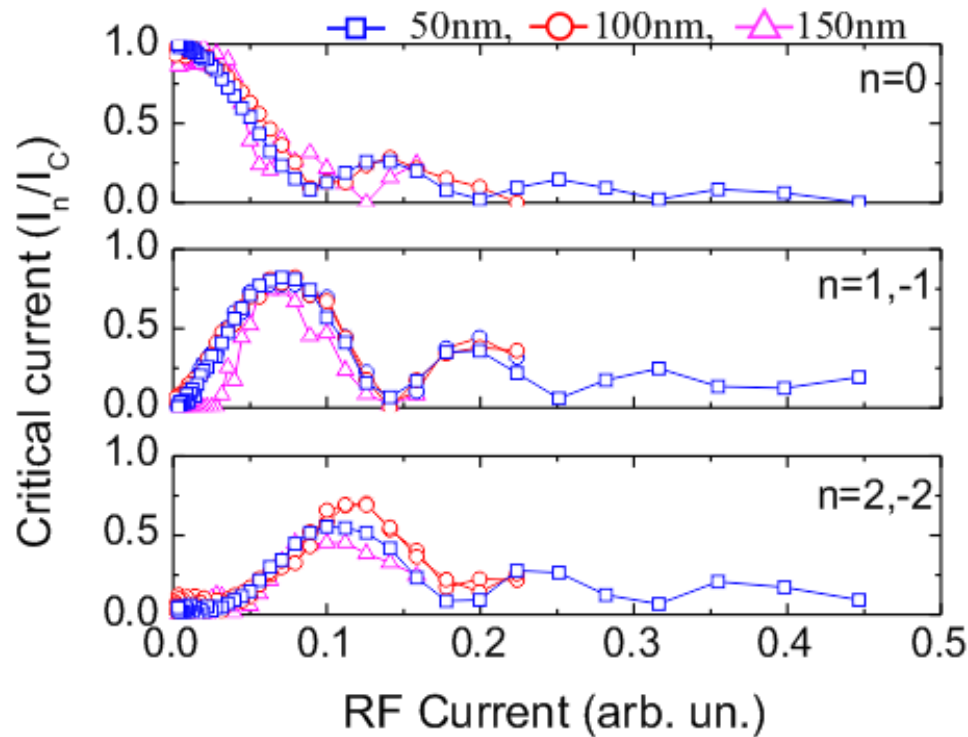
- Microwave frequency  $\omega$
- @  $2eV/\hbar = n\omega$  Shapiro steps
- (Energy Cooper pairs resonant to energy microwave)

$$\omega = 10 \text{ GHz}$$

$$V = n \times 20.7 \mu\text{V}$$

$$T = 1.6 \text{ K}$$

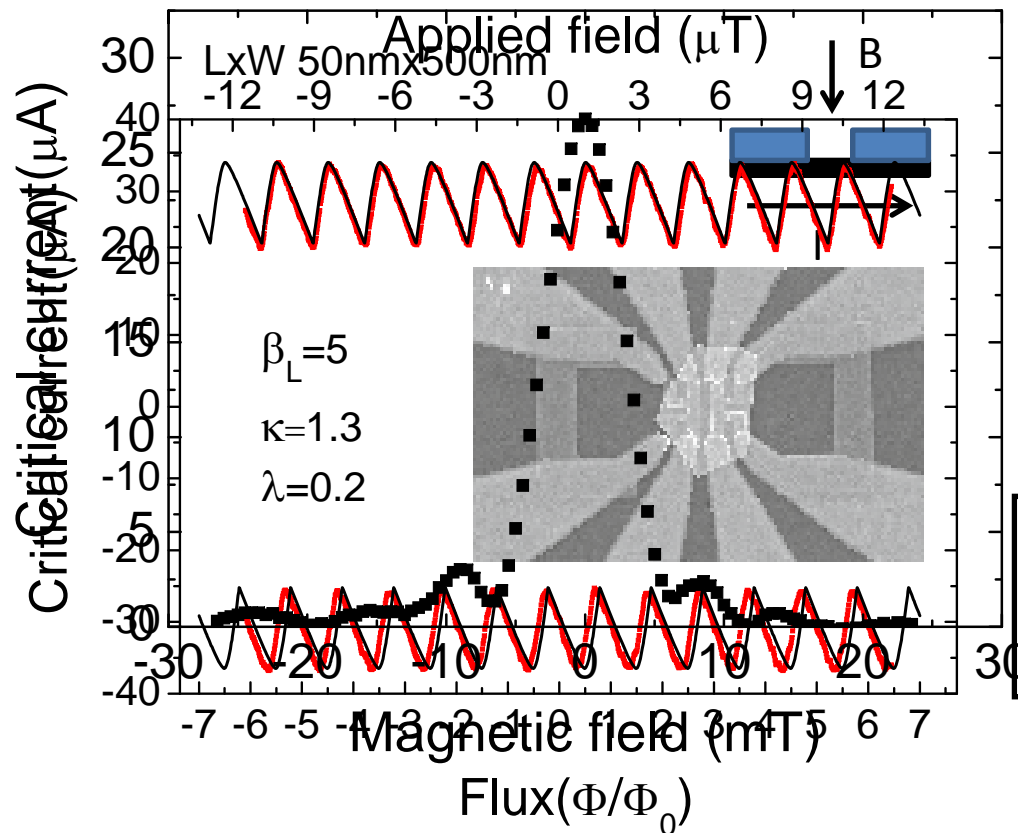
# First hallmark – Shapiro steps



Fitting Shapiro steps with Bessel functions

Only can be done if  $I_c R_n$  product is larger than the position of the steps (20.7  $\mu\text{V}$  in these measurements)

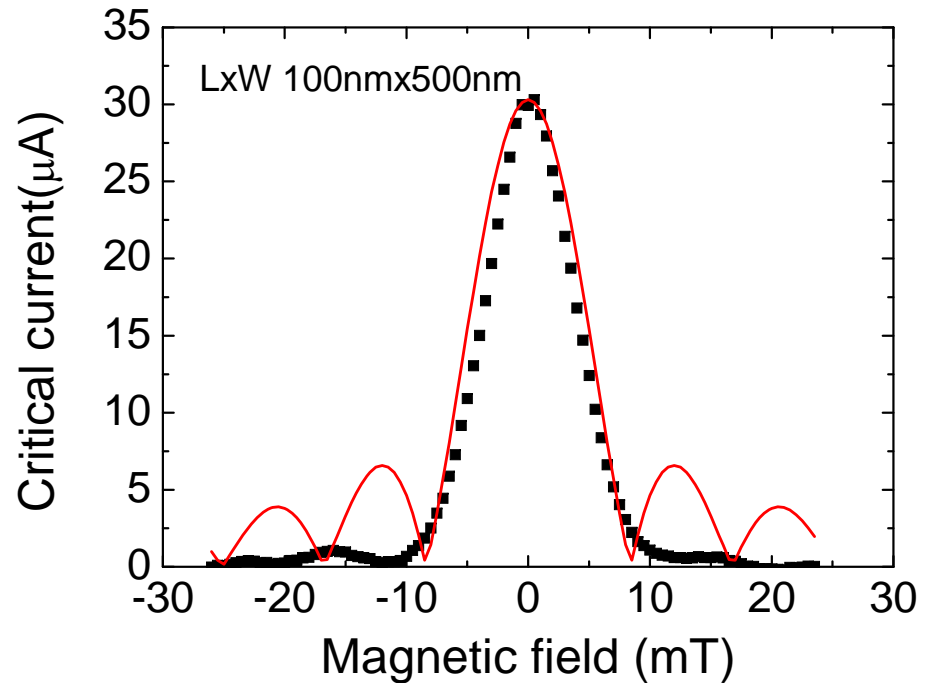
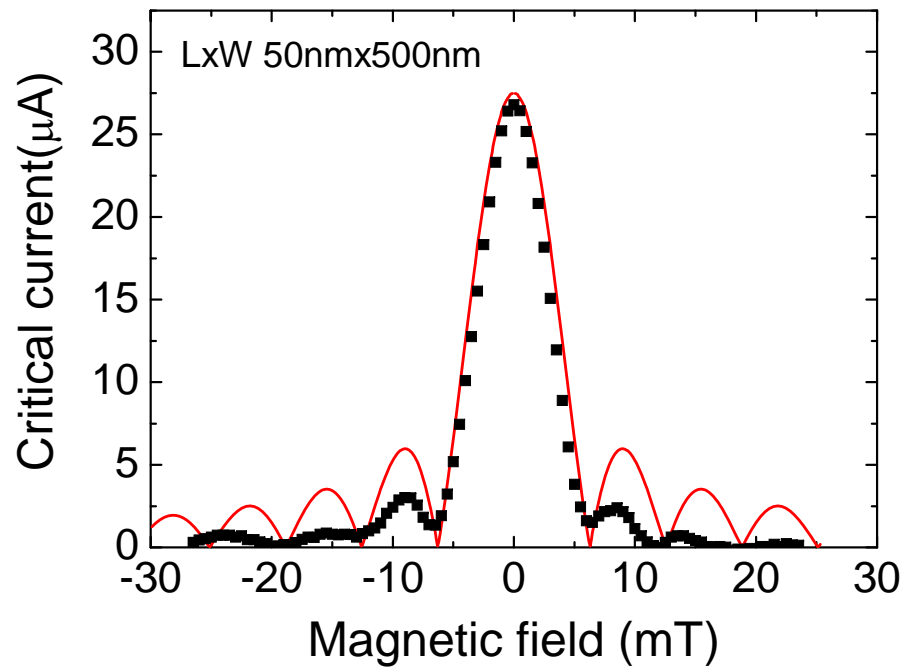
# Second hallmark – $I_c$ -B modulation



Modulation  $I_c$   
DC SQUID oscillations

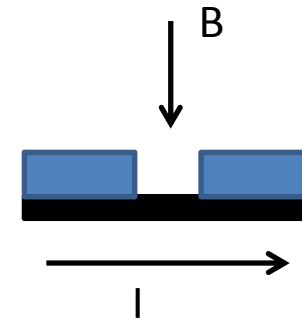
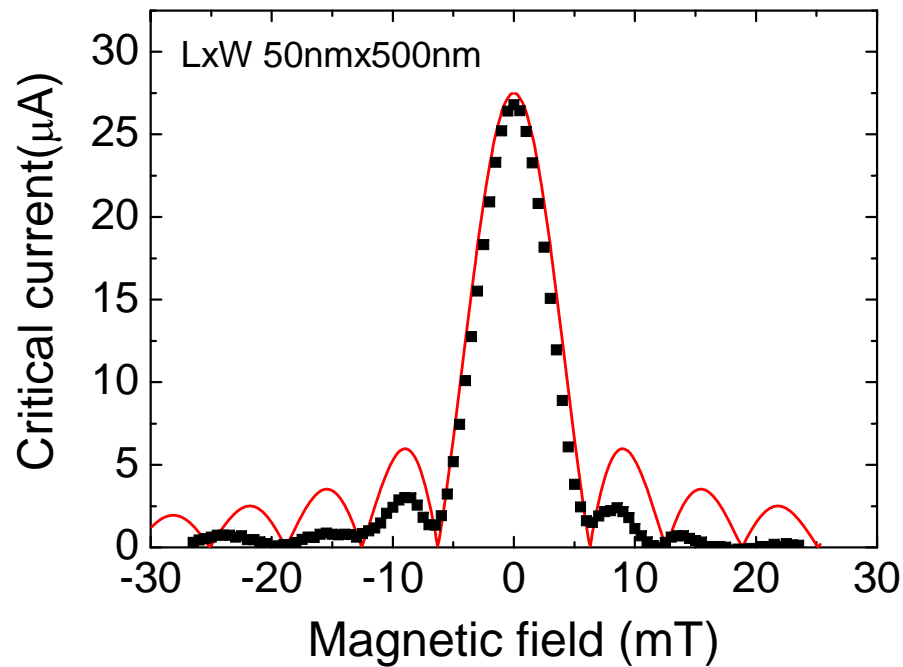
Josephson supercurrent  
present

# Second hallmark - $I_c$ -B modulation





# Second hallmark – $I_c$ -B modulation



Area uncertain:

- Penetration depth
- Flux focussing

Sinc function only valid  
for large L and W ratio

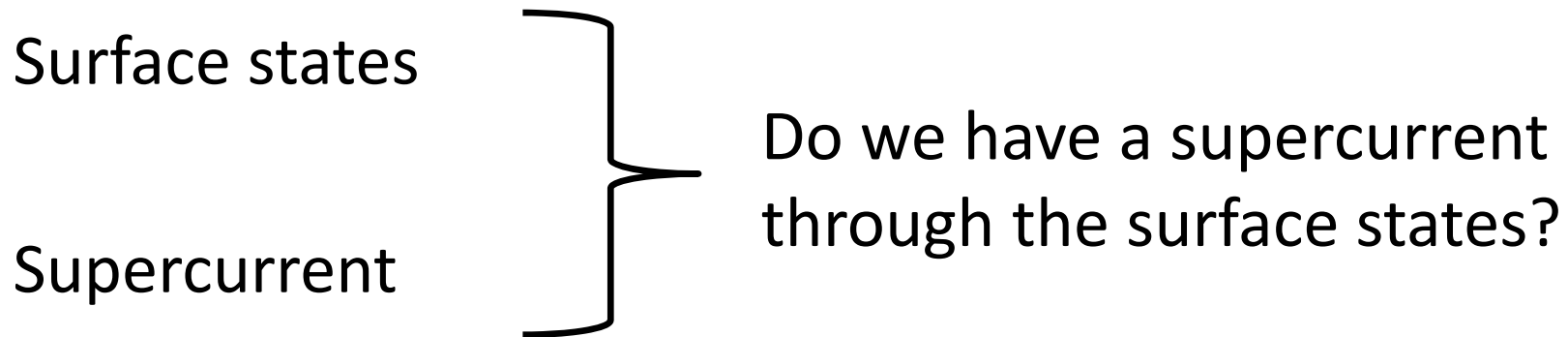
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# Link supercurrent and surface states

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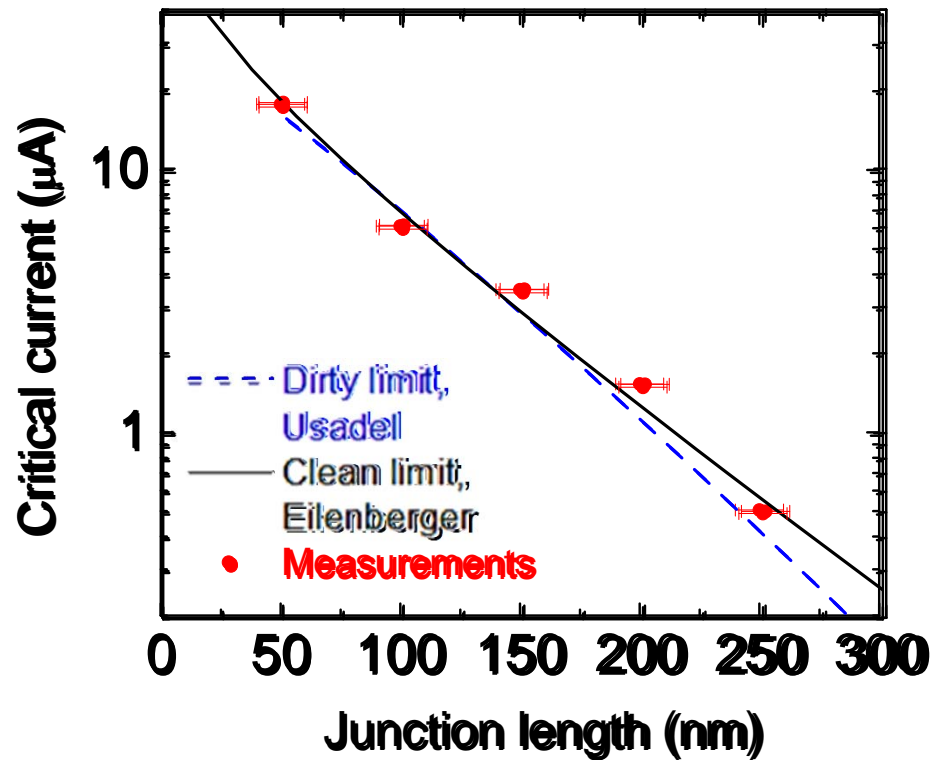


We have junctions between 50 and 250 nm and:

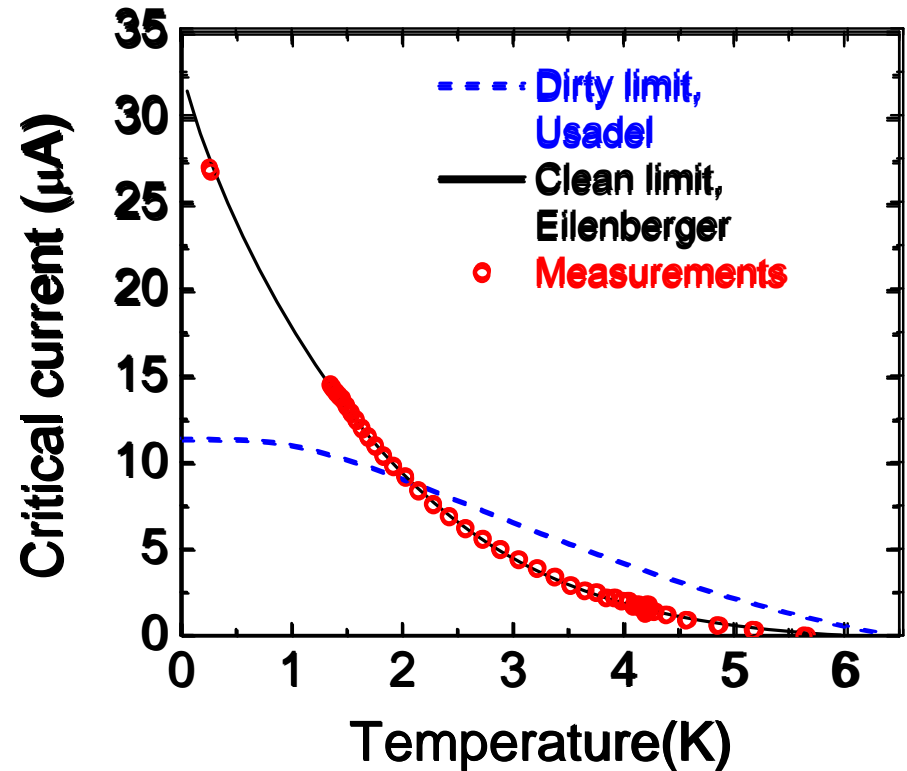
$l_{\text{mfp}} = 22$  nm (bulk states)  $\rightarrow$  diffusive transport

$l_{\text{mfp}} = 105$  nm (surface states)  $\rightarrow$  ballistic transport

# Link supercurrent and surface states



Dirty or clean?



Definitely clean

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**Josephson supercurrent has been realized  
through the surface states of  $\text{Bi}_2\text{Te}_3$**

Provides prospects for Majorana devices.....but

What is the best topological insulator for this purpose?  
(stability, insulating in the bulk, Dirac cone in the gap)

What is the smoking gun experiment with 3D  
topological insulators to observe Majorana fermions?