



# INELASTIC SCATTERING IN LUTTINGER MODELS: A NEW PARADIGMATIC MODEL

**with**

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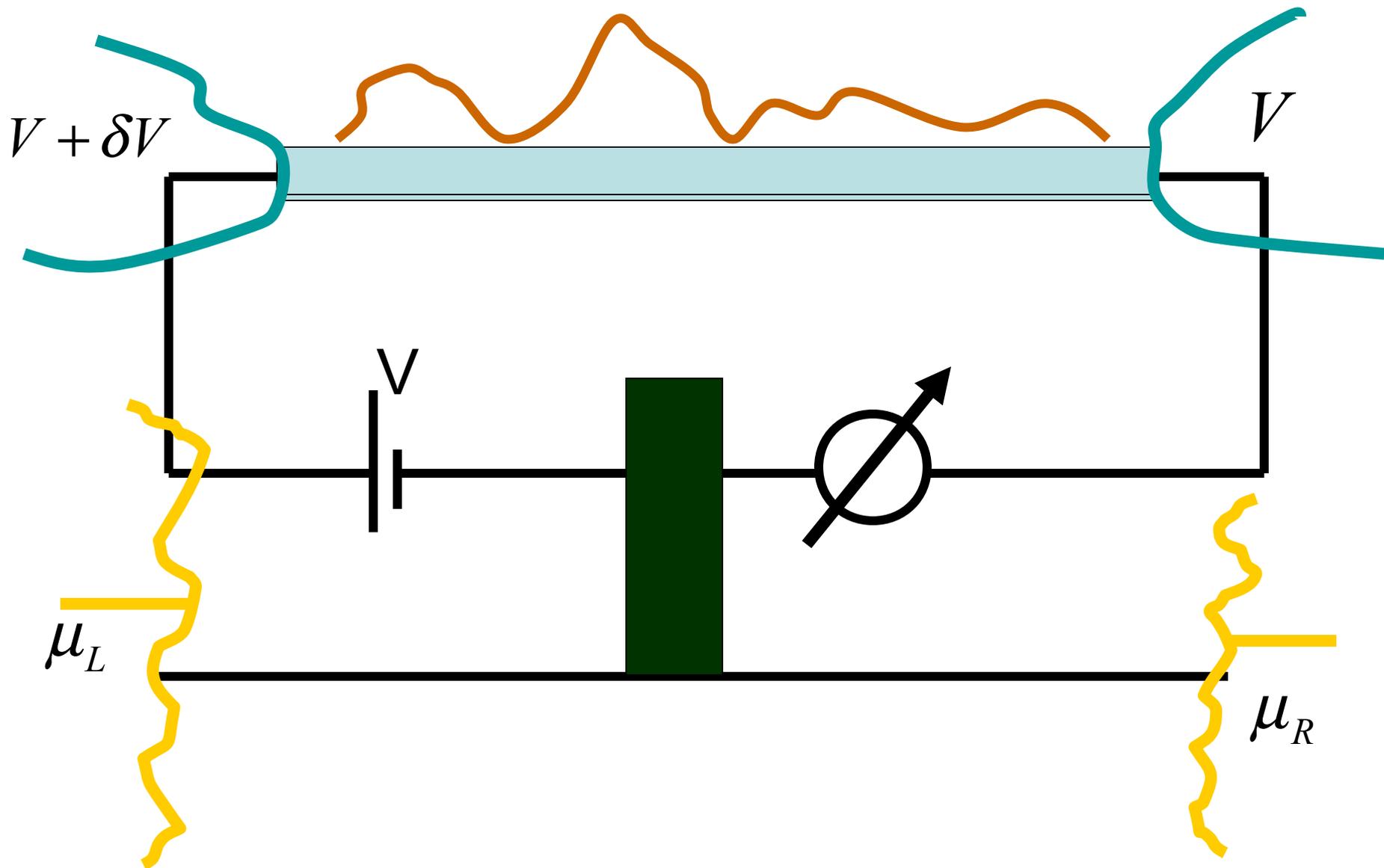
PRL 2012: Editor's selection; Focus

# **INTRODUCTION:** Noise & Transport in 1d

*Buttiker-Landauer picture*

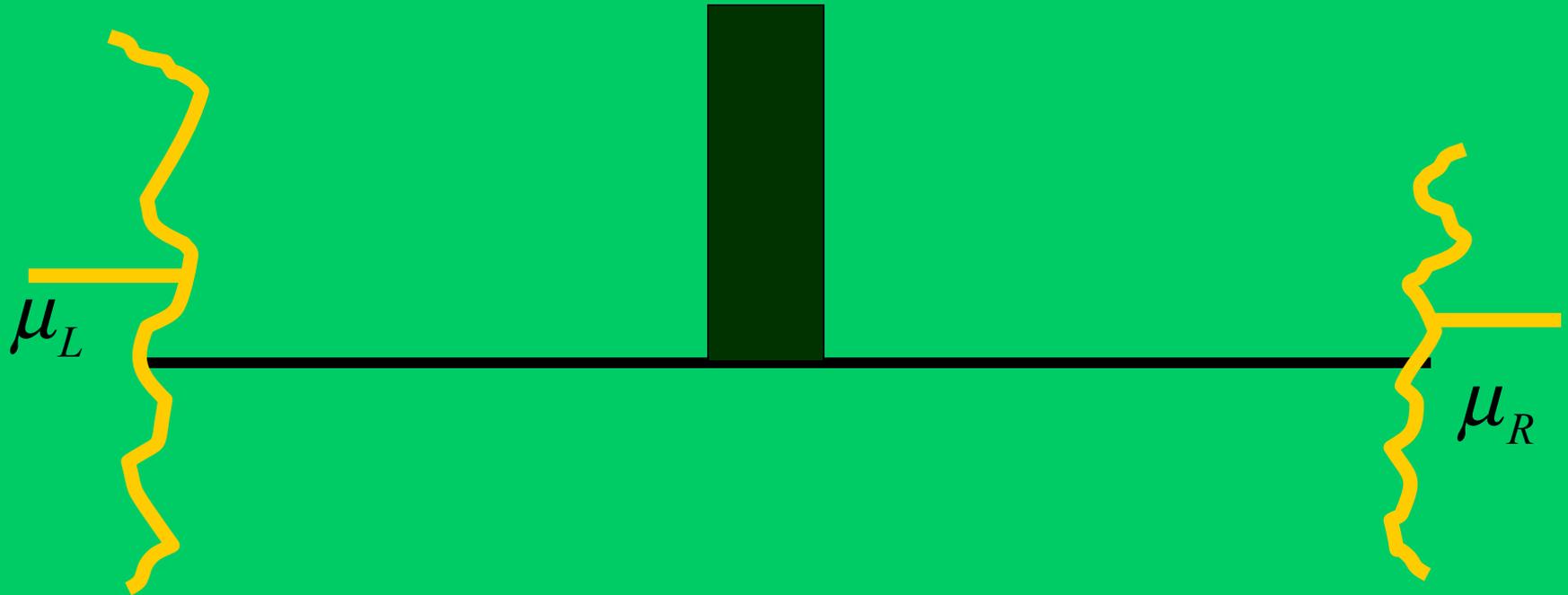
*Kane-Fisher paradigm*

# SHOT NOISE (quantum mechanical)



within the Landauer picture  
(coherent QM scatterers):

thermal noise  $\Rightarrow$  reservoirs  
shot noise  $\Rightarrow$  backscattering (partitioning)



$$\langle I_L I_L \rangle > 0; \langle I_R I_R \rangle > 0; \langle I_L I_R \rangle < 0$$

F.T.  $\langle I(t)g(t + \tau) \rangle = 2e \langle I \rangle$



**non-equilibrium shot noise**

$$S_2^{shot}(\omega = 0) = 2(\# e) \langle I \rangle$$



**FANO FACTOR**

disorder, interaction, fractional charge ...

**for fermi sea electrons:**

$$S_2 = 2e\bar{I}(1 - T)$$

**cross- current : negative**

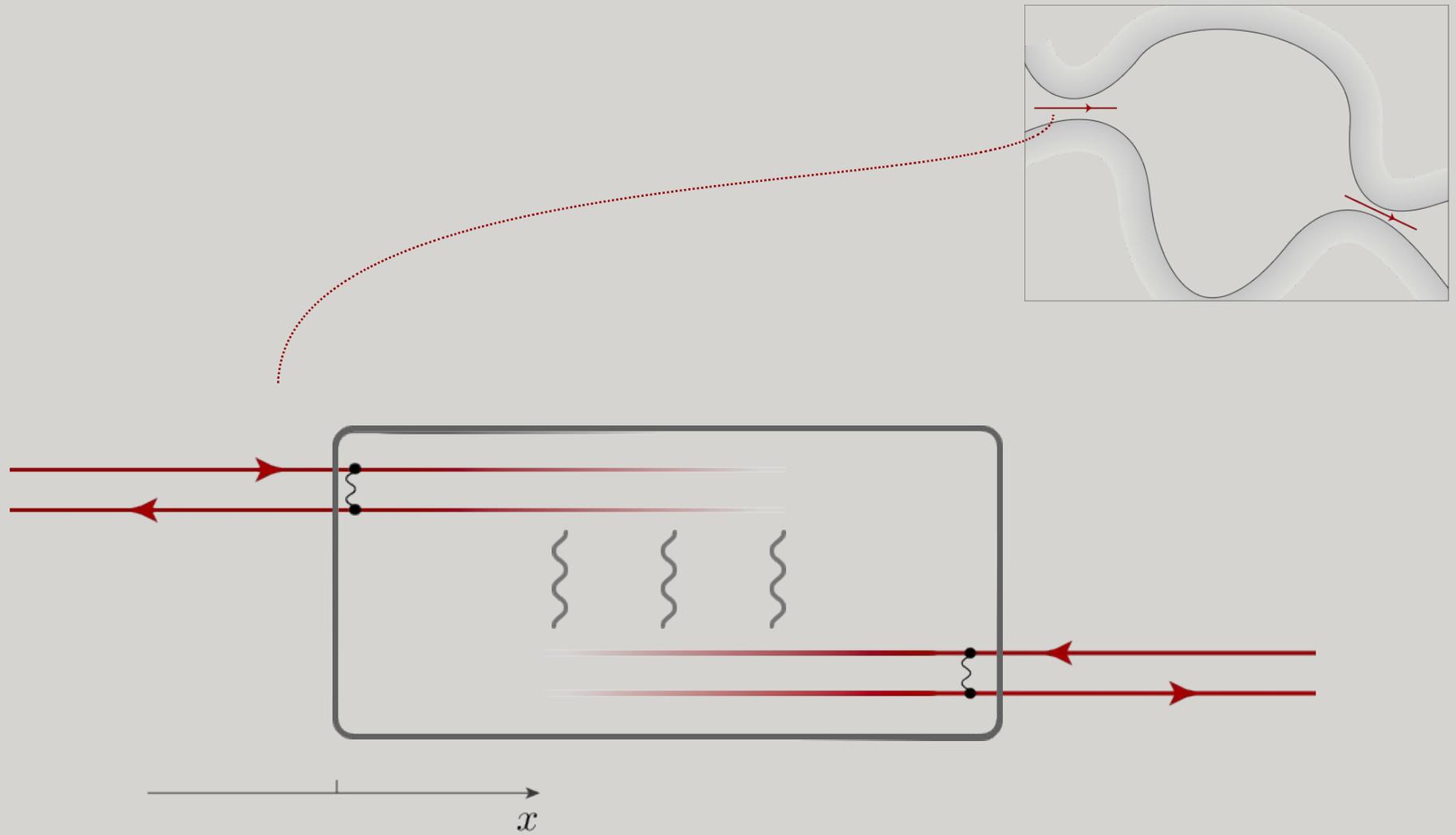
$$S_{LL} = S_{RR} = -S_{LR} = -S_{RL}$$

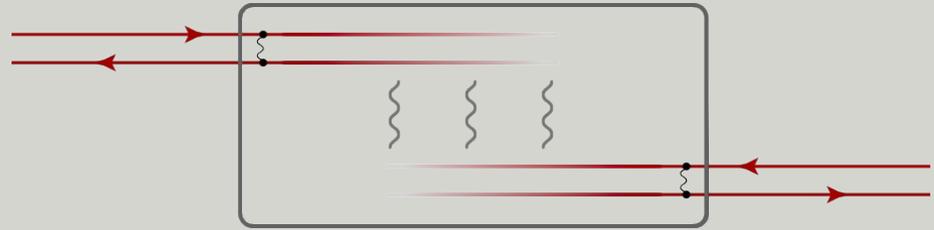
**partitioning of beam  shot noise**

# OUTLINE:

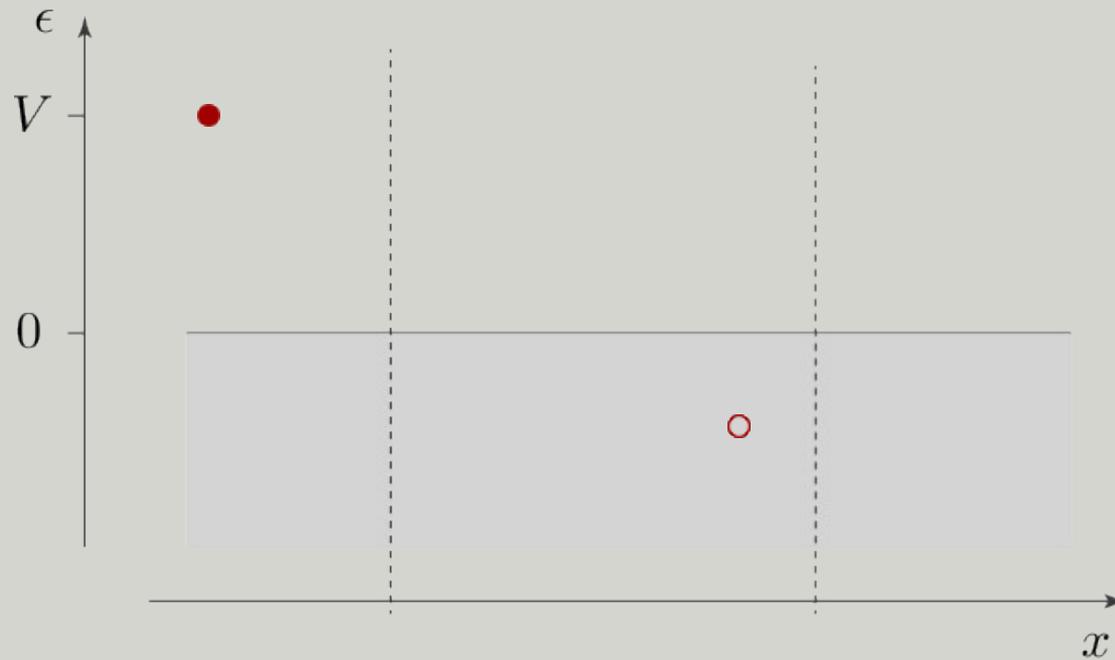
- ▷ **Matveev model**
- ▷ **generalized model**
- ▷ **scaling analysis**
- ▷ **transport properties**

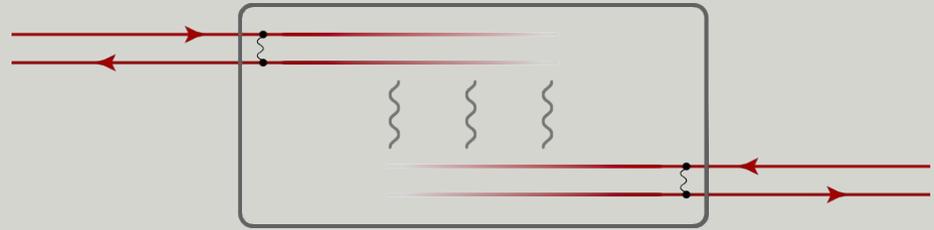
# Matveev model





- ▷ model describes interplay of backscattering and charging
- ▷ different dot/lead contacts correlated by charging interaction
- ▷ ensuing transport mechanism: **inelastic cotunneling** (cf. Kamenev, Glazman, aa, 04)



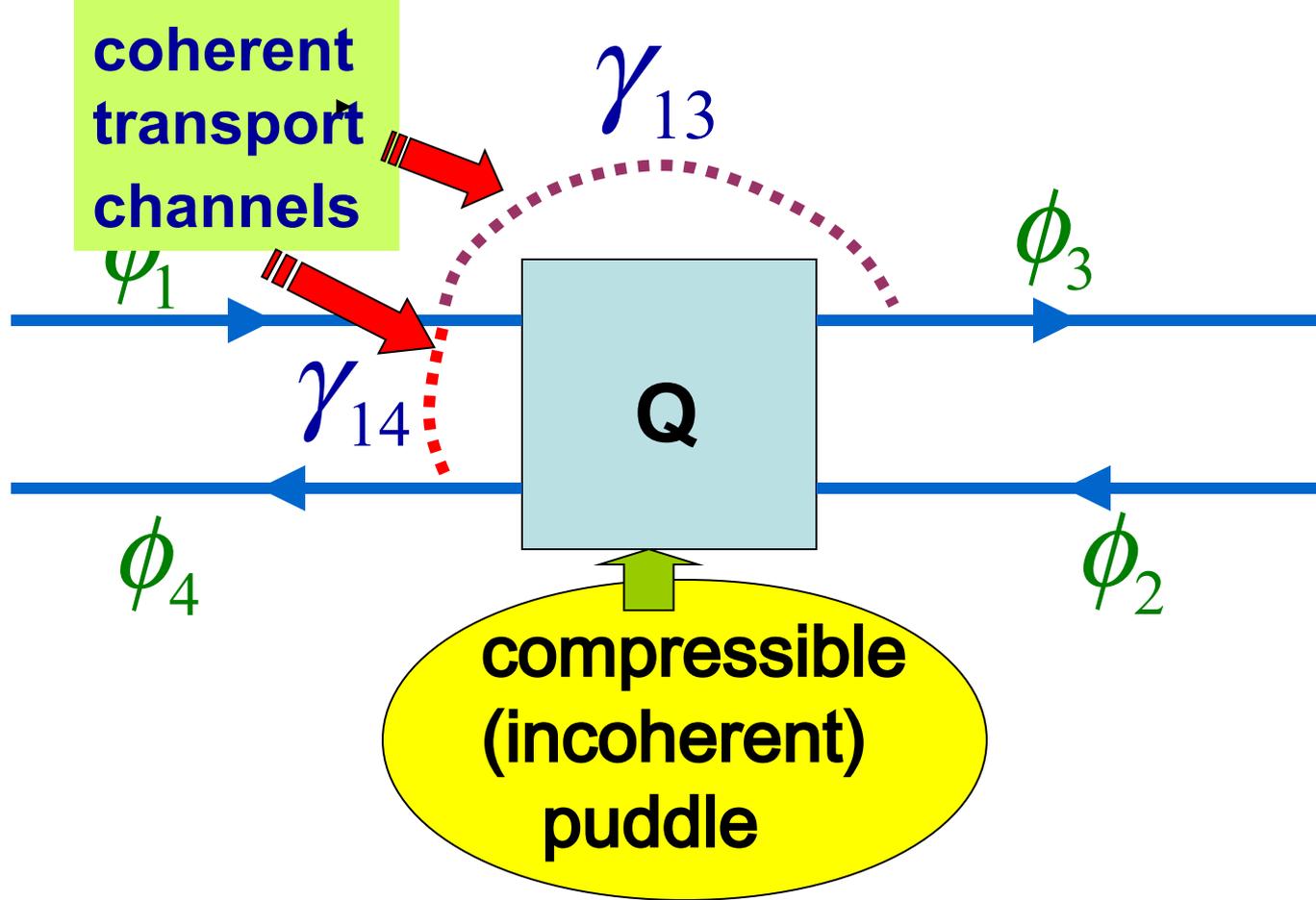


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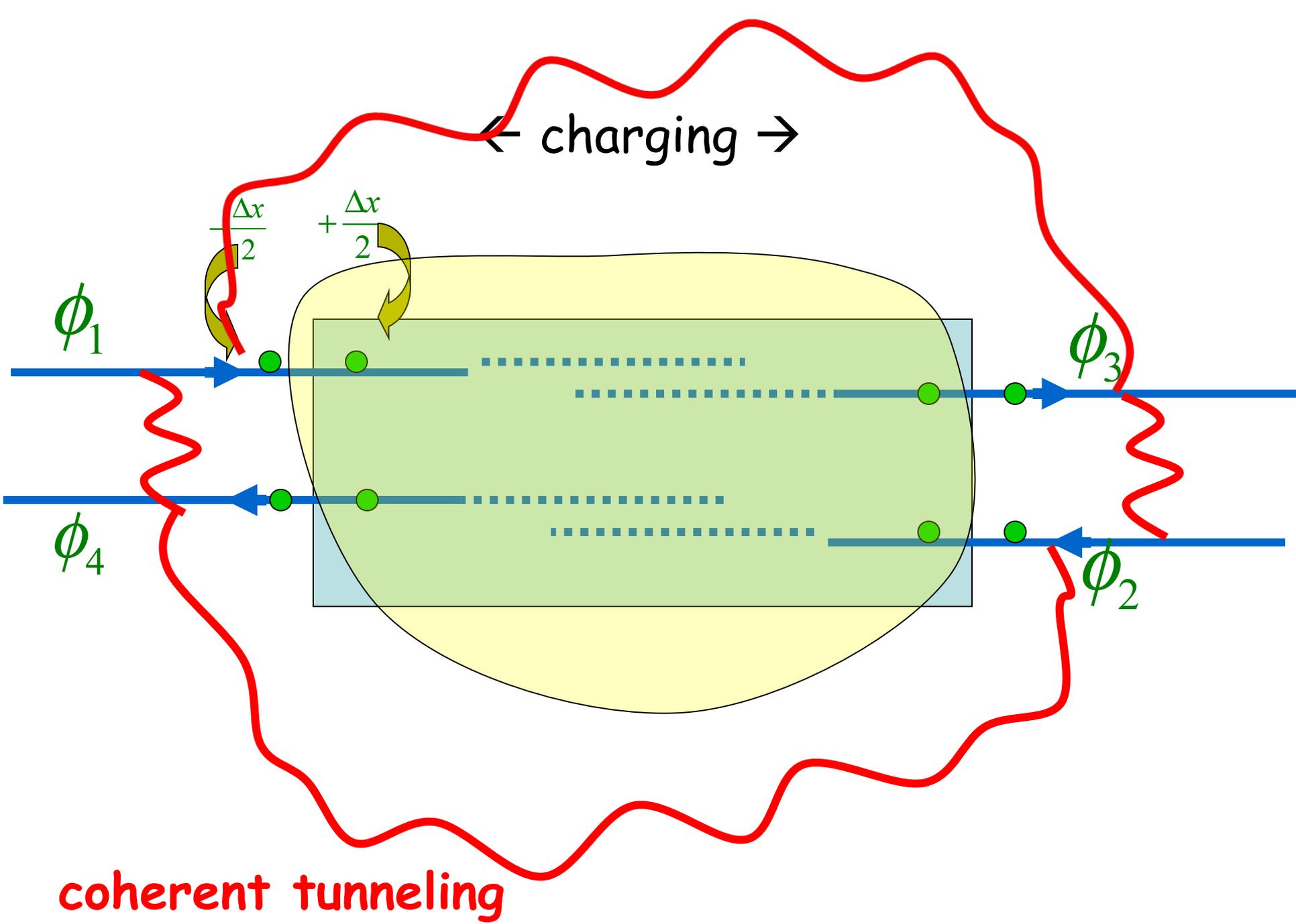
Kane/Fisher (coherent) and Matveev (incoherent) approach describe complementary channels of transport through a constriction.

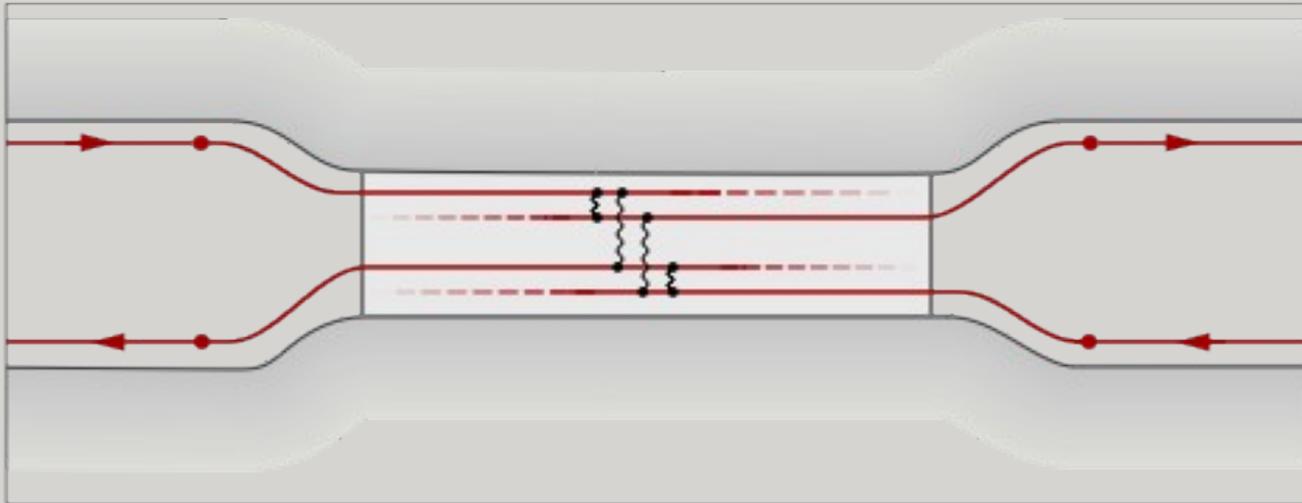
Many realistic setups may allow for a combination of both.

generalization



**ADD ELASTIC QP HOPPING**





include scattering channels between all incoming/outgoing modes.

- ▷ each chiral wire represented by '1/2' degree of freedom
- phase space of model contains **four (real) modes**.



- ▷ combination of four scattering operators and charging modes represent **three modes massive**,
- ▷ **one survivor mode** is protected by (charge conservation) symmetry.
- ▷ this mode represents the sum of the incoming currents which, by charge conservation, equals the sum of the outgoing currents.
- ▷ prevalence of this mode entails that all **transport coefficients between incoming and outgoing wires asymptote to value 1/2**.

▷ four chiral modes



$$S_0[\phi] = \frac{1}{2} \sum_{l=1}^4 \sum_{\Sigma} dx dt \sum_{i,q} \sum_{\Sigma_x} s_l \sum_{\Sigma_t} \sum_{c} S_{\text{dist}}[\phi_q]$$

$\uparrow$   $s_l = \pm 1$  for  $l = 1, 2 / 3, 4$

$\downarrow$  fixes distribution

$$S_l^{(0)}[\phi] = \frac{\hbar}{2} \int \frac{dq}{2\pi} \int \frac{d\omega}{2\pi} \phi_{i,q,\omega}^T \hat{G}_{q,\omega}^{-1} \phi_{i,-q,-\omega}$$

▷ charge in the dot

$$\frac{Q}{\Sigma} \Sigma \sum_l s_l \dot{q}_l$$

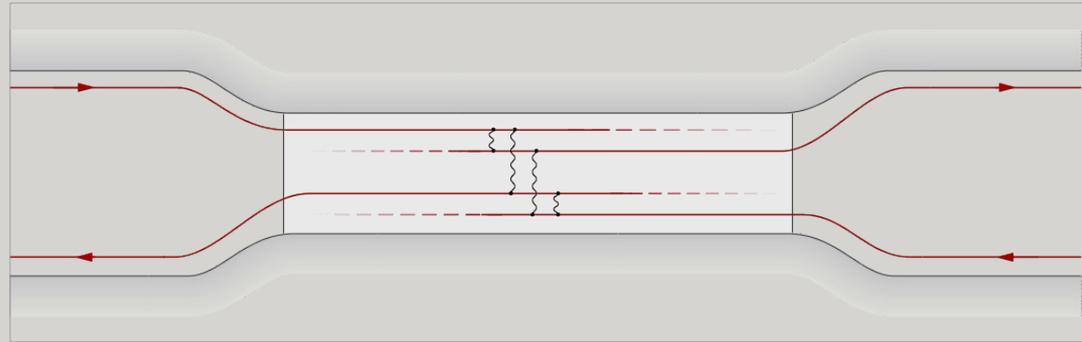
$\uparrow$   
 $\Sigma \dot{q}_l(0)$



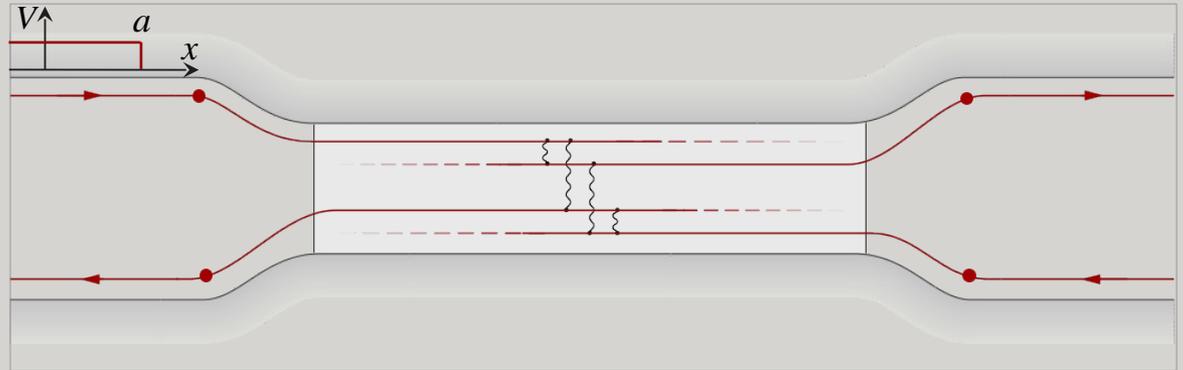
▷ charging action

$$\Sigma s_c \dot{q}_c \Sigma \frac{1}{2C} \Sigma dt Q_c Q_q$$

▷ add quasi-particle hopping at scattering hot-spots



$$S_{\text{scatt}}[\dots] \sum_{i \Sigma} \sum_{o \Sigma} \sum_{s \Sigma} \sum_{\Sigma} \dots_{i,o} \sum dt \cos \sum_{\Sigma} \dots_{i,c} \sum_{\Sigma} \dots_{o,c} \sum \frac{S}{2} \sum_{\Sigma} \dots_{i,q} \sum_{\Sigma} \dots_{o,q} \sum_{\Sigma}$$



▷ couple wire #1 to voltage kink

$$S_{\text{W}}[\rho] = \sum_{\sigma} \int dt \sum_{i,j} \rho_{i,j}^{\sigma}(a) \sum_{\sigma} \rho_{i,j}^{\sigma}(a) \mathcal{V}(t)$$

▷ current at observation points:  $j_i = \frac{1}{2} \sum_{\sigma} \rho_{i,i}^{\sigma}(a)$  (assuming gauge  $\sum_{\sigma} \rho_{i,i}^{\sigma}(a) = 0$ )

▷ strategy: integrate over field fluctuations outside tunneling hot spots

▷ effective action:

$$\int_{\mathcal{C}} \sum S_{\text{diss}} \sum S_c \sum S_{\text{scatt}} \sum S_V$$

$$\sum_{\mathcal{C}} S_{\text{diss}}[\mathcal{C}] \sum \frac{1}{2} \sum_{\Sigma} \sum_i \int dt \mathcal{C}_{i,c} \mathcal{C}_{i,q} \sum S_{\text{dist}}[\mathcal{C}_q]$$

$$\sum_{\mathcal{C}} S_c[\mathcal{C}] \sum \frac{1}{2C} \int dt (\mathcal{C}_1 \mathcal{C}_2 \mathcal{C}_3 \mathcal{C}_4)_c (\mathcal{C}_1 \mathcal{C}_2 \mathcal{C}_3 \mathcal{C}_4)_q$$

$$\sum_{\mathcal{C}} S_{\text{scatt}}[\mathcal{C}] \sum \sum_{i \in \{1,2\}} \sum_{o \in \{3,4\}} \sum_{s \in \Sigma} \mathcal{C}_{i,o} \int dt \cos \sum_{i,c} \mathcal{C}_{o,c} \sum \frac{S}{2} \sum_{i,q} \mathcal{C}_{o,q}$$

$$\sum_{\mathcal{C}} S_V[\mathcal{C}] \sum \int dt \mathcal{C}_{1,q} V$$

▷ action contains one exact zero mode (by **gauge invariance**)

$$\sum_{\vec{r}_0} \frac{1}{\sqrt{2}} (\vec{r}_1 \sum_{\vec{r}_2} \sum_{\vec{r}_3} \sum_{\vec{r}_4})$$

▷ express action in terms of the linear combinations

$$\sum_{\vec{r}_0} \frac{1}{\sqrt{2}} (\vec{r}_1 \sum_{\vec{r}_2} \sum_{\vec{r}_3} \sum_{\vec{r}_4})$$

average of incoming and outgoing current

$$\sum_{\vec{r}_1} \frac{1}{\sqrt{2}} (\vec{r}_1 \sum_{\vec{r}_2} \sum_{\vec{r}_3} \sum_{\vec{r}_4})$$

total charge in scattering region

$$\sum_{\vec{r}_2} \vec{r}_1 \sum_{\vec{r}_2}$$

difference of incoming currents

$$\sum_{\vec{r}_3} \vec{r}_3 \sum_{\vec{r}_4}$$

difference of outgoing currents

▷ express action in terms of the linear combinations

$$\sum_{\vec{r}} S_{\text{diss}} \sum_{\vec{r}_c} S_c \sum_{\vec{r}_{\text{scatt}}} S_{\text{scatt}} \sum_{\vec{r}_V} S_V$$

$$S_{\text{diss}}[\vec{r}] \sum_{\vec{r}} \frac{1}{2} \sum_{\vec{r}_l} \sum_{\vec{r}_c} dt \vec{r}_{l,c} \sum_{\vec{r}_q} S_{\text{dist}}[\vec{r}_q]$$

$$S_{\text{sc}}[\vec{r}] \sum_{\vec{r}} \frac{1}{2C} \sum_{\vec{r}_1} dt (\vec{r}_1 \sum_{\vec{r}_2} \sum_{\vec{r}_3} \sum_{\vec{r}_4})_c (\vec{r}_1 \sum_{\vec{r}_2} \sum_{\vec{r}_3} \sum_{\vec{r}_4})_q$$

$$S_{\text{scatt}}[\vec{r}] \sum_{\vec{r}} \sum_{i \in \{1,2\}} \sum_{o \in \{3,4\}} \sum_{s \in \Sigma} v_{i,o} \sum_{\vec{r}} dt \cos \frac{\vec{r}_{i,c} \cdot \vec{r}_{o,c}}{2} \sum_{\vec{r}_q} S_{\vec{r}_q}$$

$$S_V[\vec{r}] \sum_{\vec{r}} dt \vec{r}_{1,q} V$$

$$\mathcal{S} = \sum S_{\text{diss}} + \sum S_c + \sum S_{\text{scatt}} + \sum S_V$$

$$\sum S_{\text{diss}}[\mu] = \sum \frac{1}{2} \sum_a \sum_c \int dt \sum_q S_{\text{dist}}[\mu_{a,c,q}]$$

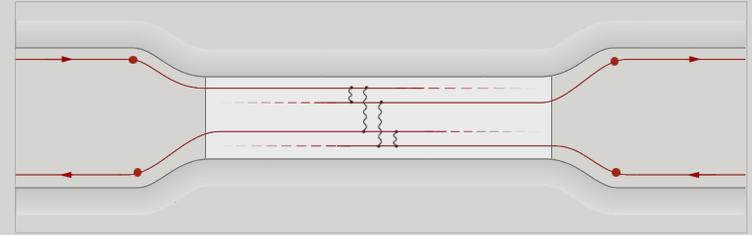
$$\sum S_c[\mu] = \frac{1}{C} \sum \int dt \sum_{c,q} S_{1,c,q}$$

$$\sum S_{\text{scatt}}[\mu] = \sum S_{\text{scatt}}[\mu_1, \mu_2, \mu_3]$$

$$\sum S_V[\mu] = \sum \int dt \frac{1}{\sqrt{8}} (\sum_0 \sum_1) \sum_2 \frac{1}{2} \sum_3$$

transport

- ▷ after integration over bulk field fluctuations and discarding massive modes: moments of currents in incoming/outgoing leads become correlation functions of zero mode.



- ▷ average outgoing current:

$$I_{\rho}(\Sigma) = 2i \sum_{c, \Sigma} \left\langle F_{\Sigma}^{c, \Sigma} \right\rangle \sum_{q, \Sigma} \frac{V_{\Sigma}}{2\Sigma} \quad \text{distribution function}$$

- ▷ equals 1/2 of the response current of an uninterrupted wire.
- ▷ corrections scale as  $\Sigma T^{2(\Sigma-1)}$
- ▷ **low temperature limit of scattering center represents (1/2,1/2) ‘beam splitter’.**

▷ current correlators

$$\sum_{\mu, \nu} X_{\mu\nu} \sum_{l, c, \Sigma} j_{l, c, \Sigma} j_{l, c, \Sigma} \sum_{\Sigma} \sum_{\Sigma} \sum_{\Sigma} \sum_{\Sigma} \sum_{\Sigma} \sum_{\Sigma}$$

$$\sum_{\mu, \nu} X_{\mu\nu} \sum_{\Sigma} \frac{1}{2 \sum_{\Sigma}} (\sum_{\Sigma} G(\sum_{\Sigma} x_l) (G(0))^{\sum_{\Sigma} - 1})_{c, \Sigma} \sum_{\Sigma} \frac{\sum_{\Sigma} V_{\Sigma}}{2 \sum_{\Sigma}}$$

Green function between system and observation point

▷ noise correlators  $\sum_{\mu, \nu} X_{\mu\nu} \sum_{\Sigma} \langle I_l \rangle \langle I_{l\Sigma} \rangle$

$$\sum_{\mu, \nu} X_{\mu\nu} \sum_{\Sigma} 2 \sum_{\Sigma} \sum_{\Sigma} F_{\Sigma} \qquad S_{o, o, \omega} = 2\pi\nu \omega F_{\omega}$$

$$S_{1, o, \omega} = 2\pi\nu \frac{\omega F_{\omega}}{2}, \quad o = 3, 4$$

all other noise correlators vanish.

## Equilibrium Noise

intra-wire noise  $X_{ll,\omega} = \frac{\nu}{2\pi} \omega F_\omega$   
outgoing/outgoing  $X_{oo',\omega} = 0$ ,  
incoming/outgoing  $X_{io,\omega} = \frac{\nu}{4\pi} \omega F_\omega$ ,  
incoming/incoming  $X_{ii',\omega} = 0$ .

## Shot Noise

no intra shot noise

no cross-correlation shot noise



**partitioning noise/ no negative cross correlations**

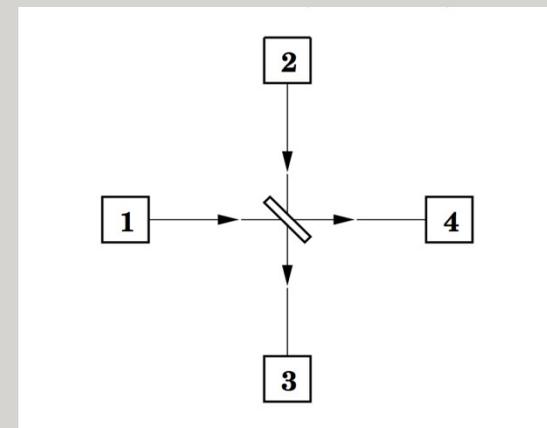


## cf: shot noise in quantum diffractive beam splitters

▷ cross-correlations between different outgoing leads  
generically negative (**partition noise**)

$$S_{ll',\omega} = \frac{1}{2} \int dt e^{i\omega t} \langle I_l(t)I_{l'}(0) + I_l(0)I_{l'}(t) \rangle$$

negative for  $l \neq l'$  (cf. Blanter & Büttiker, 99.)



summary

▷ conspiracy of elastic and inelastic scattering channels stabilizes  $(1/2, 1/2)$  low frequency fixed point.

▷ a robust phenomenon,

▷ however, realization of scaling limit depends on device setup.

p-n graphene; QHE strip + puddle

▷ asymptotic transport regime exhibits no shot noise (absence of partition noise.)



THANK YOU